Return by tue 28.9. by 2 pm.

- 1. Calculate the pair-correlation function $g(\mathbf{r}) = g_{\uparrow\uparrow}(\mathbf{r}) + g_{\uparrow\downarrow}(\mathbf{r})$ for jellium.
 - (a) Start from definition

$$g_{\sigma_1 \sigma_2}(\mathbf{r}_1, \mathbf{r}_2) = \frac{N(N-1)}{n(\mathbf{r}_1)n(\mathbf{r}_2)} \sum_{\sigma_3 \dots \sigma_N} \int \dots \int |\Psi(\mathbf{r}_1 \sigma_1, \mathbf{r}_2 \sigma_2, \dots, \mathbf{r}_N \sigma_N)|^2 d\mathbf{r}_3 \dots d\mathbf{r}_N$$
(1)

and show that for a single Slater determinant we have

$$g_{\sigma_{1}\sigma_{2}}(\mathbf{r}_{1}, \mathbf{r}_{2}) = \frac{1}{n(\mathbf{r}_{1})n(\mathbf{r}_{2})} \sum_{i,j} \{|\psi_{i}(\mathbf{r}_{1}\sigma_{1})|^{2} |\psi_{j}(\mathbf{r}_{2}\sigma_{2})|^{2} - \psi_{i}^{*}(\mathbf{r}_{1}\sigma_{1})\psi_{j}^{*}(\mathbf{r}_{2}\sigma_{2})\psi_{j}(\mathbf{r}_{1}\sigma_{1})\psi_{i}(\mathbf{r}_{2}\sigma_{2})$$
$$- \psi_{i}^{*}(\mathbf{r}_{1}\sigma_{1})\psi_{i}^{*}(\mathbf{r}_{2}\sigma_{2})\psi_{i}(\mathbf{r}_{1}\sigma_{1})\psi_{j}(\mathbf{r}_{2}\sigma_{2}) + |\psi_{j}(\mathbf{r}_{1}\sigma_{1})|^{2} |\psi_{i}(\mathbf{r}_{2}\sigma_{2})|^{2} \}. \tag{2}$$

Alternatively you can just verify this result in case of N=3.

(b) In the jellium the single-electron states are $\psi_j(\mathbf{r}\sigma) = \exp(i\mathbf{k}_j \cdot \mathbf{r})/\sqrt{V}\chi_{s_i}(\sigma)$ and for the spinor $\chi_{s_i}(\sigma)$ we have $\chi_{s_i}^*(\sigma)\chi_{s_j}(\sigma) = \delta_{s_is_j}$. Furthermore the density n of the jellium is constant N/V. Put this information into eq. (2) and show that

$$g_{\uparrow\downarrow}(\mathbf{r}_1, r_2) = \frac{1}{2}$$

$$g_{\uparrow\uparrow}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{2}[1 - \phi(\mathbf{r}_1 - \mathbf{r}_2)^2],$$

where

$$\phi(\mathbf{r}) = \frac{3}{(rk_F)^3} [\sin(rk_F) - (rk_F)\cos(rk_F)] = \frac{3}{rk_F} j_1(rk_F)$$

Use handy formula $\sum_{\mathbf{k}} F(\mathbf{k}) = V/(2\pi)^3 \int F(\mathbf{k}) d\mathbf{k}$.

(c) Define exchange hole density $n_x(\mathbf{r}) = N/V(g(\mathbf{r}) - 1)$ and show that it satisfies the sum rule

$$\int n_x(\mathbf{r}) d\mathbf{r} = -1$$

and that the Coulomb interaction energy between an electron and $n_x(\mathbf{r})$,

$$\int \frac{e^2 n_x(\mathbf{r})}{r} d\mathbf{r},$$

yields the exchange energy of the jellium.

Hints: $\int_0^\infty j_1(x)^2 dx = \pi/6$ and $\int_0^\infty j_1(x)^2/x dx = 1/8$ (3 points)

2. Calculate Fourier transformation $\mathcal{F}[\phi_s(\mathbf{r})]$ of the screened Coulomb potential (Yukawa potential) $\phi_s(\mathbf{r}) = -e \exp(-k_0 r)/r$, $k_0 > 0$:

$$\mathcal{F}[\phi_s(\mathbf{r})] := \tilde{\phi}_s(\mathbf{k}) = \int e^{-i\mathbf{k}\cdot\mathbf{r}} \phi_s(\mathbf{r}) d\mathbf{r}.$$

What is the Fourier transformation of the bare (unscreened) Coulomb potential $\phi(\mathbf{r})$? Check your result by solving the Poisson's equation of a point charge,

$$\nabla^2 \phi(\mathbf{r}) = 4\pi e \delta(\mathbf{r}),$$

with Fourier transformation. Recall that $\mathcal{F}[\nabla^2 \phi(\mathbf{r})] = -k^2 \tilde{\phi}(\mathbf{k})$. (1 point)

3. Show that near the band minimum ($\mathbf{k}=0$) the Hartree-Fock one-electron energy is parabolic in k:

$$\varepsilon(\mathbf{k}) pprox rac{\hbar^2 k^2}{2m^*},$$

where

$$\frac{m^*}{m} = \frac{1}{1 + 0.22(r_s/a_0)}.$$

(2 points)