

SCATTERING THEORY

FALL

2007

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1. SPECTRAL THEORY FOR UNBOUNDED SELF-ADJOINT OPERATORS

Scattering theory is analysis of the Hamiltonian operators

$$H = \Delta + \mathbb{L}^2 + V$$

with certain boundary cond. for solutions at infinity.

Here $H_0 = \Delta + \mathbb{L}^2$ is the free Hamiltonian and V the scattering potential.

H_0 is a self-adjoint unbounded operator in L^2 . Thus we need to study

1.1. UNBOUNDED OPERATORS IN A HILBERT SPACE

Let $T: \mathcal{D}(T) \rightarrow H$
be a linear (not cont.)
operator, where $\mathcal{D}(T) \subset H$
is a linear sub-space. The
set $\mathcal{D}(T)$ is called the *domain*
of T .

We shall assume in the sequel
that $\mathcal{D}(T)$ is *dense* in H .

T is *closed* if

$$\left. \begin{array}{l} \mathcal{D}(T) \ni x_n \rightarrow x \\ Tx_n \rightarrow y \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} x \in \mathcal{D}(T) \\ Tx = y \end{array} \right. ,$$

This is equivalent to
 $(\mathcal{D}(T), \|\cdot\|_T)$ is Hilbert space

where $\|x\|_T = (\|x\|^2 + \|Tx\|^2)^{1/2}$.

Proof Ex.

Adjoint operator: Let

$$A: \underset{\substack{\uparrow \\ H}}{D(A)} \longrightarrow H$$

(densely defined). Then

$$D(A^*) = \{y \in H \mid \exists y^* \in H \text{ s.t.}$$

$$(Ax, y) = (x, y^*), \quad \forall x \in D(A)\}$$

and we set

$$A^*y = y^*, \quad \text{for } y \in D(A^*).$$

Remark A^* is well-defined since A is densely defined.

1.2 SYMMETRIC AND SELF-ADJOINT OPERATORS

Def (i) A is symmetric if

$$(Ax, y) = (x, Ay)$$

$$\forall x, y \in D(A)$$

(ii) A is self adjoint if $A = A^*$.

Remark 1 Self adjoint operators are symmetric.

Remark 2 A is symmetric iff $(Ax, x) \in \mathbb{R}, \forall x \in D(A)$.

1.1. Proposition Assume $A : D(A) \rightarrow H$ is symmetric. Then A is self-adjoint, iff

$$(1) \quad R(A - \lambda I) = H \\ \forall \lambda \in \mathbb{C} \setminus \mathbb{R}.$$

Proof Assume (1) holds for λ and $\bar{\lambda}$ with $\text{Im } \lambda \neq 0$.

Then if $(A^* - \lambda I)u = 0$

we have for $\forall v \in D(A)$:

$$0 = ((A^* - \lambda I)u, v) \\ = (u, (A - \bar{\lambda} I)v) = 0$$

$$\Rightarrow u = 0.$$

Thus

$$(2) \quad \text{Ker}(A^* - \lambda I) = 0$$

If now $u \in \mathcal{D}(A^*)$, we have by (1), that $\exists v \in \mathcal{D}(A)$ s.t.

$$(A^* - \lambda I)u = (A - \lambda I)v$$

Since $A \subset A^*$ we get

$$(A^* - \lambda I)(u - v) = 0$$

and by (2)

$$u = v \in \mathcal{D}(A)$$

Thus $\mathcal{D}(A^*) = \mathcal{D}(A)$

and A is self-adjoint.

The converse follows from

1.2 Prop

Let A be closed and symmetric. Then for $\lambda \in \mathbb{C} \setminus \mathbb{R}$, $A - \lambda I$ is injective with closed range and

$$(1) \quad |\text{Im } \lambda| \|u\| \leq \|(A - \lambda I)u\|$$

The codimension of $\text{Im } A - \lambda I$ depends only on $\text{Im } \lambda$.

Proof Since $A - \lambda I$ is symmetric we have for $\lambda = x + iy$

$$\begin{aligned} & \| (A - \lambda I) u \|^2 \\ &= \left((A - xI) u + iy u, (A - xI) u + iy u \right) \\ &= \| (A - xI) u \|^2 + y^2 \| u \|^2 \\ &\quad + \left((A - xI) u, iy u \right) + \left(iy u, (A - xI) u \right) \\ &\geq y^2 \| u \|^2 \end{aligned}$$

which proves (1).

