

The EPPS16 NLO nuclear PDFs

H. Paukkunen

K. J. Eskola, P. Paakkinen, C. A. Salgado

University of Jyväskylä, Finland
Helsinki Institute of Physics, Finland
Universidade de Santiago de Compostela, Spain

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Theoretical Foundations on collinear factorization

- Factorization formula

Parton distribution functions (PDFs)

$$d\sigma = \sum_{i,j} f_i(Q_f^2) \otimes d\sigma_{ij}(Q_f^2, Q_r^2) \otimes f_j(Q_f^2) + \mathcal{O}(Q_f^{-2n})$$

The diagram illustrates the factorization formula for a cross-section $d\sigma$. The formula is $d\sigma = \sum_{i,j} f_i(Q_f^2) \otimes d\sigma_{ij}(Q_f^2, Q_r^2) \otimes f_j(Q_f^2) + \mathcal{O}(Q_f^{-2n})$. The terms $f_i(Q_f^2)$ and $f_j(Q_f^2)$ are highlighted in blue boxes, and the central term $d\sigma_{ij}(Q_f^2, Q_r^2)$ is highlighted in a red oval. A bracket labeled "Parton distribution functions (PDFs)" points to both $f_i(Q_f^2)$ and $f_j(Q_f^2)$.

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Coefficient functions (calculable by perturbative methods)

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Coefficient functions (calculable by perturbative methods)

- PDFs obey the DGLAP equations

$$Q^2 \frac{\partial f_i(x, Q^2)}{\partial Q^2} = \sum_j P_{ij}(Q^2) \otimes f_j(x, Q^2)$$

Splitting functions (calculable by perturbative methods)

Previous analyses

	HKN07	EPS09	DSSZ	NCTEQ15
Order in α_s	LO & NLO	LO & NLO	NLO	NLO
Neutral current DIS $\ell+A/\ell+d$	✓	✓	✓	✓
Drell-Yan dilepton $p+A/p+d$	✓	✓	✓	✓
RHIC pions $d+Au/p+p$		✓	✓	✓
Neutrino-nucleus DIS			✓	
Q^2 cut in DIS	1 GeV	1.3 GeV	1 GeV	2 GeV
datapoints	1241	929	1579	708
free parameters	12	15	25	17
error analysis	Hessian	Hessian	Hessian	Hessian
error tolerance $\Delta\chi^2$	13.7	50	30	35
Free proton baseline PDFs	MRST98	CTEQ6.1	MSTW2008	CTEQ6M-like
Heavy quark treatment	ZM-VFNS	ZM-VFNS	GM-VFNS	GM-VFNS
Cites	272	674	129	43

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Neutrino-nucleus DIS		✓	✓	
LHC p+Pb jet data		✓		
LHC p+Pb W, Z data		✓		
Drell-Yan dilepton $\pi+A$		✓		
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Framework: Parametrization

- Our definition of nuclear PDFs

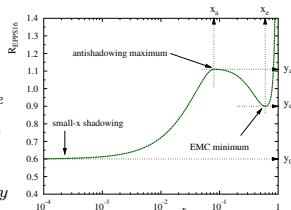
$$f_i^{p/A}(x, Q^2) \equiv R_i^A(x, Q^2) f_i^p(x, Q^2)$$

Nuclear modifications

Free proton baseline
(CT14NLO)

- Fit function x dependence at $Q_0^2 = m_{\text{charm}}^2$

$$R(x, Q_0^2) = \begin{cases} a_0 + a_1(x - x_a)^2 & x \leq x_a \\ b_0 + b_1x^\alpha + b_2x^{2\alpha} + b_3x^{3\alpha} & x_a \leq x \leq x_e \\ c_0 + (c_1 - c_2x)(1 - x)^{-\beta} & x_e \leq x \leq 1, \end{cases}$$

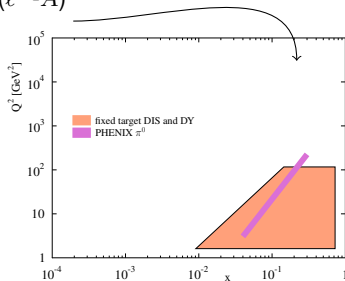


- The A dependence parametrized at $x \rightarrow 0$, x_a , and x_y
- Freedom allowed for all quark flavours separately — we no longer impose flavour-independent nuclear effects for valence and sea quarks (as e.g. in EPS09).
- 20 free parameters (15 in EPS09)

Experimental input

- EPS09:

