

# Gravitational waves from phase transitions in the early Universe

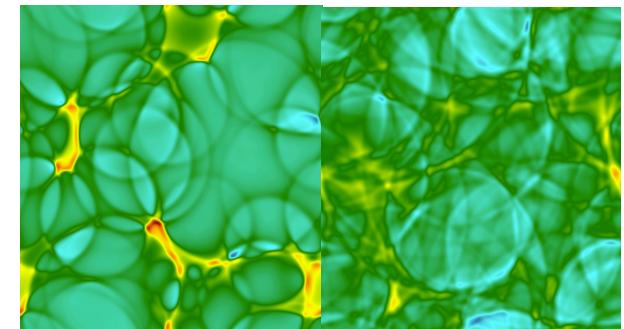
**Mark Hindmarsh**

Helsinki Institute of Physics & Dept of Physics, University of Helsinki  
and

Department of Physics & Astronomy, University of Sussex

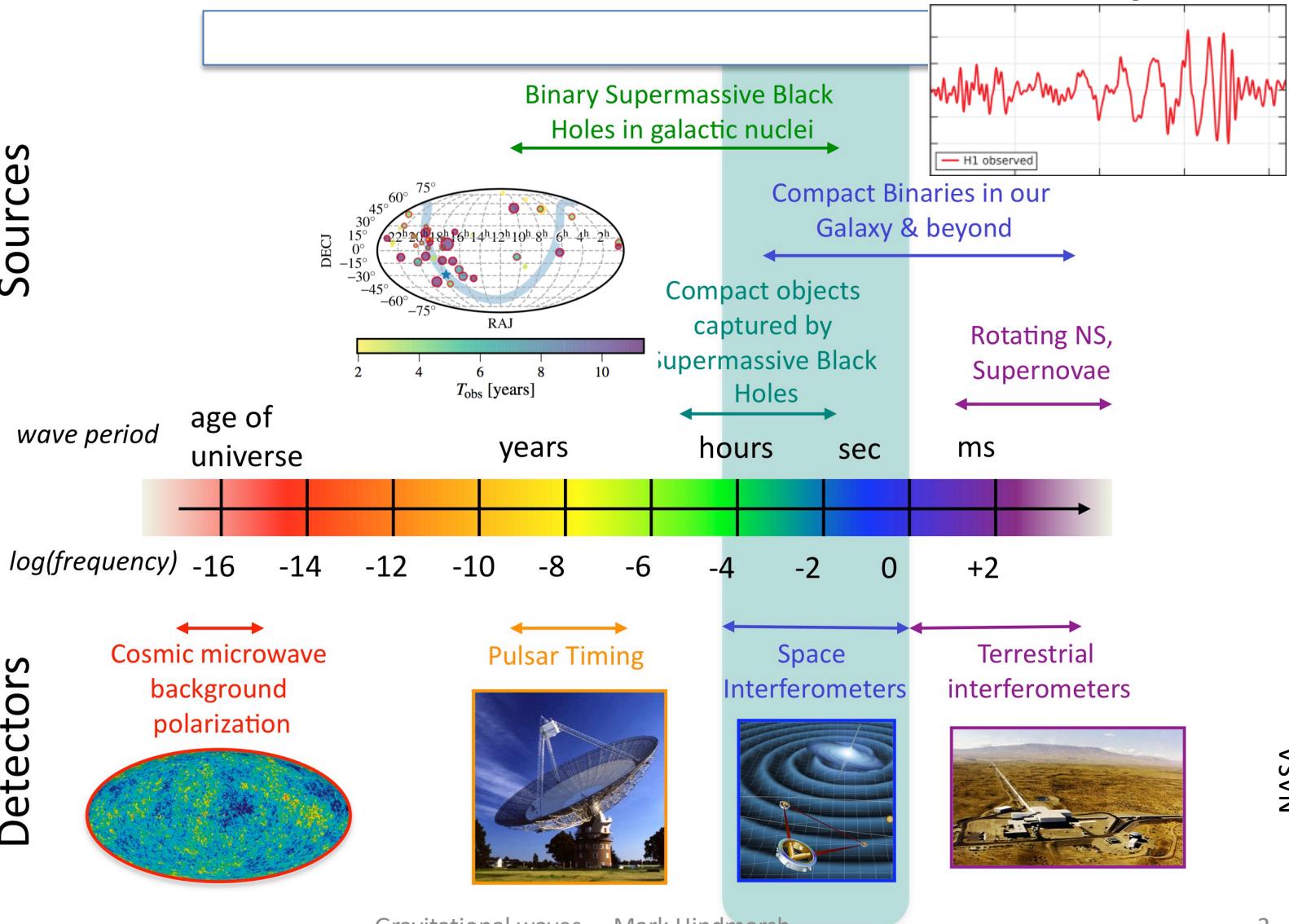


IDS workshop  
Jyväskylä  
3. kesäkuuta 2019



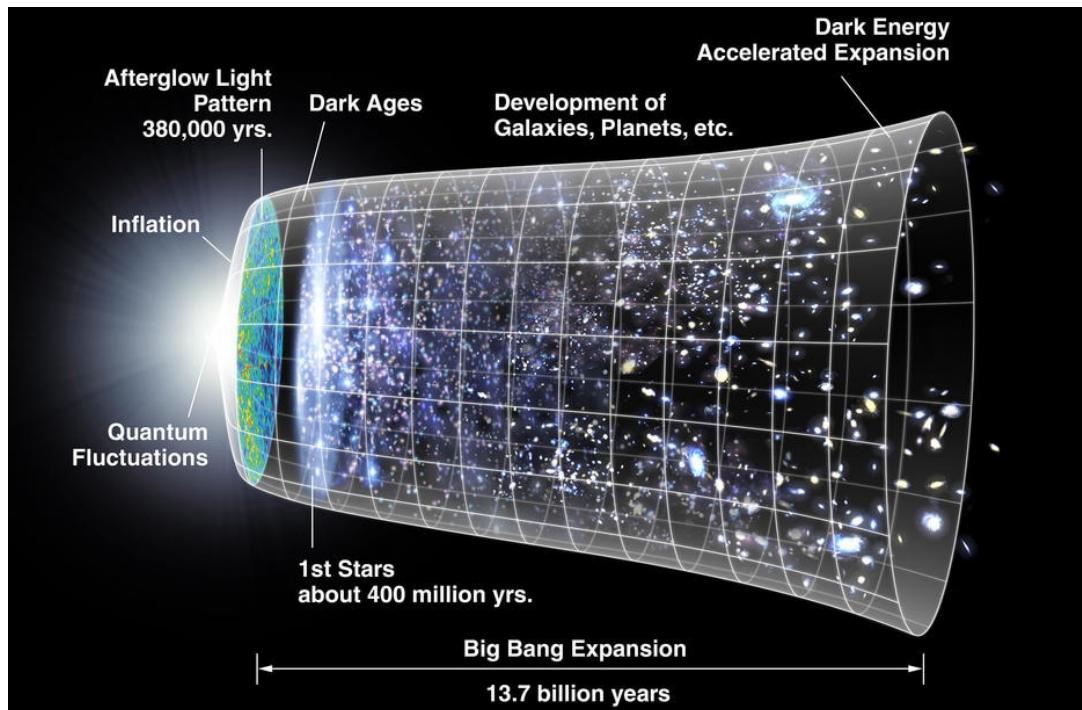
# Gravitational wave astronomy

## Sources



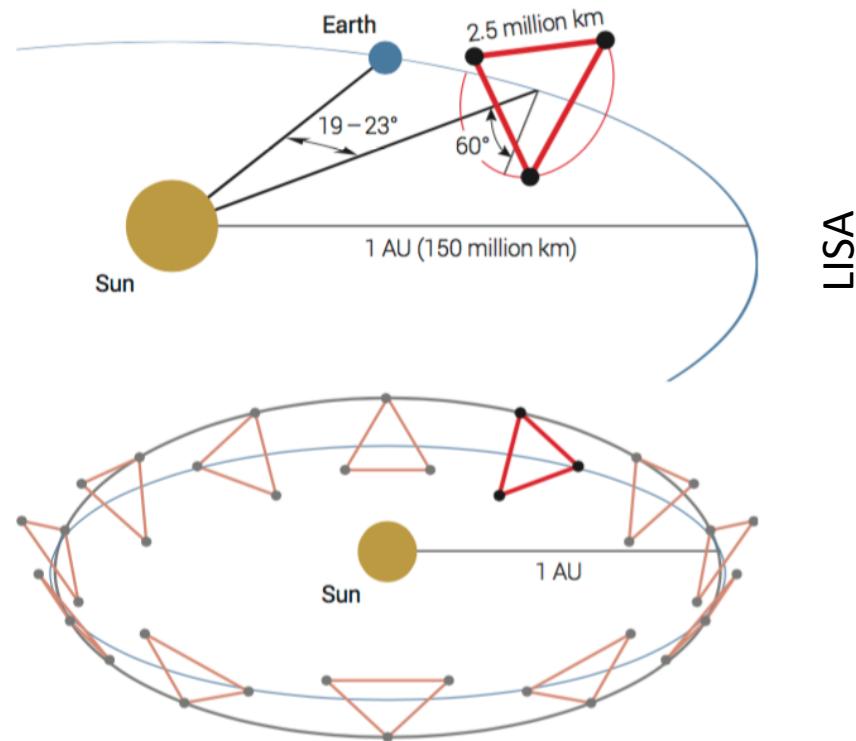
# Gravitational wave cosmology

- Gravitational waves are hard to observe
- Once made, not absorbed by intervening matter
- Complete history of the universe visible in GWs



# Space-based gravitational wave detectors

- Approved:
  - LISA (ESA L3 2034)
- LISA sensitivity
  - Peak:  $10^{-3}$  -  $10^{-2}$  Hz
  - Arm length  $l \approx 10^9$  m
  - $\Delta l \approx 10^{-12}$  m