Let's prove that the range is closed:

If $u_n \in \mathcal{D}(A)$ and $(A - \lambda) u_n \rightarrow f$.

By (1) applied to $u_n - u_m$ we get

$$\| u_n - u_m \| \leq \frac{1}{|\text{Im } \lambda|} \| (A - \lambda)(u_n - u_m) \|$$

we have $u_n \rightarrow u$. Since $A - \lambda$ is closed we obtain

$$u \in \mathcal{D}(A) \text{ and } (A - \lambda)u = f.$$

$\Rightarrow R(A - \lambda)$ is closed.

Moreover

$$u \in \text{Ker}(A - \lambda) \Leftrightarrow Au = \lambda u$$

$$\Rightarrow (u, Au) = \bar{\lambda} \|u\|^2$$

Since A is symmetric is l.h.s. is real and hence $u = 0$

This shows that $A - \lambda$ is injective for $\text{Im } \lambda \neq 0$.

We omit the proof of the last claim. It follows from the stability of the index of bounded Fredholm operators. \square

1.3 Definition The defect indices $n_{\pm} = \text{codim } R(A - \lambda)$, $\pm \text{Im } \lambda > 0$

Note that

$$n_{\pm} = \dim \text{Ker}(A^* - \bar{\lambda}), \quad \mp \text{Im } \lambda > 0$$

1.4 Theorem If A is self-adjoint then the resolvent $R(\lambda) = (A - \lambda)^{-1}$ is bounded operator for $\text{Im} \lambda \neq 0$ and

$$\|R(\lambda)\| \leq (\text{Im} \lambda)^{-1}$$

Moreover $R(\lambda)^* = R(\bar{\lambda})$ and $(\lambda - \mu)R(\lambda)R(\mu) = R(\lambda) - R(\mu)$

Proof Since range of $R(\mu)$ is in $D(A)$

$$((A - \mu) - (A - \lambda))R(\mu)$$

$$= (\lambda - \mu)R(\mu)$$

Multiplying by $R(\lambda)$ from the left we get

$$R(\lambda) - R(\mu) = (\lambda - \mu)R(\mu)$$

□

1.5 Theorem Assume A is closed and symmetric. Then

(1) A is self-adjoint $(\Leftrightarrow) n_{\pm}(A) = 0$

(2) h has an self adjoint extension (\Leftrightarrow)

$$n_{+}(A) = n_{-}(A)$$

Proof We prove (1) and sketch
the proof of (2):

Since $n_{\pm} = \dim \text{Ker}(A^* - \lambda)$,
 $\pm \text{Im} A < 0$

If $A^* = A$ we have $n_{\pm} = 0$
by Prop 1.2. The proof of the
converse claim follows exactly as
the proof of Prop. 1.1.

The idea behind the proof of (2)
is that

$T: (A + i)u \mapsto (A - i)u$
is isometric if A is symmetric.
Moreover

$$R(A - i) = R(T)$$

$$R(A + i) = \mathcal{D}(T)$$

If $n_+ = n_-$, T has a
unitary extension (\Leftrightarrow) A has
a self-adjoint extension.

□

Friedrich's extension:

A is semi-bounded if for some $c \in \mathbb{R}$

$$(1) \quad A \geq cI \quad (\Leftrightarrow)$$

$$(Au, u) \geq c \|u\|^2$$

$\forall u \in D(A)$. Note that (1) (\Leftrightarrow)

$$c := \inf_{\substack{u \in D(A) \\ \|u\|=1}} (Au, u) > -\infty$$

1.6 Theorem If A is symmetric and semi-bounded, then \exists self-adjoint extension A_D with $(A_D u, u) \geq c \|u\|^2$, $u \in D(A_D)$

Outline of proof:

By replacing A with

$$A + (1-c)I$$

we may assume $c=1$. Then

(1) $\|u\| \leq \|u\|_D$ where the new norm is defined by inner product

$$(2) \quad (u, v)_D = (Au, v)$$

Define D be the completion of D_A w.r.t. (ρ) . By (1)

$$D \subset H.$$

Finally define A_D is the restriction of A to $D_A^* \cap D$.

