Heavy Ion Physics

Jan Rak and Sami Räsänen Jyväskylä University & Helsinki Institute of Physics, Finland













FYSH550 Relativistic Heavy Ion Physics

Lectures: Jan Rak

Sami Räsänen

Exercises: Sami Räsänen

Grading

One exam

provides "base grade" 80% questions + 20% exercises

Seminar in groups of two
-1, 0 or +1 to the base grade

Lectures

- Start Thursday 15th January 2009
- Continue:
 - Mondays 8-12 or 10-14?
 - Tuesdays 8-12?
- Total of 12 weeks, i.e. 48 hours
- Exceptional dates:
 - 2.2. 9.2. (4th international workshop High-pT physics at LHC 09)
 - 28.3. 6.4. (Quark Matter 2009)

Lecture topics

Lecture material: <u>https://trac.cc.jyu.fi/projects/alice/wiki/jan</u>

- Week I: Introduction part 1
- Week II: Introduction part 2
- Week III: Collective phenomena
- Week IV: Photons and π^0
- Week V: Correlations
- Week VI: Jet quenching
- Weeks VII-X: Seminar working on chosen topic
- Weeks XI and XII: Presentations

Topics for seminar works are introduced during weeks IV-VI. We all gather together in weeks VII-X to see progress and to discuss open issues.

ALICE @ JYFL

- Jan Rak
- DongJo Kim
- Sami Räsänen
- Norbert Novitzky
- Jiri Kral







Opportunities

Visit ALICE experiment at CERN

- dead line for applications 31st January
- send free form applications to Sami via email
- see details from

http://www.jyu.fi/static/fysiikka/sekalaista/StudentApplication.pdf

CERN summer student positions (HIP)

- dead line for applications 31st January
- total of 17 positions, all available for JYFL students
- ALICE groups position is 14. related with EMcal
- further info, instructions and application form

http://www.hip.fi/education/kesaharjoittelu.html

JYFL summer student position

- stay tuned, places should open ~February

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What is the world made of? What holds the world together? Where did we come from?



Courtesy: Y.K. Kim

- 1. Are there undiscovered principles of nature: New symmetries, new physical laws?
- 2. How can we solve the mystery of dark energy?
- 3. Are there extra dimensions of space?
- 4. Do all the forces become one?
- 5. Why are there so many kinds of particles?
- 6. What is dark matter?
 - How can we make it in the laboratory
- 7. What are neutrinos telling us?
- 8. How did the universe come to be?
- 9. What happened to the antimatter?

From "Quantum Universe"

Evolved Thinker

Courtesy: Y.K. Kim

High Energy (ultra-relativistic) Heavy Ion physics



Heavy lon physics:

- nonperturbative QCD phenomena at high-temperature and high-energy density.
- addresses the properties of QCD vacuum relevant to the structure of the early universe
- the origin of particle masses, QCD phase transition
- thanks to the Maldacena discovery of the duality between super gravity in Anti de Sitter space and conformal field theory, heavy ion physics provides a unique testbed for applied string physics.

Some expectations for the LHC

• With coming era of the CERN Large Hadron Collider (LHC), the largest highenergy particle accelerator ever built, many major discoveries are expected.

• It is known that the Standard Model (SM) of particles and fundamental forces cannot be a complete theory. Although exceptionally successful in explaining a vast variety of physics phenomena, the SM leaves us with many open questions and problems. Examples include:

- the disparity between the characteristic strengths of the weak and strong nuclear forces (the "hierarchy problem"),
- numerical quantities which are not predicted by the theory such as the mass of each particle,
- structure of the QCD vacuum, no explanation for quark confinement etc.

• Many extensions to the SM such as string theory, extra dimensions, black holes creation and many others are considered to be within the experimental reach of the LHC.

Recall basic scales:

Structure within

the Atom

e⁻

Electron

Ouark

Size < 10^{-19} m

Nucleus

- Atom ~ 1 Å = 10⁻¹⁰ m
- Nucleus ~ 10 fm = 10⁻¹⁴ m
- Nucleon (= p tai n) ~ 1 fm
- elementary particles are point like (r < 10⁻¹⁸ m) in SM







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Lattice QCD calculations verify that deconfinement transition should be a real phase transition.

Heavy ion collisions provide an opportunity to study thermodynamics of strongly interacting matter in laboratory



Lattice QCD, Lect. Notes Phys 583, 209 (2002)

Calculations and simulation by Harri Niemi

Some more details when we consider collective phenomena

Shuryak 1980

Jan F

 Shuryak publishes first "review" of thermal QCDand coins a phrase:

> "Because of the apparent analogy with similar phenomena in atomic physics, we may call this phase of matter the QCD (or quark-gluon) plasma."