  - Characteristic strain  $h \approx 10^{-21}$
- Proposed:
  - DECIGO (Japan, ?)
  - Taiji, Tianqin (China, ?)
  - Big Bang Observer (USA, ?)



# Gravitational waves from the early universe

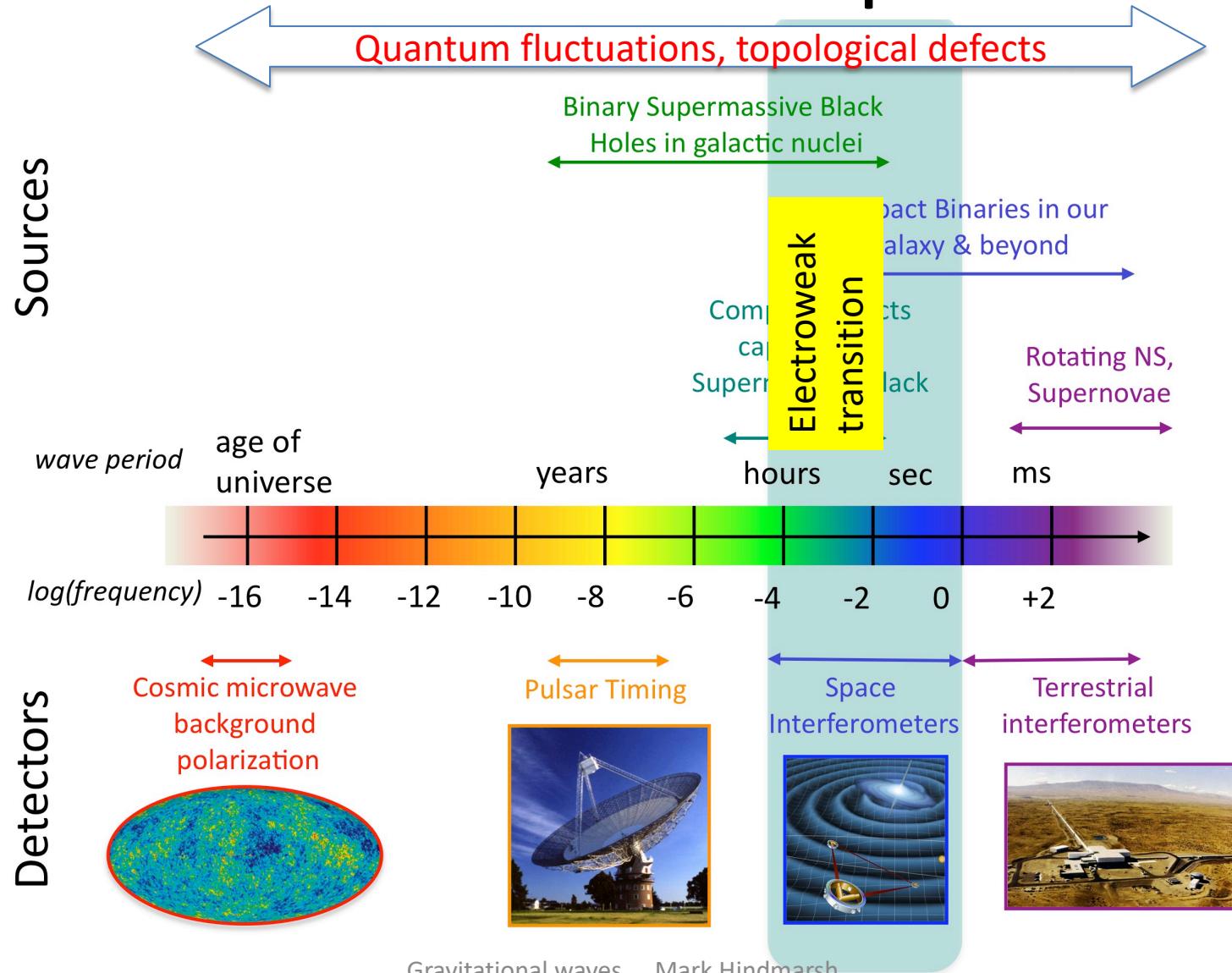
- Events at time  $t$  generate waves with minimum frequency  $f \approx 1/t$  (Hubble rate)
- Redshifted to a frequency now:  $f_0 = [a(t)/a(t_0)] f$
- Minimum frequencies (redshifted Hubble rates):

Event	Time/s	Temp/GeV	$f_0/\text{Hz}$
QCD phase transition	$10^{-3}$	0.1	$10^{-8}$
EW phase transition	$10^{-11}$	100	$10^{-5}$
?	$10^{-25}$	$10^9$	100 <b>LIGO</b>
End of inflation	$\geq 10^{-36}$	$\leq 10^{16}$	$\leq 10^8$

- Inflation and topological defects: waves on all scales

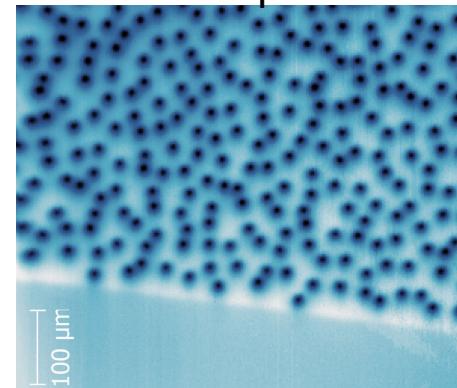
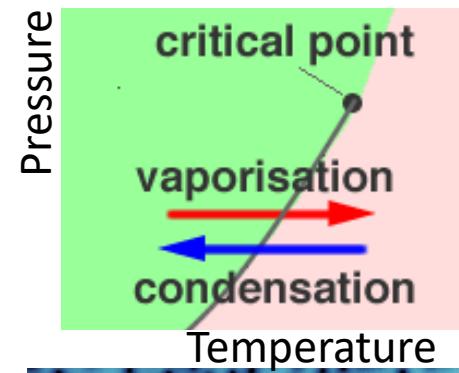
Caprini, Figueroa (2018), Christensen (2019)