□

1.3 PROJECTION VALUED MEASURES AND SPECTRAL THEOREMS

A projection valued measure is mapping from Borel-sets of the real line \mathbb{R} to projections of H provided that the following holds:

$$E : B(\mathbb{R}) \longrightarrow L(H);$$

$$\text{denote } E(\Omega) = E_\Omega.$$

a) E_Ω is orthogonal projection
for all $\Omega \in \mathcal{B}(\mathbb{R})$

b) $E_\emptyset = 0$; $E_{\mathbb{R}} = I$

c) If $\Omega = \bigcup_{n=1}^{\infty} \Omega_n$, $\Omega_n \cap \Omega_m = \emptyset$
for $n \neq m$, then
$$E_\Omega = s\text{-lim}_{N \rightarrow \infty} \sum_{n=1}^N E_{\Omega_n}$$

d) $E_{\Omega_1} E_{\Omega_2} = E_{\Omega_1 \cap \Omega_2}$

Remarks (i) c) means that
$$\| E_\Omega x - \sum_{n=1}^N E_{\Omega_n} x \| \rightarrow 0$$

as $N \rightarrow \infty$.

(ii) It is easy to see that
for $x \in H$

$$(x, E_\Omega x)$$

is a well-defined Borell
measure that we denote by

$$d(x, E_\lambda x)$$

(similarly one defines the complex Borell - measure $d(y, E_\lambda x)$).

(iii) If we define an operator

$$A: H \rightarrow H, \text{ by}$$

$$(y, Ax) = \int_{-\infty}^{\infty} \lambda d(y, E_\lambda x) \quad (*)$$

and

$$D(A) = \{x \mid \int_{-\infty}^{\infty} |\lambda|^2 d(x, E_\lambda x) < \infty\}$$

then A is self-adjoint.

(iv) If there exists a number $\epsilon > 0$ such that

$$E_{(-\epsilon, \epsilon)} = \bar{I}$$

then A is bounded.

(v) If (*) holds for every $x, y \in H$ we write symbolically

$$A = \int \lambda dE_\lambda$$

The converse is also true

Spectral theorem For any self-adjoint operator $A \in L(H)$ there exists a p.v.m E_λ s.t.

$$A = \int_{-\infty}^{\infty} \lambda dE_\lambda \quad \text{and}$$

$$D(A) = \left\{ x \mid \int_{-\infty}^{\infty} |\lambda|^2 d(x, E_\lambda x) < \infty \right\}$$

Remark (i) This is one of the most powerful tools in analysis

(ii) If g is any Borel-function on \mathbb{R} , we may define the operator $g(A)$ by

$$g(A) = \int g(\lambda) dE_\lambda$$

and

$$D_{g(A)} = \left\{ x \mid \int |g(\lambda)|^2 d(x, E_\lambda x) < \infty \right\}$$

Note that $D_{g(A)}$ is dense in H
and that $g(A)$ is self-adjoint
if g is **real-valued**.

(iii) If g is bounded then

$$\begin{aligned} & \int_{-\infty}^{\infty} |g(\lambda)|^2 (x, dE_{\lambda} x) \\ & \leq M \int_{-\infty}^{\infty} (x, dE_{\lambda} x) \\ & = M (x, \underbrace{\int_{-\infty}^{\infty} dE_{\lambda}}_{=I} x) \\ & = M \|x\|^2 < \infty \end{aligned}$$

Thus $D_{g(A)} = H$

for all self-adjoint A
and consequently $g(A)$ is
in $L(H)$.

1.4 Examples of self-adjoint operators

1.4.1 Example Let M be locally compact topological space and $d\mu$ a positive (Borel) measure on M . Take

$$H = L^2(d\mu).$$

Assume $\beta: M \rightarrow \mathbb{R}$ is measurable and define

$$A: H \rightarrow H; \quad Af = \beta f,$$

$$D(A) = \{f \in L^2 \mid \beta f \in L^2\}$$

Then

(i) A is self-adjoint

(ii) E_λ is multiplication with the char. function of the set

$$\{x \mid \beta(x) < \lambda\}$$

Proof (i) If $f, g \in D(A)$, then $\int \overline{Ag} = \int \overline{\beta g} \in L^1$ and (f, Ag) is well defined.