- fixed-target deep-inelastic scattering (DIS) ($\ell^\pm$ -A)
- low-mass Drell-Yan dilepton data (p-A)
- inclusive pion production (d-Au)



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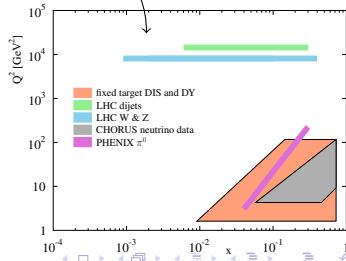
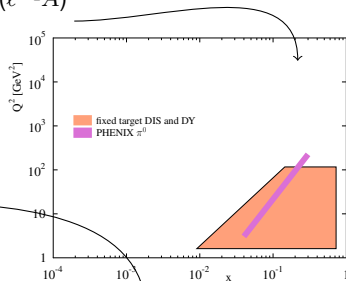
- Additional data in EPPS16:

- low-mass Drell-Yan dilepton data ($\pi^\pm$ -A)
- Neutrino DIS data (ν -Pb, $\bar{\nu}$ -Pb)
- LHC data for dijets, $W^\pm$ and Z (p-Pb)

⇒ Completely new region open in (x, Q^2) plane!

⇒ Larger variety of data than in any other contemporary analysis!

⇒ First true test of nuclear PDF universality!



Experimental input: Corrections for old $\ell^- A$ DIS data

- Recover the true structure functions from the “isoscalarized” ones ($\ell^- A$ DIS):

“Isoscalarized” structure functions
reported by the experiments
(used e.g. in EPS09 analysis):

$$\hat{F}_2^A = \frac{1}{2} F_2^{p,A} + \frac{1}{2} F_2^{n,A}$$

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The true structure functions
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$$F_2^A = \frac{Z}{A} F_2^{\text{p},A} + \frac{N}{A} F_2^{\text{n},A}$$

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Relation between the two:

$$\hat{F}_2^A = \left[\frac{A}{2} \left(1 + \frac{F_2^n}{F_2^p} \right) / \left(Z + N \frac{F_2^n}{F_2^p} \right) \right] F_2^A,$$

Can find these from the literature

⇒ Better sensitivity for flavour separation

Experimental input: CHORUS Neutrino DIS data

- The ν -Pb and $\bar{\nu}$ -Pb DIS data available as absolute cross sections $\frac{d\sigma_{i,\text{exp}}^{\nu,\bar{\nu}}}{dxdy}$
 $\implies$ Sensitive to both the free proton baseline & nuclear modifications

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⇒ Sensitive to both the free proton baseline & nuclear modifications

- To be less sensitive to the free proton PDF use normalized cross sections instead:

$$\frac{d\tilde{\sigma}_{i,\text{exp}}^{\nu,\bar{\nu}}}{dxdy} \equiv \frac{d\sigma_{i,\text{exp}}^{\nu,\bar{\nu}}}{dxdy} / \sigma_{\text{exp}}^{\nu,\bar{\nu}}(E_i),$$
$$\sigma_{\text{exp}}^{\nu,\bar{\nu}}(E) = \sum_i \frac{d\sigma_{i,\text{exp}}^{\nu,\bar{\nu}}}{dxdy} \Delta_i^{xy} \delta_{E,E_i} \approx \text{integrated xsec at fixed } E$$

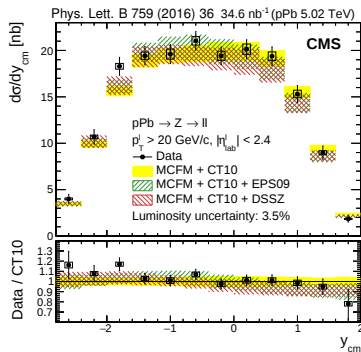
Size of the (x,y) bin (rectangle)

- Correlated systematic uncertainties propagated to the normalized cross sections