# Gravitational wave spectrum



# Phase transitions in the early Universe

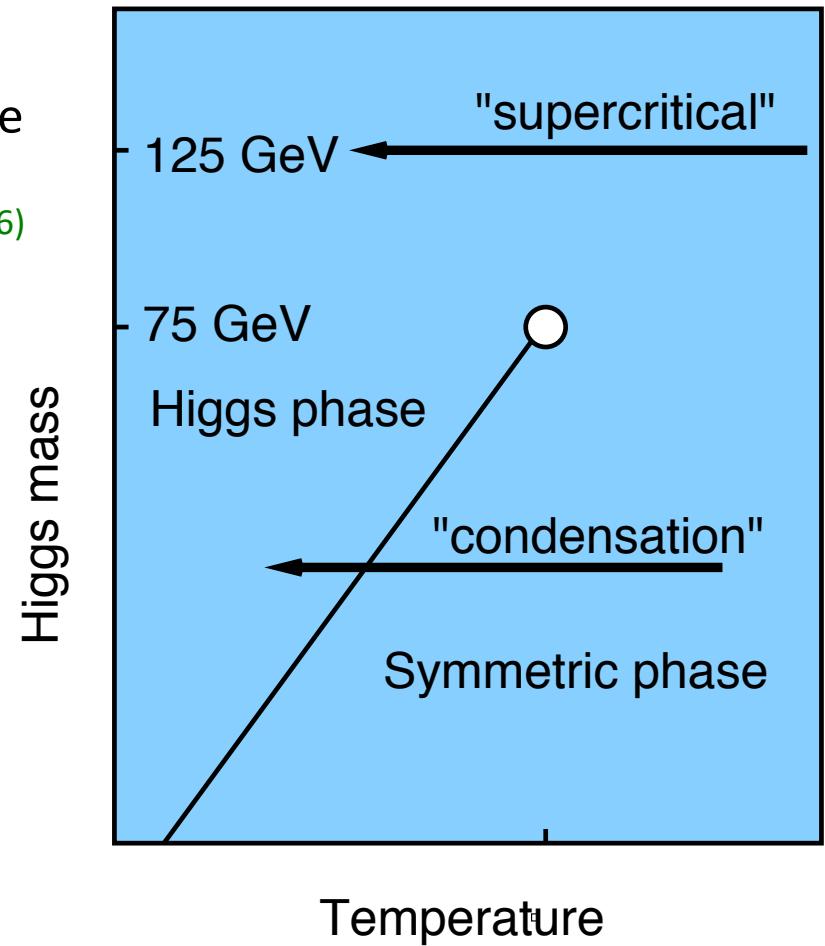
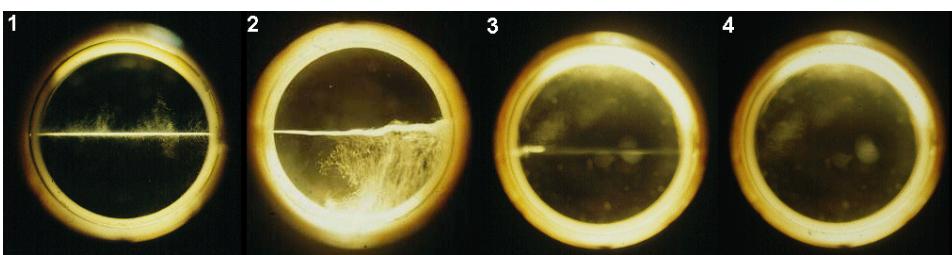
- At very high temperatures and pressures, the state of matter in the Universe changes
  - $T_c \sim 100$  MeV QCD (cross-over)
  - $T_c \sim 100$  GeV Higgs/electroweak
  - $T_c \gg 100$  GeV ???
- Departures from equilibrium and homogeneity (shear stress)
  - First order phase transition: relativistic condensation or 'fizz'  
Steinhardt (1982)
  - Formation of topological defects  
Kibble (1976)



Abrikosov vortices

# Electroweak transition

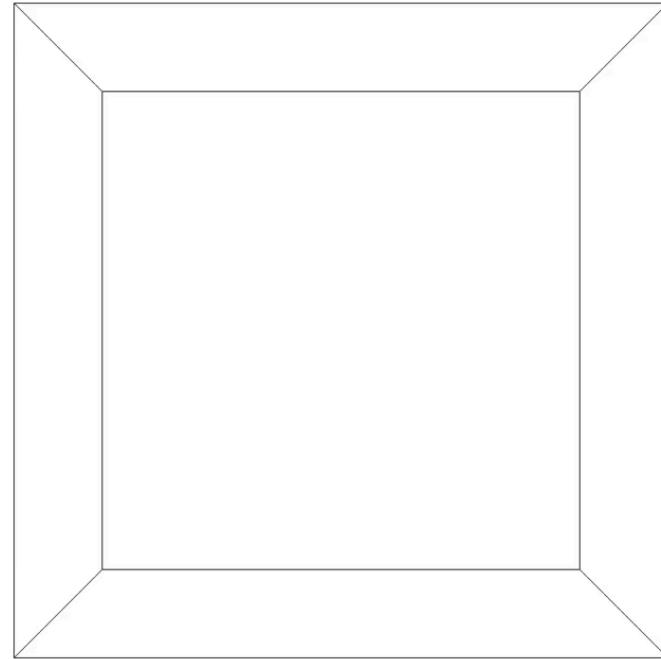
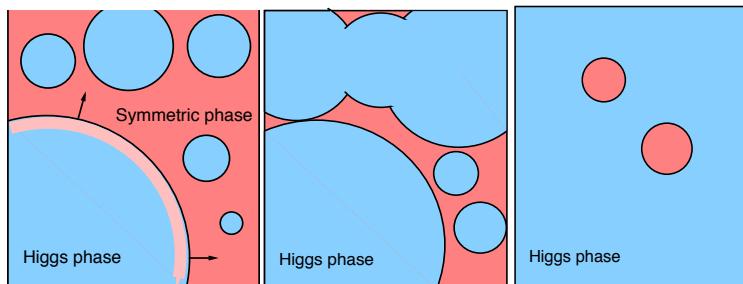
- SM is not weakly coupled at high T
- Non-perturbative techniques:
  - Dimensional reduction to 3D effective field theory + 3D lattice  
Kajantie, Laine, Rummukainen, Shaposhnikov (1995,6)
  - SU(2)-Higgs on 4D lattice  
Czikor, Fodor, Heitger (1998)
- SM transition at  $m_h \approx 125$  GeV is a cross-over - **a supercritical fluid**



- Search for 1<sup>st</sup> order transition is a search for physics beyond SM

# First order phase transitions

- 1st order transition  
proceeds by nucleation of  
bubbles of Higgs phase
- Expanding bubbles generate  
pressure waves in hot fluid
- Shear stresses - detectable  
gravitational waves?



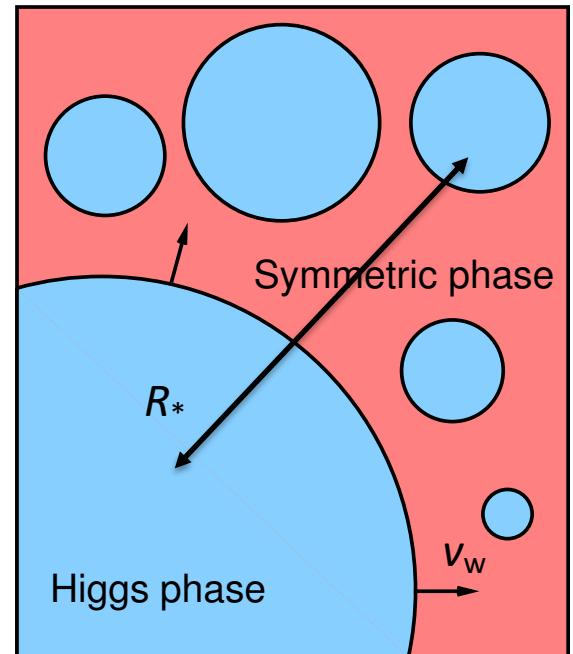
Kinetic energy density

Steinhardt (1982); Gyulassy et al (1984);  
Witten (1984); Enqvist et al (1992); Gravitational waves ... Mark Hindmarsh

