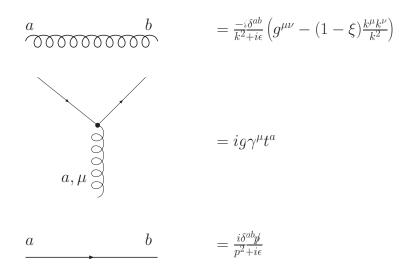
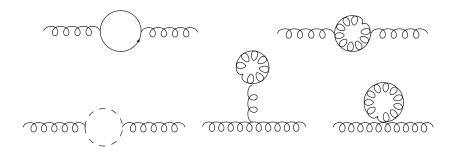
**1.** Compute the QCD correction to the quark self-energy for massless quarks at 1-loop order. The relevant Feynman rules are



where  $t^a$ 's are the QCD-generators. Do the calculation in the dimensional regularization scheme, using this time the Landau gauge  $\xi = 0$ .

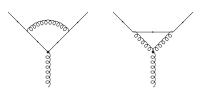
2. Compute the gluon self-energy to 1-loop order, using the dimensional regularization and the Feynman gauge. For the Feynman rules and algebra for SU(3), see the Appendix of Peskin & Schroeder. If you are unfamiliar with the color algebra, you may want to also consult the section 15.4 from P&S. Note that there are now several diagrams at 1-loop:



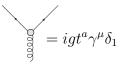
After the smoke has cleared, you should find the current-conserving structure  $(q^2 g^{\mu\nu} - q^{\mu}q^{\nu})$ . The Feynman rule for the counter term is given by

$$\bigcirc \bigcirc \bigcirc \bigcirc \bigcirc \bigcirc = -i \left( q^2 g^{\mu\nu} - q^\mu q^\nu \right) \delta^{ab} \delta_3$$

Perform the renormalization in the  $\overline{MS}$ -scheme. In this particular scheme one subtracts only the pole  $\sim 2/\epsilon - \gamma_E + \log(4\pi)$ . The finite part should now explicitly depend on  $\mu^2$ , the mass dimensional parameter introduced by the dimensional regularization. **3.** Compute the UV divergent parts of the 1-loop corrections to the QCD quark-quark-gluon 1PI-function. Take the quarks to be massless. There is now two diagrams to consider:



The Feynman rule for the counter term is given by



Perform the renormalization in the  $\overline{MS}$ -scheme by subtracting the pole  $\sim 2/\epsilon - \gamma_E + \log(4\pi)$ . Note that in order to extract the counter term  $\delta_1$  you need only the UV-divergent terms. There is also IR-divergences present in the diagrams above which can as well show up as a  $1/\epsilon$ -type poles. So be careful not to mix up these terms with UV-poles!

4. Show that the BRST transformation.

$$\begin{array}{rcl} QA^a_\mu &=& D^{ab}c^b\\ Q\psi &=& igc^at^a\psi\\ Qc^a &=& -\frac{1}{2}gf^{abc}c^bc^c\\ Q\bar{c}^a &=& B^a\\ QB^a &=& 0\,, \end{array}$$

where  $D^{ab}_{\mu} = \partial_{\mu} \delta^{ab} + g f^{abc} A^{c}_{\mu}$ , is nilpotent, *ie.*  $Q^{2} \phi = 0$  for all fields  $\phi = A^{a}_{\mu}, \psi, c^{a}, \bar{c}^{a}, B^{a}$ . Show also in detail that the transformation  $\phi \to \phi + \epsilon Q \phi$  (where  $\epsilon$  is an anticommuting number), is a continuous symmetry of the QCD Lagrangean

$$\mathcal{L}_{QCD} = \bar{\psi}(iD - m)\psi - \frac{1}{4}F^{a}_{\mu\nu}F^{a,\mu\nu} + \frac{\xi}{2}(B^{a})^{2} + B^{a}(\partial_{\mu}A^{a}_{\mu}) + \bar{c}^{a}(-\partial^{\mu}D^{ac}_{\mu})c^{c}.$$