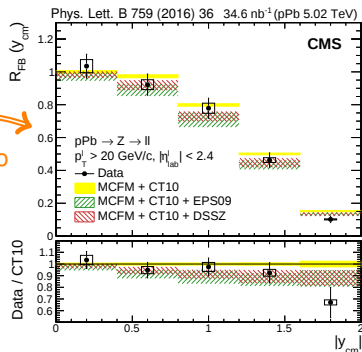
Experimental input: The LHC p-Pb data

- The LHC p-Pb data included as forward-to-backward ratios

$$R_{\text{FB}} = \frac{d\sigma(\eta > 0)}{d\sigma(\eta < 0)}$$



Take the ratio



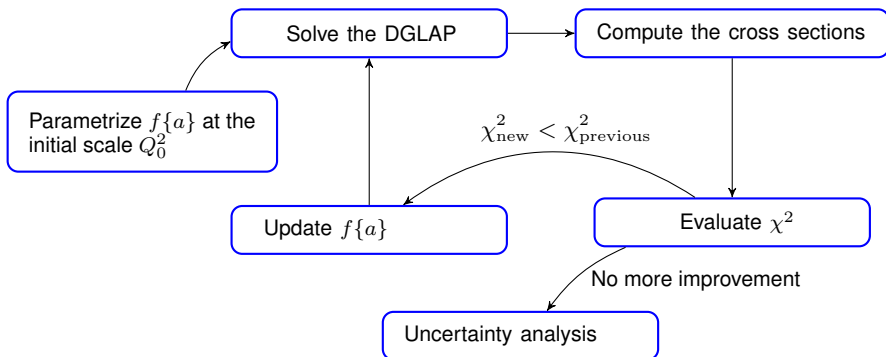
- Cancel experimental & theory uncertainties...but lose some information also
- $R_{\text{FB}} \neq 1$ for: nuclear mods in PDFs + isospin and phase-space effects

Analysis procedure

- We consider χ^2 figure-of-merit function

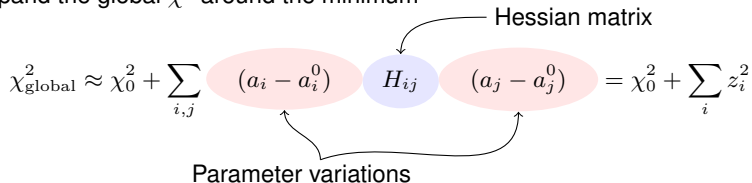
$$\chi_{\text{global}}^2 \equiv \sum_{i,j} [T_i(\vec{a}) - D_i] C_{ij}^{-1} [T_j(\vec{a}) - D_j]$$

- Minimize χ^2 by built-in Levenberg-Marquardt algorithm



Uncertainty analysis by Hessian method

- Expand the global χ^2 around the minimum



The diagram illustrates the Hessian method for uncertainty analysis. It shows the expansion of the global χ^2 around a minimum. The equation is $\chi_{\text{global}}^2 \approx \chi_0^2 + \sum_{i,j} (a_i - a_i^0) H_{ij} (a_j - a_j^0) = \chi_0^2 + \sum_i z_i^2$. In the diagram, the terms $(a_i - a_i^0)$ and $(a_j - a_j^0)$ are enclosed in red ovals and labeled "Parameter variations" with a bracket. The term H_{ij} is enclosed in a blue oval and labeled "Hessian matrix" with an arrow. The summation indices i, j and i are also present.

$$\chi_{\text{global}}^2 \approx \chi_0^2 + \sum_{i,j} (a_i - a_i^0) H_{ij} (a_j - a_j^0) = \chi_0^2 + \sum_i z_i^2$$

Parameter variations

Hessian matrix

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Parameter variations

Hessian matrix

- The z_i coordinates (linear combinations of a_i) are $\sim$ uncorrelated and one can use the standard law of error propagation

$$(\delta X)^2 = \sum_i \left(\frac{\partial X}{\partial z_i} \times \delta z_i \right)^2, \quad \delta z_i = \frac{\delta z_i^+ + \delta z_i^-}{2}$$

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- Define the PDF uncertainty sets $S_i^\pm$

$$S_1^\pm \equiv \pm \delta z_1^\pm (1, 0, \dots, 0)$$

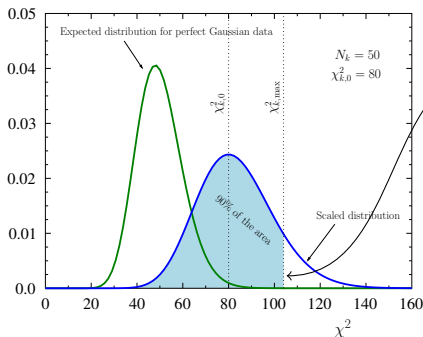
$$\vdots$$

$$S_N^\pm \equiv \pm \delta z_N^\pm (0, 0, \dots, 1)$$

$$\Rightarrow (\delta X)^2 = \frac{1}{4} \sum_i [X(S_i^+) - X(S_i^-)]^2$$

Uncertainty analysis by Hessian method

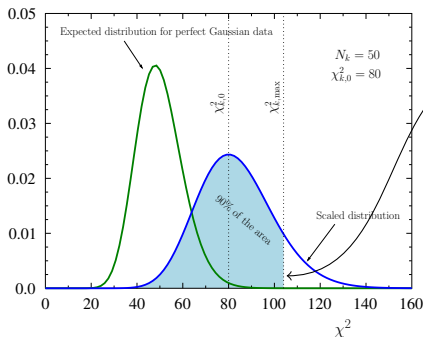
- The deviations $\delta z_i^\pm$ are determined by “hypothesis testing” finding “90% confidence limits” from scaled χ^2 distributions