# GWs from first order phase transitions

- Parametrise transition:
  - $T_n$  = nucleation temperature
  - $\alpha$  = (scalar potential)/(thermal energy)
  - $R_*$  = mean bubble centre separation
  - $v_w$  = bubble wall speed
  - $g_{\text{eff}}$  = effective d.o.f.
- Notes on calculating parameters:
  - Bubble nucleation rate/volume =  $p(T)$
  - Transition rate parameter  $b = d \ln(p)/dt$
  - $v_w$  non-equilibrium
  - Calculate  $T_n, \alpha, R_*$  from  $V_T(f)$ ,  $b$  and  $v_w$
- Aim: GW power spectrum

$$\frac{d\Omega_{\text{gw}}}{d \ln f}(T_n, R_* H_n, \alpha, v_w, g_{\text{eff}}^\pm)$$



# GWs from phase transitions

- Gravitational waves generated by shear stress fluctuations

$$\Omega_{\text{GW}} \sim \frac{1}{G\rho} \left\langle \left| \dot{h}_{ij}(t) \right|^2 \right\rangle$$

$$\dot{h}_{ij} \sim G \int dt' \cos[k(t-t')] T_{ij}^{TT}(k, t')$$

$$T_{ij}^{TT} \sim (\rho + p) U_i U_j \sim \rho K$$

–  $K(\alpha, v_w)$  = fluid kinetic energy fraction

- Shear stress  $\sim$  kinetic energy

- Kinetic energy from scalar potential energy  $T_{ij}^{TT} \sim (\rho + p) U_i U_j \sim \rho K$

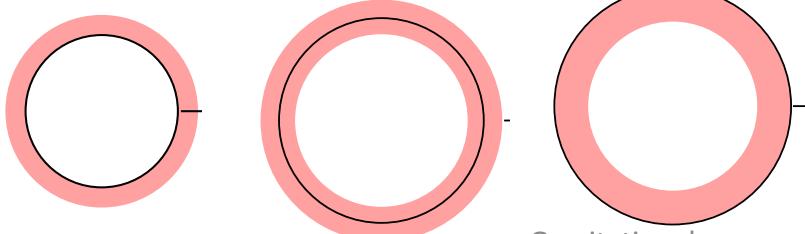
- Timescales  $\tau_v$  and  $\tau_c$

- $\tau_v$  duration of stresses from fluid velocity
- $\tau_c$  coherence time of stress fluctuations

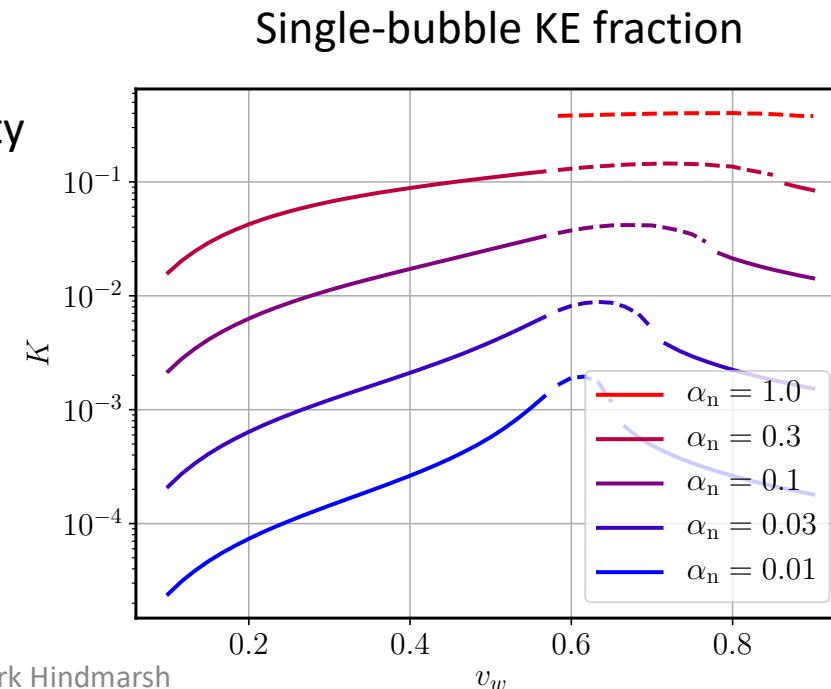
$$\Omega_{\text{gw},0} \sim \Omega_{\text{rad},0} (H_n \tau_v) (H_* \tau_c) K^2$$

- KE fraction estimate from single bubble

Kamionkowski et al 1994, Espinosa et al 2010



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# Simulations of thermal phase transitions

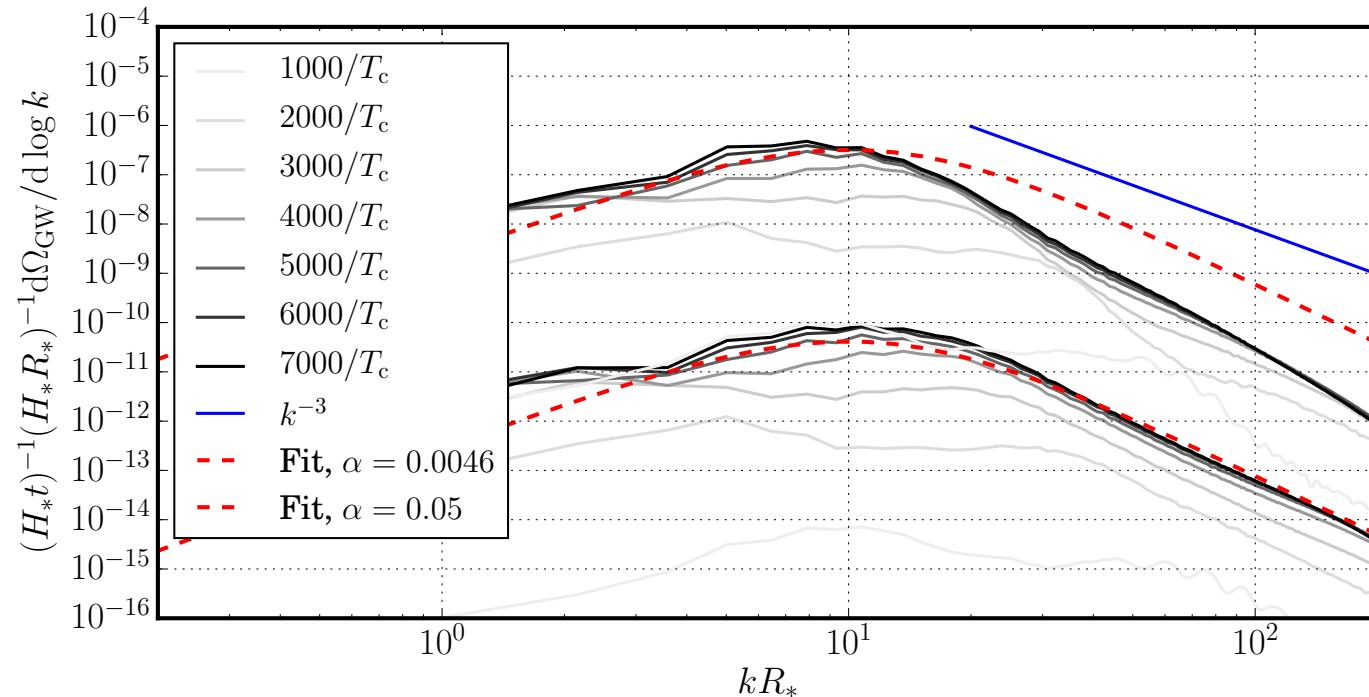
- Preparatory: 1M hrs CSC, Finland
- 2015/6: 17M CPU-hours  
Tier-0 (Hazel Hen, Stuttgart)
- $4200^3$  lattice on 24k cores
- GW density fraction power spectrum

$$\frac{d\Omega_{\text{gw}}}{d \ln k} = \frac{1}{12H^2} \frac{k^3}{2\pi^2} \langle |\dot{h}_{ij}^{\text{tt}}(\mathbf{k})|^2 \rangle$$