Moreover

$$\begin{aligned}(f, Ag) &= \int_M f \overline{\beta g} \, d\nu = \\ &= \int_M \beta f g = (Af, g)\end{aligned}$$

and thus A is symmetric. Finally,

$$D(A^*) = \{g \in L^2 \mid \exists g^* \in L^2 \text{ s.t.}$$

$$(Af, g) = (f, g^*) ; \forall f \in D(A)\}$$

$$= \{g \in L^2 \mid \exists g^* \in L^2 \text{ s.t.}$$

$$\int \beta f \bar{g} \, d\nu = \int f \bar{g}^* \, d\nu\}$$

$$\stackrel{(i)}{=} \{g \in L^2 \mid \beta g = g^* \in L^2\}$$

$$= D(A)$$

Note that (i) follows from the density of $D(A)$.

(ii) We need to show

$$(f, Af) = \int_{\mathbb{R}} \lambda (f, d\chi_{\{\beta(x) < \lambda\}} f)$$

Here

$$\text{r.h.s.} = \int_{\mathbb{R}} \lambda (f, d\chi_{\{\beta(x) < \lambda\}} f)_M$$

$$=: \int_{\mathbb{R}} \lambda d\nu(\lambda)$$

Here the function $\nu(\lambda)$ is given by

$$\begin{aligned} \nu(\lambda) &= (f, \chi_{\{\beta(x) < \lambda\}} f)_M \\ &= \int_M |f(x)|^2 \chi_{\{\beta(x) < \lambda\}} d\nu(x) \end{aligned}$$

Since for all $s \in \mathbb{R}$

$$\int \lambda d\chi_{\{s < \lambda\}} = s$$

we have

$$\begin{aligned} \text{r.h.s.} &= \int_M \beta(x) |f(x)|^2 d\nu(x) \\ &= (1, Af) \end{aligned}$$

□

1.4.2 Ex Same as above but $M = \mathbb{R}$ and $\beta(\xi) = \xi$.

Moreover if λ is an **eigen value** of A , then **for almost all** $\xi \in \mathbb{R}$

$$A f(\xi) = \underbrace{\xi f(\xi)} = \lambda f(\xi)$$

If $\mu(\{\xi_0\}) > 0$, this is possible, by defining

$$f(\xi) = \begin{cases} 1, & \xi = \xi_0 \\ 0, & \xi \neq \xi_0 \end{cases}$$

Thus eigenvalues correspond to such points where $d\mu$ has positive mass.

1.5 Spectrum

If $A : H \rightarrow H$ is dense by definition we denote by $\rho(A)$ the **resolvent set**

$$\rho(A) = \{ \lambda \in \mathbb{C} \mid (A - \lambda)^{-1} \in L(H) \}$$

The complement of $\rho(A)$ in \mathbb{C} is called the **spectrum** of A and is denoted by $\sigma(A)$.

Recall λ is **eigenvalue** of A
if $Ax = \lambda x$, for $x \neq 0$.

The set of eigenvalues is called
point spectrum and is denoted
by $\sigma_p(A)$.

The **residual spectrum** is
 $\sigma_r(A) = \{\lambda \mid R(A - \lambda) \subset H$
is not dense and $\lambda \notin \sigma_p(A)\}$

Remark If A is self adjoint, then

(i) $\sigma(A) \subset \mathbb{R}$ and

(ii) $\sigma_r(A) = \emptyset$.

Proof (i) follows from Theorem 1.4

(ii) If $u \in H$ and $\lambda \in \mathbb{R}$ is
such that

$$(A - \lambda)f, u) = 0, \quad \forall f \in D(A)$$

Then by definition

$$u \in \mathcal{D}(A^*) = \mathcal{D}(A)$$

and $(f, (A - \lambda)u) = 0 \quad \forall f \in \mathcal{D}(A)$

Since $\mathcal{D}(A)$ is dense

$$Au = \lambda u$$

and $u \in \sigma_p(A)$

□