1. Find $\chi^2_{k,max}$ for each data set k

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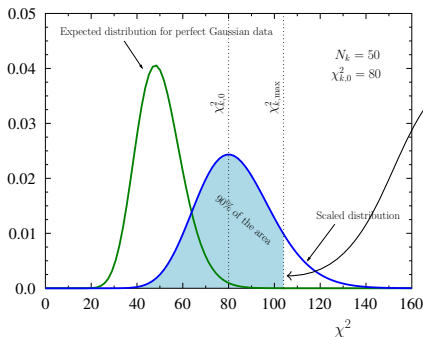


1. Find $\chi^2_{k,max}$ for each data set k

2. Find the corresponding variations $\delta z_{i,k}^\pm$ for all directions i

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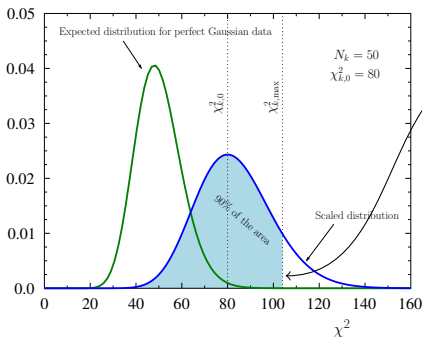
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3. Find the most limiting of all $\delta z_{i,k}^\pm \Rightarrow \delta z_i^\pm$

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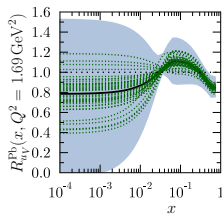
2. Find the corresponding variations $\delta z_{i,k}^\pm$ for all directions i

3. Find the most limiting of all $\delta z_{i,k}^\pm \Rightarrow \delta z_i^\pm$

- Uncertainty sets $S_i^\pm$ are **limiting cases** that still reproduce all the data within “90% confidence limit”
- Procedure is *ad hoc* but used by several PDF fitters

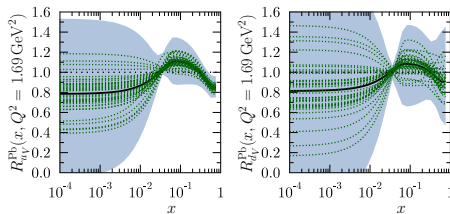
Results: Nuclear modification for ^{208}Pb at $Q^2 = m_{\text{charm}}^2$