(QGP)

PHYSICS REPORTS



QUANTUM CHROMODYNAMICS AND THE THEORY OF SUPERDENSE MATTER

Edward V. Shuryak

Volume 61 Number 2

May 1980

PRPLCM 61(2) 71-158 (1980)



NORTH-HOLLAND PUBLISHING COMPANY Amsterdam ed quantum r fields, the nology and earts by the nics (QED). elying upon al hadronic hs, etc.), but talogy with tark-gluon) he methods

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discussions Linde, A.B. harov and



Heavy Ion Physics - new paradigm Anti de-Sitter/CFT duality conjecture



quantum gravity that is

"blown up and slowed down by a factor of 10²⁰".

Relativistic kinematics

$$p_{\parallel}^{CMS} = \gamma^{CMS} \left(p_{\parallel}^{LAB} - \beta^{CMS} E^{LAB} \right) \qquad p_{\parallel}^{LAB} = \gamma^{CMS} \left(p_{\parallel}^{CMS} + \beta^{CMS} E^{CMS} \right)$$
$$E^{CMS} = \gamma^{CMS} \left(E^{LAB} - \beta^{CMS} p_{\parallel}^{LAB} \right) \qquad E^{LAB} = \gamma^{CMS} \left(E^{CMS} + \beta^{CMS} p_{\parallel}^{CMS} \right)$$
Energy/momentum transformation between any inertial frames e.g. CMS and LAB

Useful to work with four-vectors -> vectors in Minkowski 4D space and metric tensor of Mink. space

$$x^{\mu} = (x^{0}, x^{1}, x^{2}, x^{3}) = (t, \vec{x}) = (t, x, y, z)$$
$$p^{\mu} = (E, \vec{p}) = (E, \vec{p}_{T}, p_{z}) = (E, p_{x}, p_{y}, p_{z})$$
$$p^{\mu}p_{\mu} = E^{2} - p^{2} = m_{0}^{2}$$
Famous Invariant

$$g_{\mu\eta} = g^{\mu\eta} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$
$$x_{\mu} = g_{\mu\eta} x^{\eta} = (t, -x, -y, -z)$$
$$a \cdot b = a^{\mu} b_{\mu} = a_{\mu} b^{\mu} = g_{\mu\eta} a^{\mu} b^{\eta} = a^{0} b^{0} - \vec{a} \cdot \vec{b}$$

Mandelstam variables

$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2 \qquad s = 4p^2 = 4E_{beam}^2 = E_{cm}^2$$

$$t = (p_1 - p_3)^2 = (p_2 - p_4)^2 \qquad t = -2p^2(1 - \cos\theta)$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2 \qquad u = -2p^2(1 + \cos\theta)$$

 $\sqrt{s} = E_{cm}$ = center of mass energy available for particles production!

$$s + t + u = 3p_1^2 + p_2^2 + p_3^2 + p_4^2 + 2p_1 \cdot p_2 - 2p_1 \cdot p_3 - 2p_1 \cdot p_4$$
$$= \sum p_i^2 + 2p_1 \cdot (p_1 + p_2 - p_3 - p_4) = \sum m_i^2$$



Kinematics of relativistic particle interaction



$$s = (p_1 + p_2)^2 = (p_3 + p_4)^2$$

$$t = (p_1 - p_3)^2 = (p_2 - p_4)^2$$

$$u = (p_1 - p_4)^2 = (p_2 - p_3)^2$$

 $s + t + u = \sum_{i=1}^{4} p_i^2 = \sum_{i=1}^{4} m_i^2$

s-channel

 \mathbf{p}_2

momentum conservation: p₁+p₂=p₃+p₄

p_i = four momenta of in/out going particles
s,t,u
Lorentz invariant quantities

For a given s, both t and u depend linearly on the cos of the CMS deflection angle

$$-t = -(p_1 - p_3)^2 = 2E_1^* E_3^* - m_1^2 - m_3^2 - 2p_1^* p_3^* \cos \Theta^*$$

$$-u = -(p_1 - p_4)^2 = 2E_2^* E_3^* - m_2^2 - m_3^2 - 2p_2^* p_3^* \cos \Theta^*$$