Fluid kinetic energy in slice ( $1200^3$ )

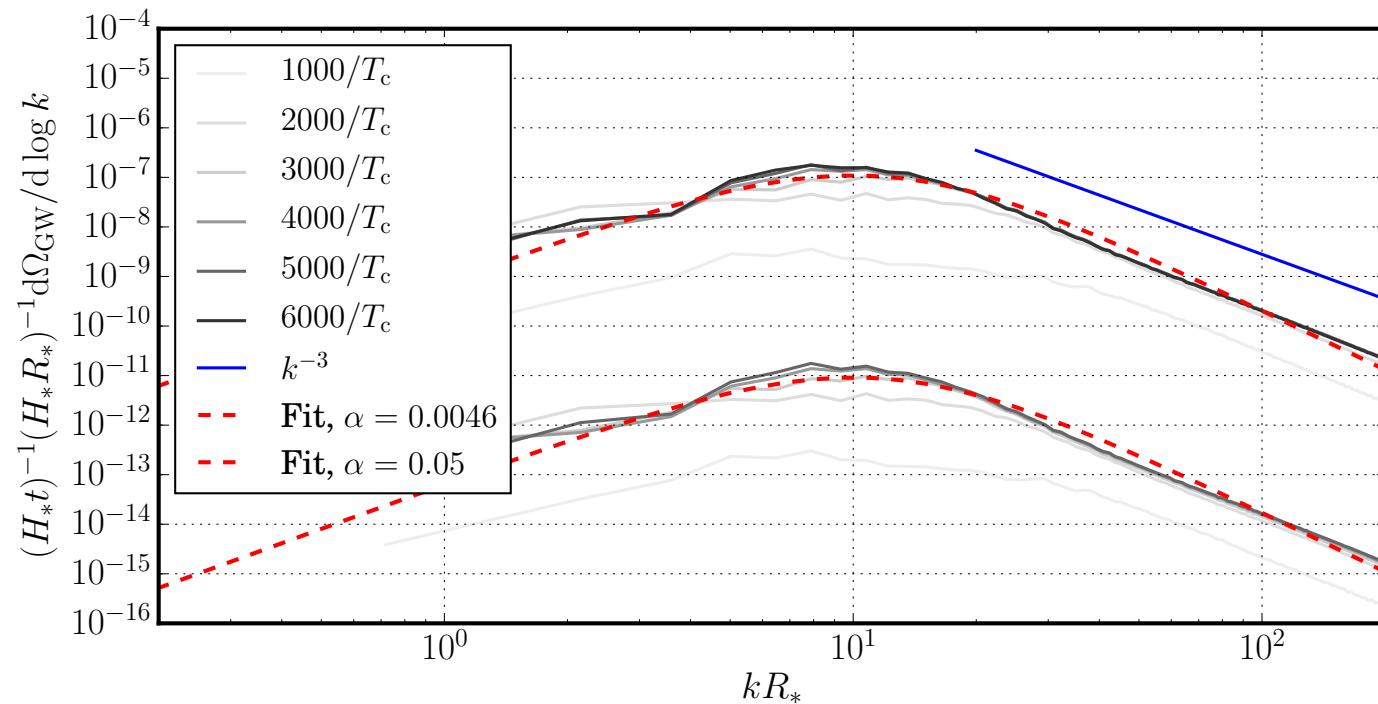
# GW power spectra: deflagration



- Transition strength:  
 $\alpha = 0.0046, 0.05$
- Wall speed:  
 $v_w = 0.44$
- Mean bubble separation:  
 $R_* = 1900/T_c$
- Domed peak at  $kR_* \sim 10$
- Approx  $k^{-4}$  spectrum at high  $k$

Hindmarsh, Huber, Rummukainen, Weir (2017)

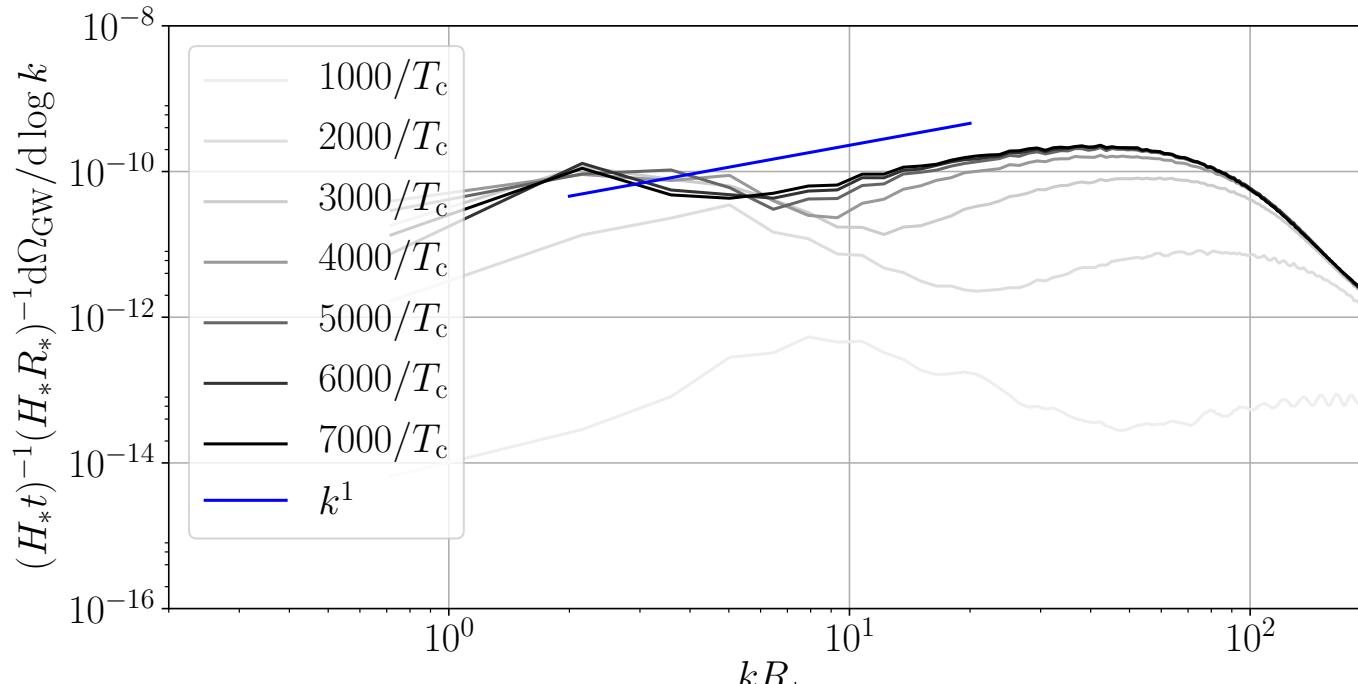
# GW power spectra: detonation



- Transition strength:  
 $\alpha = 0.0046, 0.05$
- Wall speed:  
 $v_w = 0.92$
- Mean bubble separation:  
 $R_* = 1900/T_c$
- Domed peak at  $kR_* \sim 10$
- Approx  $k^{-3}$  spectrum at high  $k$

Hindmarsh, Huber, Rummukainen, Weir (2017)

# GW power spectra: near $c_s$



- Transition strength:  
 $\alpha = 0.0046$
- Wall speed:  
 $v_w = 0.56$
- Mean bubble separation:  
 $R_* = 1900/T_c$

- Peak at  $kR_* \sim 40$
- Rising plateau at  $kR_* \sim 10$   
(approx  $k$  spectrum)
- Peak length scale from  
sound shell thickness

Hindmarsh, Huber, Rummukainen, Weir (2017)