- Total uncertainties shown as blue bands, individual errorsets in green



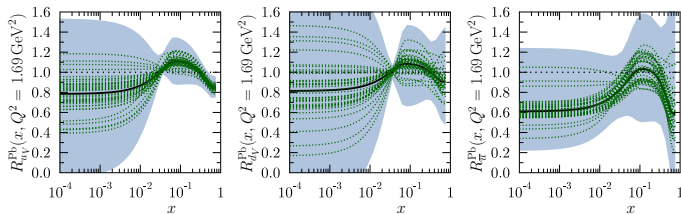
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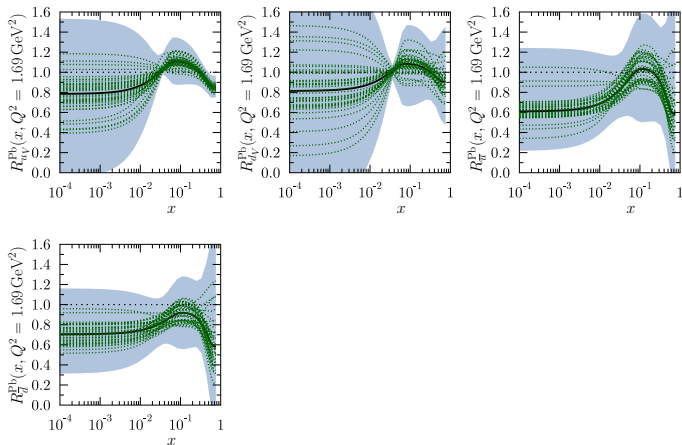
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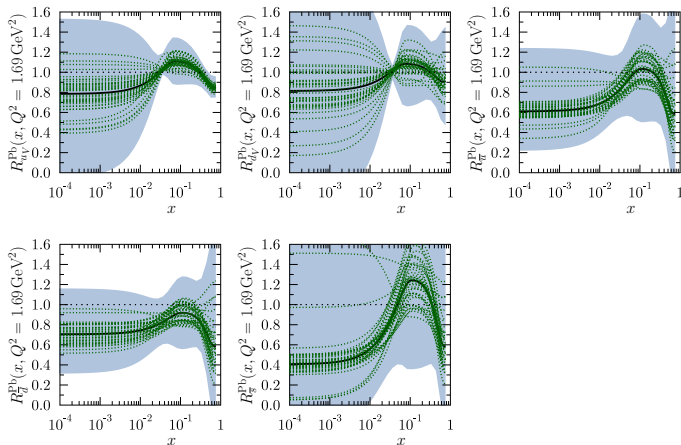
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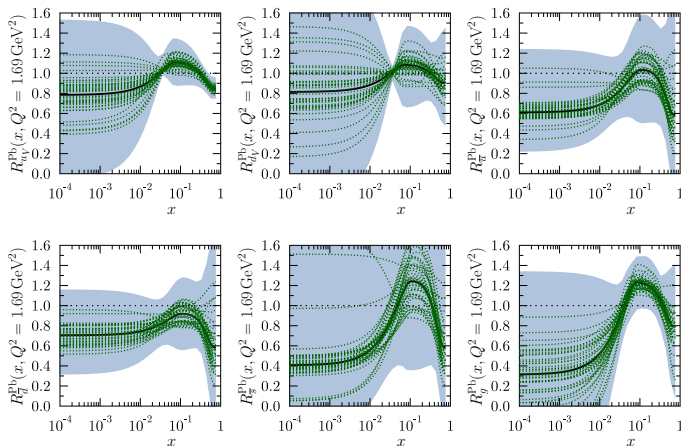
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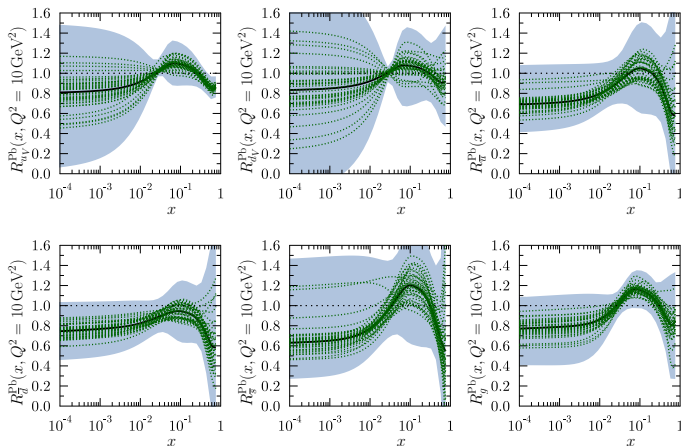
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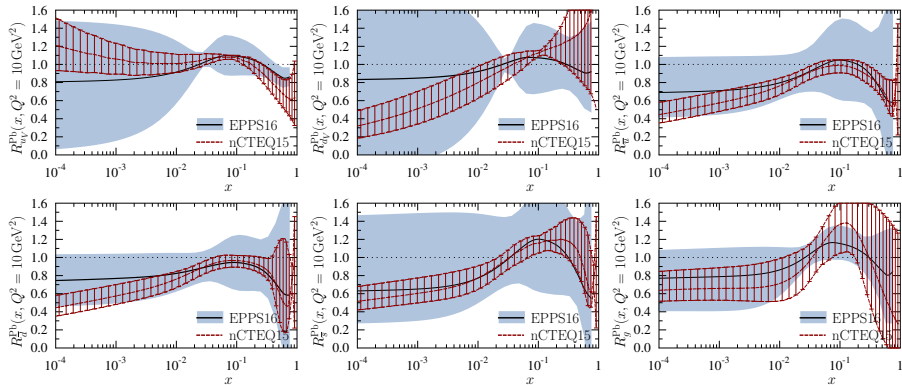
Results: Nuclear modification for ^{208}Pb at $Q^2 = 10 \text{ GeV}^2$

- Total uncertainties shown as blue bands, individual errorsets in green



Results: Comparison with nCTEQ15 nuclear PDFs

- EPPS16 consistent with nCTEQ15
- Typically smaller uncertainties in nCTEQ15 $\Leftarrow$ less freedom in the parametrization
- Larger high- x gluon uncertainties in nCTEQ15 $\Leftarrow$ looser cuts and no LHC data