In the case of elastic scattering and for fixed s $t = -2p^2(1 - \cos\Theta^*)$



p₄



u-channel

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Annihilation vs scattering

In the partonis CMS

$$\frac{e^{+}e^{+} \rightarrow \mu^{+}\mu^{+}}{\frac{d\sigma}{d\Omega}} = \frac{\alpha^{2}}{8p^{2}} \left(\frac{t^{2}+u^{2}}{s^{2}}\right) = \frac{\alpha^{2}}{8p^{2}} \left[\sin^{4}(\vartheta/2) + \cos^{4}(\vartheta/2)\right]$$

$$= \frac{\alpha^{2}}{4s} \left[1 + \cos^{2}\vartheta\right]$$

$$2\sin^{2}(\vartheta/2) = 1 - \cos\vartheta$$
$$2\cos^{2}(\vartheta/2) = 1 + \cos\vartheta$$
$$s = 4p^{2}$$

$$e^{-}\mu^{+} \rightarrow e^{-}\mu^{+}$$

Crossed channels t \Leftrightarrow s

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{8p^2} \left(\frac{s^2 + u^2}{t^2} \right) = \frac{\alpha^2}{8p^2 \sin^4(\vartheta/2)} \left[1 + \cos^4(\vartheta/2) \right]$$

Fixed target vs collider



 $\mathsf{CMS}_{\mathsf{NN}} \neq \mathsf{LAB}$

much of the energy goes towards forward motion of the particles that result from the impact with the target

$$\sqrt{s} \approx \sqrt{m_0 2 E_{beam}}$$



CMS _{NN}	Ξ	LAB
CMS _{ee}	=	LAB
CMS _{partonic}	¥	LAB

the total energy of the two beams is available for producing new particles

$$\sqrt{s} = 2E_{beam}$$

Often used tricks

Fix target CMS != LAB. Fraction of the beam energy is converted (wasted) into a kinetic energy of the center of mass system in LAB frame.

$$\begin{split} P^{\mu}\Big|_{\text{beam}} &= (E, \vec{0}_{T}, p_{\parallel}) \qquad P^{\mu}\Big|_{\text{target}} = (m, \vec{0}_{T}, 0_{\parallel}) \\ P^{\text{LAB,tot}}_{\parallel} &= P^{\text{beam}}_{\parallel} + P^{\text{target}}_{\parallel} = \left(E^{\text{beam}}_{\parallel} + m_{0}, \vec{0}_{T}, p^{\text{beam}}_{\parallel}\right) \\ s &= P^{\text{LAB,tot}, \mu}_{\parallel} P^{\text{LAB,tot}}_{\parallel, \mu} = 2m_{0}^{2} + 2m_{0}E^{\text{beam}}_{\parallel} \cong 2m_{0}E^{\text{beam}}_{\parallel} \end{split}$$



Rapidity, pseudorapidity

Rapidity y is a Lorentz equivalent of a non-relativistic velocity

$$y = \frac{1}{2} \ln \left(\frac{E + p_{\parallel}}{E - p_{\parallel}} \right) = \frac{1}{2} \ln \left(\frac{1 + \beta_{\parallel}}{1 - \beta_{\parallel}} \right)$$

y is Lorentz additive quantity!

Assume hadron *h* in CMS of y_h , CMS moves with y_{CMS} with respect to LAB. Then y_h in LAB = $y_h + y_{CMS}$ Pseudorapidity η approximation for y in the high-energy limit.

$$\eta = -\ln\left(\tan\frac{\theta}{2}\right) = \frac{1}{2}\ln\left(\frac{|p| + p_{\parallel}}{|p| - p_{\parallel}}\right)$$

If no PID we do not no the rest mass and thus no E, we can measure and angle θ wrt beam.

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Useful



where $m_T = \sqrt{m_0^2 + p^2}$ is so called transverse mass



HI - Center Of Mass Energy regimes



Brookhaven Nat. Lab. Long Island, USA 27



Alternating Gradient Synchrotron Complex



At Brookhaven National Lab. Long Island (NY area) USA

Fix target √s≈4 GeV 3 Nobel prizes



SPS - Super Proton Synchrotron



At CERN

Fix target √s≈17 GeV (E_{beam}=160 GeV)

- Circumference : 6.9 km
- 2.5 km of secondary beam lines.
- protons for fixed target physics at 400 GeV/c
- protons for CNGS experiment at 400GeV/c
- protons for LHC at 450GeV/c
- lead ions for fixed target physics at 400 GeV/c proton equivalent
- machine studies for SPS
- machine studies for LHC
- · Injector for the LHC

