

*Return to Michael by Wednesday 18.10.2006 at noon*

**1.** Calculate the pair-correlation function  $g(r) = g_{\uparrow\uparrow}(r) + g_{\uparrow\downarrow}(r)$  for jellium.

(a) Start from definition

$$g_{\sigma_1\sigma_2}(\mathbf{r}_1, \mathbf{r}_2) = \frac{N(N-1)}{n(\mathbf{r}_1)n(\mathbf{r}_2)} \sum_{\sigma_3 \dots \sigma_N} \int \dots \int |\Psi(\mathbf{r}_1\sigma_1, \mathbf{r}_2\sigma_2, \dots, \mathbf{r}_N\sigma_N)|^2 d\mathbf{r}_3 \dots d\mathbf{r}_N \quad (1)$$

and show that for a single Slater determinant we have

$$\begin{aligned} g_{\sigma_1\sigma_2}(\mathbf{r}_1, \mathbf{r}_2) &= \frac{1}{n(\mathbf{r}_1)n(\mathbf{r}_2)} \sum_{i,j} \{ |\psi_i(\mathbf{r}_1\sigma_1)|^2 |\psi_j(\mathbf{r}_2\sigma_2)|^2 - \psi_i^*(\mathbf{r}_1\sigma_1) \psi_j^*(\mathbf{r}_2\sigma_2) \psi_j(\mathbf{r}_1\sigma_1) \psi_i(\mathbf{r}_2\sigma_2) \\ &\quad - \psi_j^*(\mathbf{r}_1\sigma_1) \psi_i^*(\mathbf{r}_2\sigma_2) \psi_i(\mathbf{r}_1\sigma_1) \psi_j(\mathbf{r}_2\sigma_2) + |\psi_j(\mathbf{r}_1\sigma_1)|^2 |\psi_i(\mathbf{r}_2\sigma_2)|^2 \}. \end{aligned} \quad (2)$$

Alternatively you can just verify this result in case of  $N = 3$ .

(b) In the jellium the single-electron states are  $\psi_j(\mathbf{r}\sigma) = \exp(i\mathbf{k}_j \cdot \mathbf{r})/\sqrt{V} \chi_{s_i}(\sigma)$  and for the spinor  $\chi_{s_i}(\sigma)$  we have  $\chi_{s_i}^*(\sigma) \chi_{s_j}(\sigma) = \delta_{s_i s_j}$ . Furthermore the density  $n$  of the jellium is constant  $N/V$ . Put this information into eq. (2) and show that

$$\begin{aligned} g_{\uparrow\downarrow}(\mathbf{r}_1, \mathbf{r}_2) &= \frac{1}{2} \\ g_{\uparrow\uparrow}(\mathbf{r}_1, \mathbf{r}_2) &= \frac{1}{2} [1 - \phi(\mathbf{r}_1 - \mathbf{r}_2)^2], \end{aligned}$$

where

$$\phi(\mathbf{r}) = \frac{3}{(rk_F)^3} [\sin(rk_F) - (rk_F) \cos(rk_F)] = \frac{3}{rk_F} j_1(rk_F)$$

Use handy formula  $\sum_{\mathbf{k}} F(\mathbf{k}) = V/(2\pi)^3 \int F(\mathbf{k}) d\mathbf{k}$ .

(c) Define exchange hole density  $n_x(\mathbf{r}) = N/V(g(\mathbf{r}) - 1)$  and show that it satisfies the sum rule

$$\int n_x(\mathbf{r}) d\mathbf{r} = -1$$

and that the Coulomb interaction energy between an electron and  $n_x(\mathbf{r})$ ,

$$\int \frac{e^2 n_x(\mathbf{r})}{r} d\mathbf{r},$$

yields the exchange energy of the jellium.

Hints:  $\int_0^\infty j_1(x)^2 dx = \pi/6$  and  $\int_0^\infty j_1(x)^2/x dx = 1/8$  **(4 points)**

**2.** (a) Derive the result for the thermal conductivity coefficient:  $\kappa_T = \frac{1}{3} c_V \tau v_F^2$ . Make a one-dimensional model which has a thermal gradient  $dT/dx$ . Consider electrons at point  $x$ . The energy of electrons depends on the position through  $T$ :  $\varepsilon = \varepsilon(T(x))$ . What is the contribution to

the thermal current density  $J_Q$  at  $x$  by electrons arriving from "left (cooler)" or "right (hotter)" regions? Expand  $J_Q$  about the point  $x$ , recognize the term consisting of the electronic heat capacity per unit volume ( $c_V$ ) and generalize the equation to 3D, which will give you the desired result.

(b) Conversely, a temperature gradient in a long, thin metal bar will induce an electric field that is directed opposed to the temperature gradient (*Seebeck effect*). That is,  $\mathbf{E} = Q \nabla T$ . By a similar consideration of a 1D model and generalization to 3D, derive the *thermopower constant*  $Q = -c_V/3ne$ . **(3 points)**

**3.** Show that near the band minimum ( $\mathbf{k}=0$ ) the Hartree-Fock one-electron energy

$$\varepsilon(k) = \frac{\hbar^2 k^2}{2m} - \frac{2e^2}{\pi} F(x)$$

where  $x = k/k_F$  and

$$F(x) = \frac{1}{4x} \left[ (1 - x^2) \ln \left| \frac{1+x}{1-x} \right| + 2x \right]$$

is parabolic in  $k$ :

$$\varepsilon(k) \approx \frac{\hbar^2 k^2}{2m^*},$$

where

$$\frac{m^*}{m} = \frac{1}{1 + 0.22(r_s/a_0)}.$$

**(3 points)**