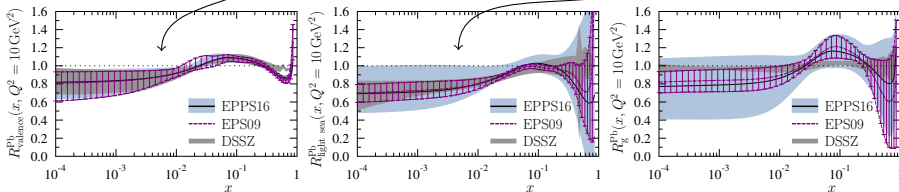


Results: Comparison with EPS09 and DSSZ

- No flavour freedom in EPS09 nor DSSZ — compare the averages

$$R_{\text{valence}} \equiv \frac{u_V^{p/Pb} + d_V^{p/Pb}}{u_V^p + d_V^p}$$

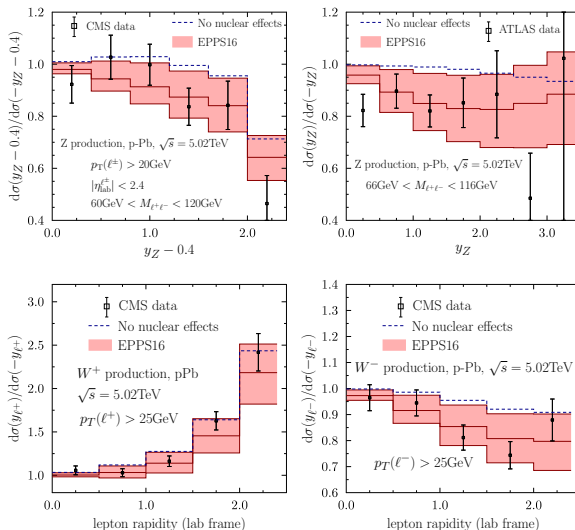
$$R_{\text{light sea}} \equiv \frac{\bar{u}^{p/Pb} + \bar{d}^{p/Pb} + \bar{s}^{p/Pb}}{\bar{u}^p + \bar{d}^p + \bar{s}^p}$$



- EPPS16 consistent with these earlier analyses
- Typically larger uncertainties in EPPS16 (for having more degrees of freedom)

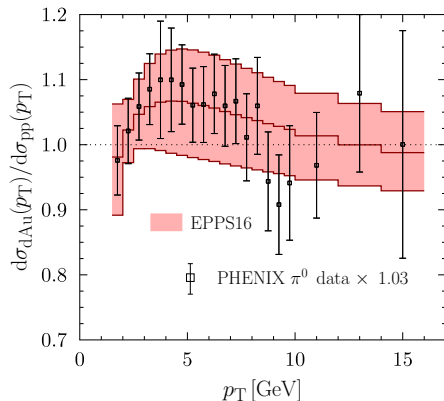
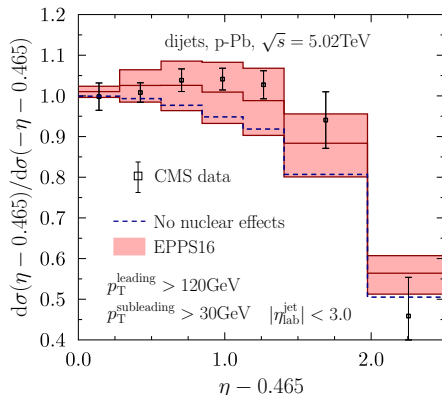
Comparison with the LHC p-Pb electroweak data

- Good agreement — lowish statistics limits the obtainable constraints



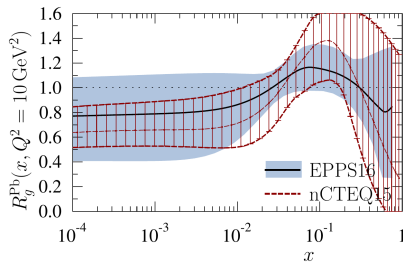
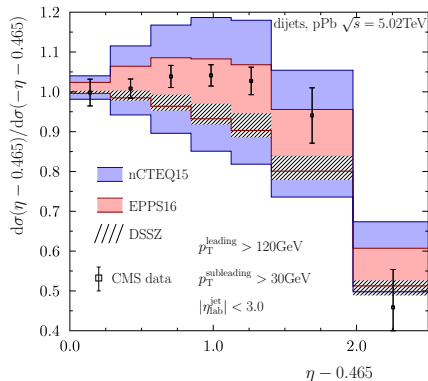
Comparison with the LHC p-Pb dijet and RHIC d-Au pion data