# Lifetime of sound waves

- Sound damped by viscosity

$$\left( \frac{4}{3} \eta_s + \zeta \right) \nabla^2 V_{\parallel}^i$$

- Shear viscosity  $\eta_s \sim T^3/e^4 \ln(1/e)$
- Lifetime of scale  $R$ :  $\tau_{\eta}(R) \sim e^4 \ln(1/e) R^2 T$ .
- Longer than Hubble time for scales

$$R \gg \frac{v_w}{H_*} \left( \frac{T_c}{m_{Pl} e^4} \right) \sim 10^{-11} \frac{v_w}{H_*} \left( \frac{T_c}{100 \text{ GeV}} \right),$$

- Effective lifetime is Hubble time, unless ...
- ... non-linearities develop, at  $\tau_{nl} \sim R_*/\bar{U}_f$
- Pure acoustic waves only if  $\bar{U}_f < R_* H_*$

# Acoustic GW power spectrum

Hindmarsh, Huber, Rummukainen, Weir (2017)

- Density fraction of GWs by acoustic production:

$$\Omega_{\text{gw},0}(f) = 0.68 F_{\text{gw},0} \tilde{\Omega}_{\text{gw}}(H_* R_*) K^2 C \left( \frac{f}{f_{p,0}} \right)$$

- Fluid kinetic energy density fraction  $K(\alpha, v_w)$   
(estimate from single-bubble fluid kinetic energy) Espinosa et al (2010)

- Dimensionless constant (from simulations)  $\tilde{\Omega}_{\text{gw}} \sim 0.01$

- Phenomenological fit  $C(s) = s^3 \left( \frac{7}{4 + 3s^2} \right)^{\frac{7}{2}}$

- Peak frequency

$$f_{p,0} \simeq 26 \left( \frac{1}{H_* R_*} \right) \left( \frac{T_n}{10^2 \text{ GeV}} \right) \left( \frac{g_{\text{eff}}}{100} \right)^{\frac{1}{6}} \mu\text{Hz},$$

- Matter-era dilution:

$$F_{\text{gw},0} = (3.57 \pm 0.05) \times 10^{-5} \left( \frac{100}{g_{\text{eff}}} \right)^{\frac{1}{3}}.$$

- NB  $\alpha$  from  $\Delta(e - 3p)/4$ , not latent heat  $\Delta(e + p)$

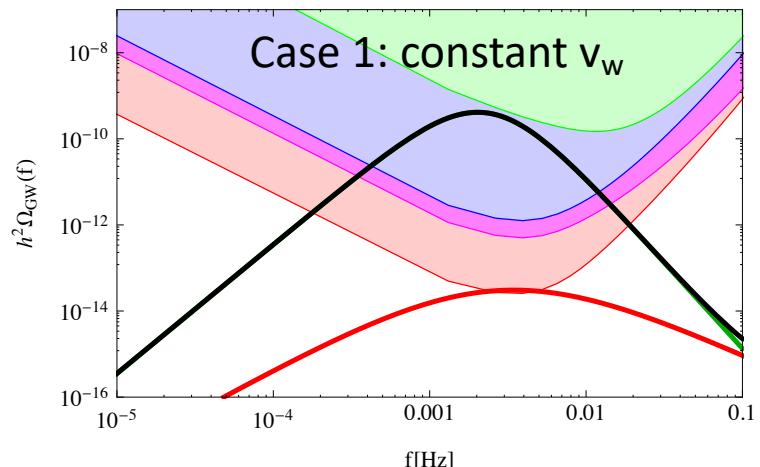
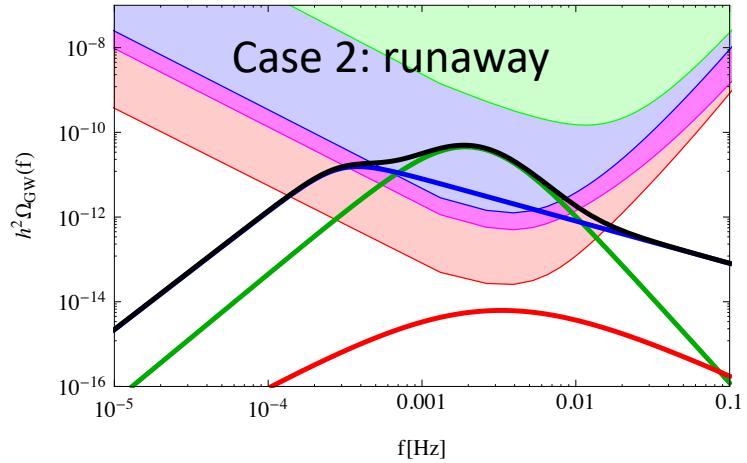
- Latent heat 4x larger at  $T_c$

# LISA CWG party line 2016

- Three contributions to total power:
  - Scalar field f
  - Acoustic ac
  - Turbulent tu

$$\Omega_{\text{gw}} = \Omega_{\text{gw}}^\phi + \Omega_{\text{gw}}^{\text{ac}} + \Omega_{\text{gw}}^{\text{tu}}$$

- Scalar field: bubble wall collisions
  - relevant only for runaway walls
  - “envelope approximation”
  - Kosowsky, Turner 1992
  - Huber, Konstandin 2008
- Acoustic production:
  - M.H. et al 2013, 2015, 2017
- Turbulent production:
  - Caprini, Durrer, Servant 2009



# Developments 1: scalar field

- Numerical simulations show differences from envelope approximation

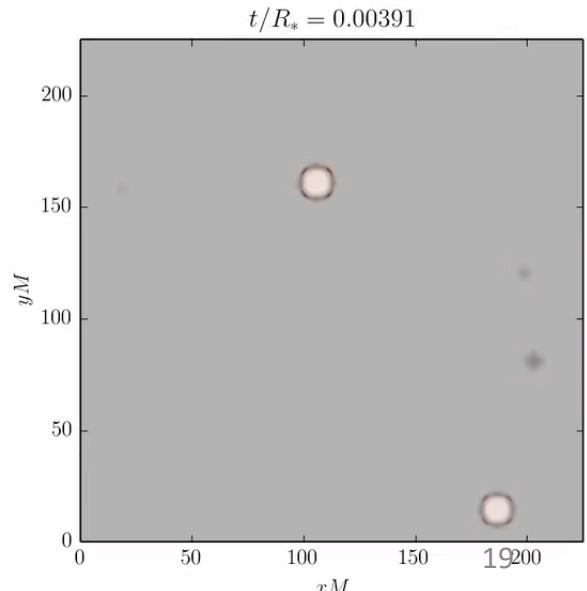
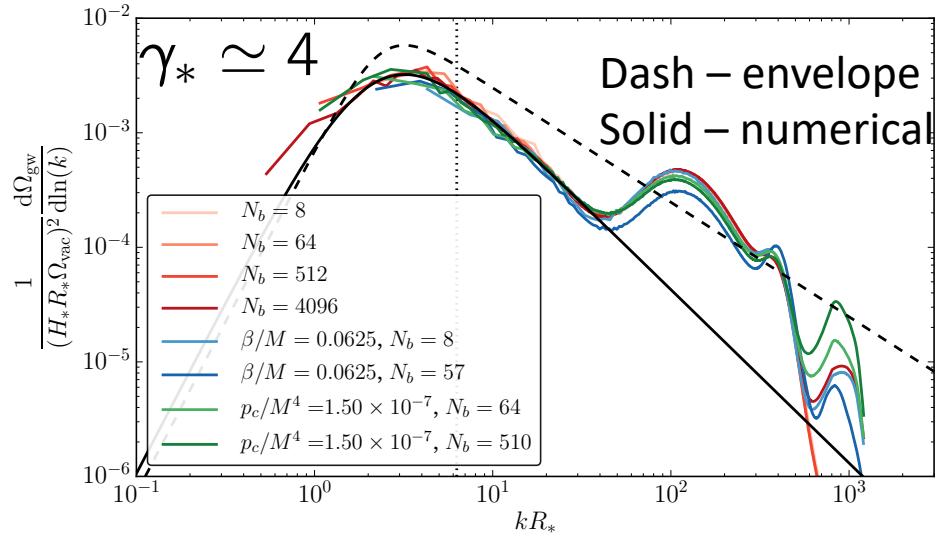
Cutting, MH, Weir 2018

