

1. The contribution from the source term to the solution of the scalar field equation depends on the choice of the propagator function:

$$\delta\phi_X(x) \equiv \int d^4y \Delta_X(x-y) J(y),$$

where  $\Delta_X = \Delta_R, \Delta_A, \Delta_F$  tai  $\Delta_{\bar{F}}$ . Show that the *differences*  $\delta\phi_X(x) - \delta\phi_{X'}(x)$  fulfill the homogenous equation, and can thus be absorbed to the free part  $\phi_0(x) \rightarrow \phi_{0X}(x)$ . From this conclude that the use of different propagators simply corresponds to different boundary conditions at spatial and temporal infinity (we assume that  $j$  vanishes at infinity).

2. Verify that the integral representation of the Feynman propagator agrees with the expectation value of the time-ordered 2-point function:

$$D_F(x-y) = \int \frac{d^4p}{(2\pi)^4} e^{-ip \cdot (x-y)} \frac{i}{p^2 - m^2 + i\epsilon} = \langle |T(\phi(x)\phi(y))| \rangle.$$

3. Show that the generator matrices  $S^{\mu\nu} = \frac{1}{2}\sigma^{\mu\nu} \equiv \frac{i}{4}[\gamma^\mu, \gamma^\nu]$  fulfill the Lorentz algebra

$$[S^{\mu\nu}, S^{\rho\sigma}] = i(g^{\nu\rho} S^{\mu\sigma} - g^{\mu\rho} S^{\nu\sigma} - g^{\nu\sigma} S^{\mu\rho} + g^{\mu\sigma} S^{\nu\rho}).$$

when  $\gamma^\mu$ -matrices obey the Clifford algebra  $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$ .

4. Show that  $\gamma^\mu$  transforms as a vector under Lorentz-transform:

$$\Lambda_{\frac{1}{2}}^{-1} \gamma^\mu \Lambda_{\frac{1}{2}} = \Lambda^\mu{}_\nu \gamma^\nu.$$

Hint: Follow Peskin-Schröder on page 42.

5. We can write a unitary representation for an arbitrary 4-vector  $V$  as follows

$$V = V^0 + \vec{\sigma} \cdot \vec{V},$$

where  $\sigma_i$  are the familiar Pauli matrices. Let us now require that  $V$ -matrix transforms according to the  $(\frac{1}{2}, \frac{1}{2})$ -representation of the Lorentz-group,

$$V \rightarrow \Lambda_{\frac{1}{2}L}^{-1}(\theta, \beta) V \Lambda_{\frac{1}{2}R}(\theta, \beta)$$

where  $\Lambda_{\frac{1}{2}L,R}$  are  $2 \times 2$ -matrices that transform the Weyl spinors under  $(\frac{1}{2}, 0)$  and  $(0, \frac{1}{2})$  representations (see the lectures on page 42). Show that this transformation law implies that the components of  $V$  transform as a Lorentz-vector. (This demonstrates that  $(\frac{1}{2}, \frac{1}{2})$  in fact is a *vector* representation of the Lorentz-group.)