- Both data sensitive to the large- x gluons — mutually compatible



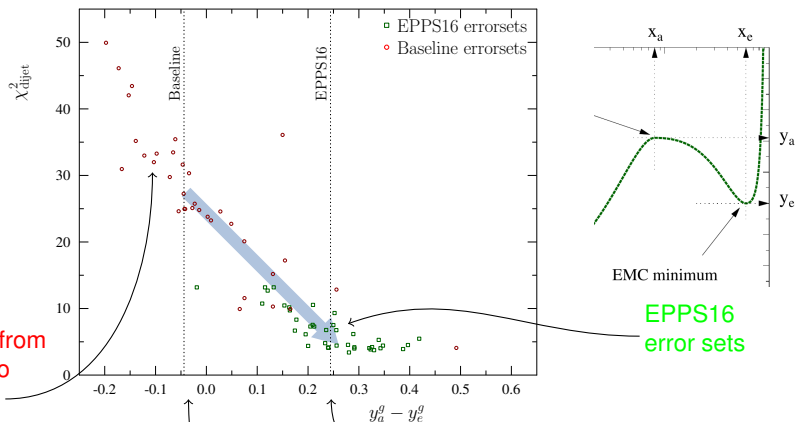
Comparison with the LHC p-Pb dijet and RHIC d-Au pion data

- Larger gluon uncertainties of nCTEQ15 clearly reflected in the dijets



Comparison with the LHC p-Pb dijet data

- For us the dijet data are essential in establishing a gluon EMC effect

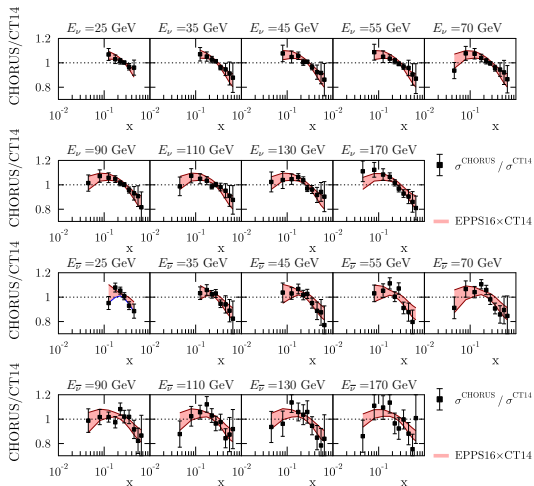


No gluon EMC effect here

Gluon EMC effect

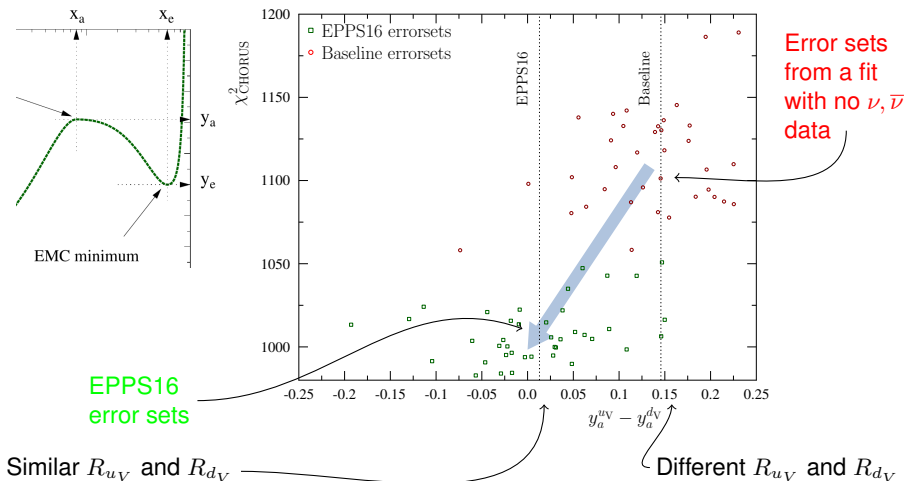
Comparison with the neutrino data

- The neutrino DIS data from CHORUS collaboration show a typical pattern of antishadowing + EMC effect (unity without nuclear mods)