$$\frac{d\Omega_{\text{gw}}^{\text{fit}}}{d\ln k} = \Omega_p^{\text{fit}} \frac{(3+b)^c \tilde{k}^b k^3}{(b\tilde{k}^{(3+b)/c} + 3k^{(3+b)/c})^c}$$

$$\Omega_p^{\text{fit}} = (3.22 \pm 0.04) \times 10^{-3} (H_n R_*)^2 \Omega_\phi^2,$$

$$\tilde{k}R_* = 3.20 \pm 0.04,$$

$$b = 1.51 \pm 0.04, \quad c = 2.18 \pm 0.15$$

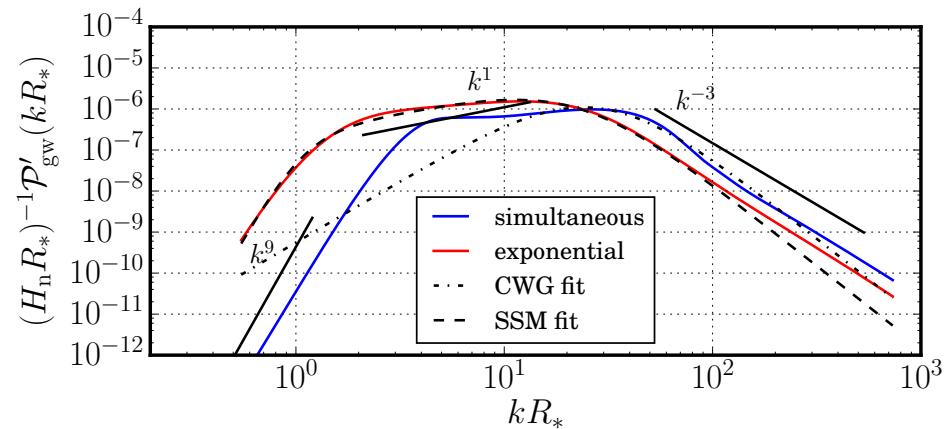
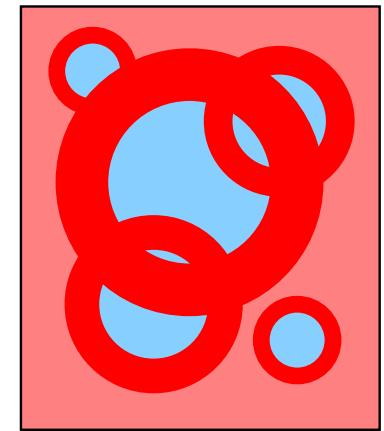
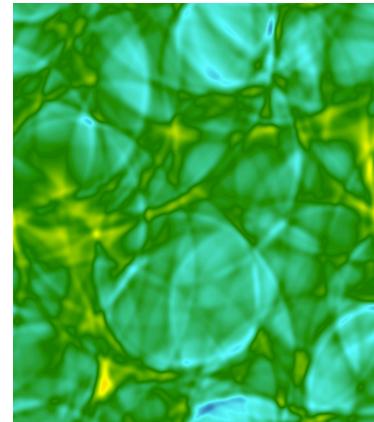


# Developments 2: sound shell model

- Gaussian velocity field from weighted addition of sound shells  $\mathbf{v}_q(t_i)$

MH 2017, MH, Hijazi (in prep 2019)

- Two length scales:
  - Bubble spacing  $R_*$
  - Shell width  $R_* |v_w - c_s|$
- Double broken power law
  - $P_{gw} \sim k^9, k^1, k^{-3}$
- Amplitude 10% - 20% agreement w/ simulations
- Systematically too large for deflagrations

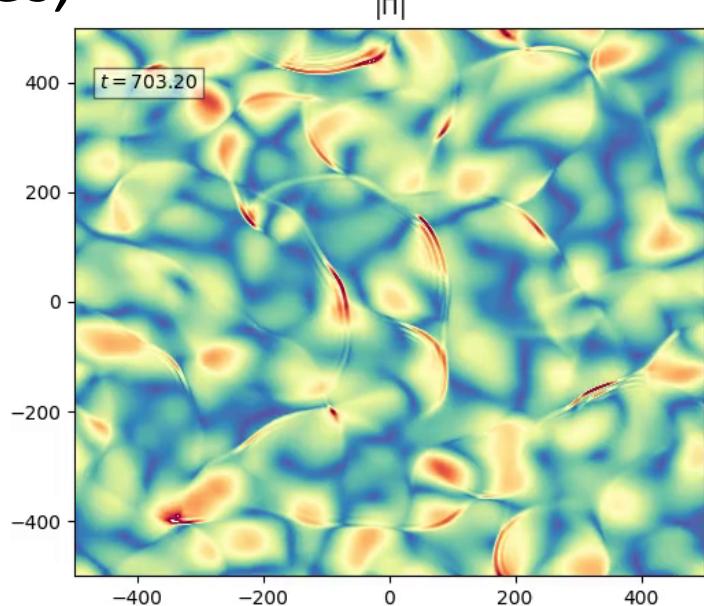
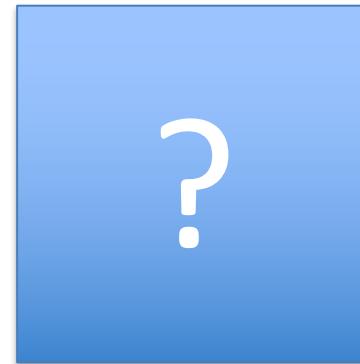


# Developments 3: non-linearities

- Fluid equations: non-linearities important after

$$\tau_{\text{nl}} \sim L_f / \bar{U}_f$$

- Energy transport to small scales, where dissipated.
- Longitudinal  $v$ 
  - Wave turbulence
  - Shocks Pen, Turok 2015
- Transverse  $v$ 
  - Vorticity production
  - Turbulence Caprini, Durrer, Servant 2008  
Gogoberidze, Kahniashvili, Kosowsky 2007



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# Non-linearities: dealing with uncertainty

- Non-linearities important after

$$\tau_{\text{nl}} \sim L_f / \bar{U}_f$$



## Acoustic issues

- CWG 2016: Non-linear dissipation ignored
- source lifetime assumed to be  $H_n^{-1}$
- **Conservative fix:** multiply PS by  $\min(1, H_n \tau_{nl})$

## Turbulence issues

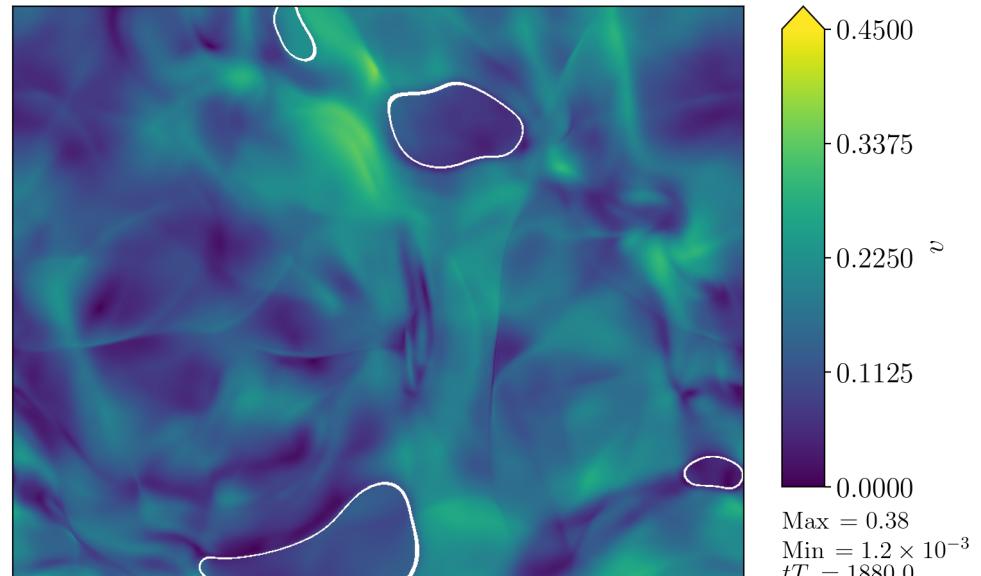
- CWG 2016:  $K_{tu} = 0.05K$
- but we don't really know  $K_{tu}/K$
- Disagreement on turbulence PS between CDS and GKK
- GW production from (longitudinal \* transverse)
- **Conservative fix:** take  $K_{tu} = 0$

# ‘Strong’ phase transitions

- Recent simulations of strong transitions with  $\alpha = 0.5$
- Relativistic velocity flows with maximum velocity  $V > 0.6$
- During collision of bubble walls rotational velocity  $V_{\perp}$  can be generated.
- Deflagrations visibly produce more  $V_{\perp}$  than detonations.
- Suppression of kinetic energy and GW signal, especially deflagrations

Cutting, Hindmarsh, Weir (in prep.)