Relativistic Heavy Ion Collider (RHIC)



RHIC

Two independent rings 3.83 km in circumference

Maximum Energy per N-N collision

 $\sqrt{s} = 500 \text{ GeV p-p} \text{ (polarized)}$

$$\sqrt{s} = 200 \text{ GeV} \text{Au-Au}$$

Design Luminosity

Au-Au 2x10²⁶ cm⁻²s⁻¹ p - p 2x10³² cm⁻²s⁻¹





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Relativistic Heavy Ion Collider (RHIC)



PHENIX

The PHENIX Experiment, main emphasis on electromagnetic probes. Focus:

- •Rare probes J/ Ψ , Ψ ', e⁺e⁻, μ ⁺ μ ⁻, Φ , direct- γ ...
- The spin structure of the nucleons

The Configuration:

- •2 Forward Muon Arms
- •2 Central Spectrometer Arms to measure photons, electrons, and hadrons

PHENIX Central Arm

PHENIX Detector - Second Year Physics Run

- proton, anti-proton < 4 GeV/c
- $\cdot \quad \Delta \phi = \pi/4$

$\pi^{\rm 0}$ measurement by EMCal

- . 1<pt<15 GeV/c
- . 6 lead- scintillator (PbSc) sectors
- . 2 lead- glass (PbGl) sectors
- . $|\eta| < 0.38$ at midrapidity, $\Delta \phi = \pi$

 global observables, event-by-event physics, HBT, strangeness, high-pt jets...

The Configuration:

·large acceptance TPC, Silicon Vertex Tracker, RICH, TOF, EMC...

Au on Au central event at √s=130GeV \$**7**AR PH*****ENIX

beam view

The LHC Machine and Experiments

The LHC is Coming!

Energy in the LHC: What is "big" vs "small"

Collision energy pp -collisions: $E_{cms} = 14 \text{ TeV}$ PbPb -collisions: $E_{cms} = 5500 \text{ A GeV}$

In pp -collision, goal is total 2808 bunches of protons in LHC beams each containing 1.15 x 10¹¹ protons

- 1) Compare energy of a single proton or lead ion to a kinetic energy of a house fly with mass 12 mg
- 2) Estimate the time between bunch crossings
- 3) Compare total energy of proton beams to a kinetic energy of Pendolino train with mass 316 tons
- 4) Compare energy per mass ratio for a particle in the LHC and a car in high way

Physics at the LHC: pp @ 14 TeV

LHC will explore directly the highly-motivated TeV-scale and say the final word about the SM Higgs mechanism and many TeV-scale New Physics predictions
Also LHC will be a great machine for: QCD, B-physics, Heavy Ions, EW precision..

WHY HEAVY IONS AT THE LHC?

... factor ~30 jump in vs ...

J. Schukraft QM2001: hotter - bigger -longer lived

Central collisions	SPS	RHIC	LHC
s ^{1/2} (GeV)	17	200	5500
dN _{ch} /dy	500	850	2–8 x10 ³
ε (GeV/fm³)	2.5	4–5	15–40
V _f (fm³)	10 ³	7x10 ³	2x10 ⁴
τ _{QGP} (fm/c)	<1	1.5–4.0	4–10
τ ₀ (fm/c)	~1	~0.5	<0.2

$$\begin{split} \epsilon_{LHC} &> \epsilon_{RHIC} > \epsilon_{SPS} \\ V_{f LHC} &> V_{f RHIC} > V_{f SPS} \\ \tau_{LHC} &> \tau_{RHIC} > \tau_{SPS} \end{split}$$

CMS: The Detector

CMS: Capabilities & Performance

- Large acceptance tracking and calorimetry:
 - -2π in azimuth (jets,..)
 - Si-tracker InI<2.5 (b-,c- physics, ...)
 - uniquely large range in rapidity ($|\eta| < 6.6$, $|\eta_{neutral}| > 8$) (\rightarrow forward physics)

CMS has the coverage to address the physics

• High resolution

- Granularity of the Si-pixel layer+4T mag field \rightarrow

 $\Delta p_T/p_T < 1.5\%$ for $p_T < 50$ GeV/c

– For jet finding, calorimeters σ_{η} =0.028 σ_{ϕ} =0.032

CMS has the detector performance needed to resolve the physics

The ATLAS Detector

The ATLAS Detector

- Full azimuthal acceptance in all detectors
- Unprecedented pseudorapidity coverage for A+A

ATLAS Calorimetry

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