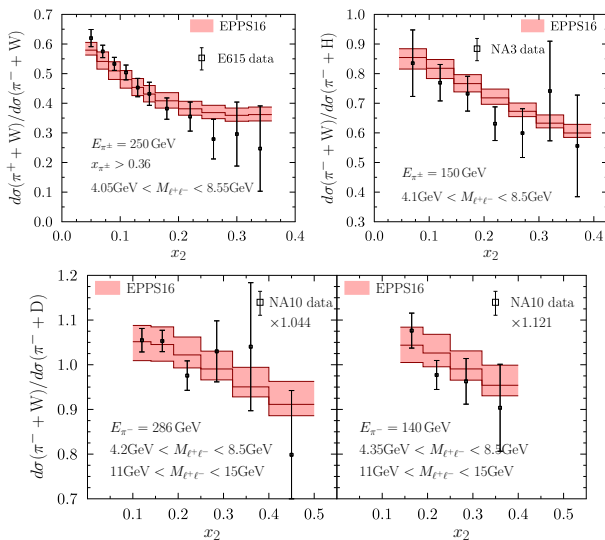
Comparison with the neutrino data

- The neutrino data are essential in establishing mutually similar R_{uV} and R_{dV}



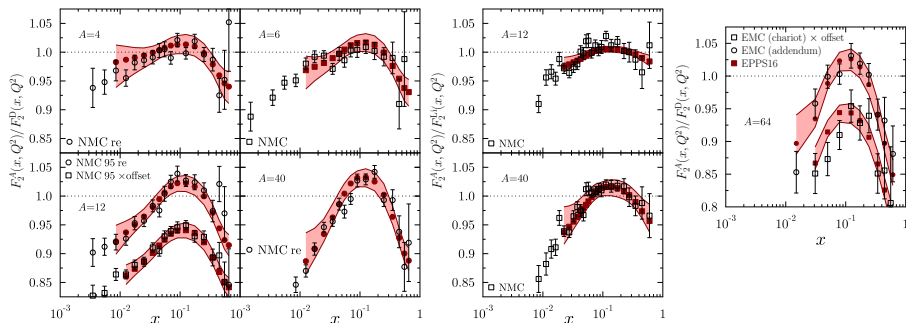
Comparison with fixed-target $\pi^\pm$ -A Drell-Yan data

- Also the $\pi^\pm$ -A Drell-Yan data well consistent with EPPS16



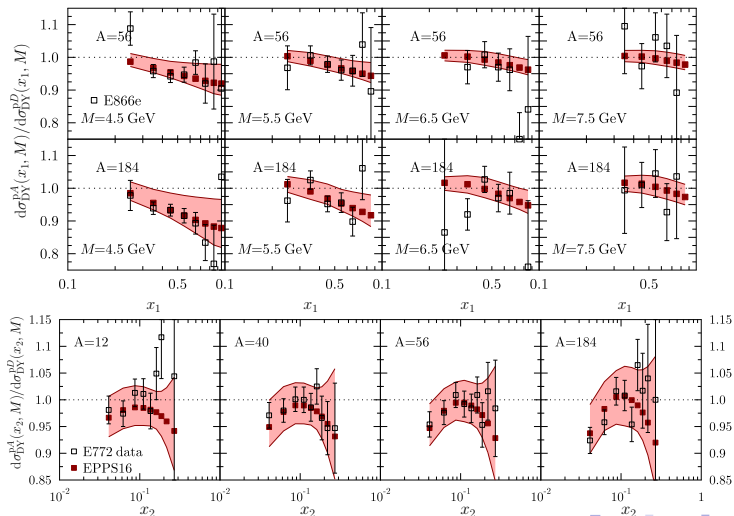
Comparison with charged-lepton DIS data

- No significant differences in the fit quality for older DIS measurements (included already in EPS09)



Comparison with Drell-Yan data

- No large differences in the fit quality for older Drell-Yan measurements neither (included already in EPS09)



Summary & Outlook

- Presented EPPS16 NLO nuclear PDFs
 - supersedes EPS09
 - More variety of data input than in other available fits
 - the most essential additions are the CMS dijet and CHORUS neutrino DIS data
 - More fit-function freedom flavour-by-flavour than in other available fits
 - less “biased” result
 - reveals large uncertainties flavour-wise
 - Excellent agreement with the data
 - supports the validity of collinear factorization in a largest region of (x, Q^2) plane examined to date
-
- More (much!) data expected in the near future
 - e.g. the second p-Pb run at $\sqrt{s} = 8 \text{ TeV}$
 - Ample room for theoretical/procedural improvements
 - NNLO, electroweak corrections (photon PDF), more flexible fit functions, including free proton data (no need for a separate baseline PDF), etc...