See D. Cutting’s talk at 14.40

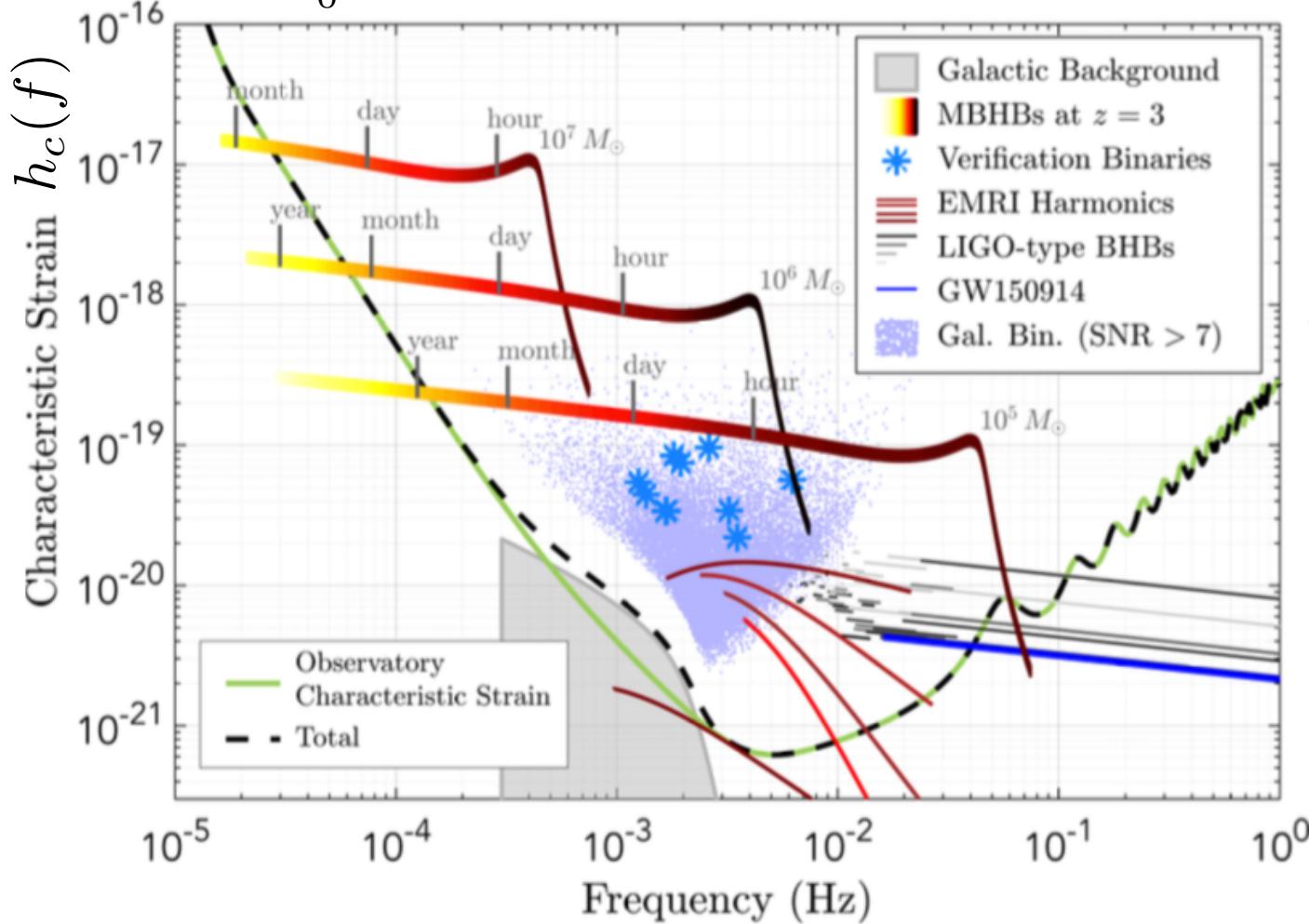


$V$  fluid (slice)  
White = bubble wall  
 $v_w=0.44$   
Lattice  $960^3$

# Laser Interferometer Space Antenna

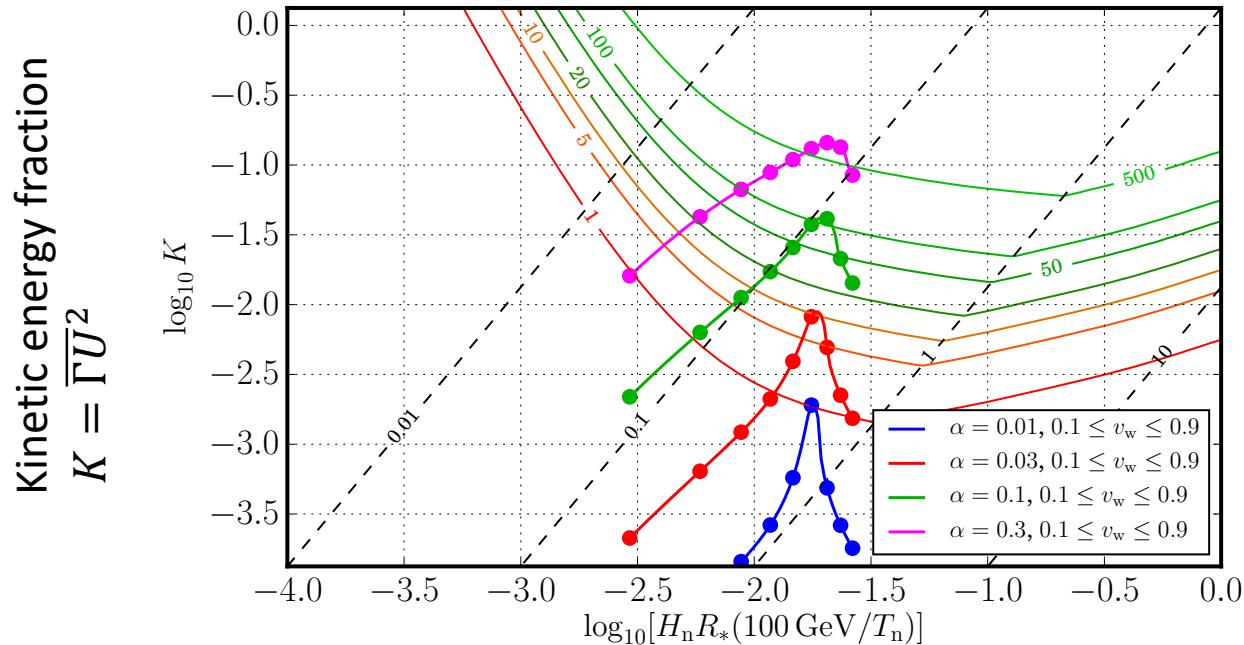
$$\Omega_{\text{gw},0}(f) = \frac{2\pi^2}{3} \frac{f^2}{H_0^2} h_c(f)^2$$

$$h_c(f) \simeq 0.87 \times 10^{-18} (f/\text{Hz})^{-1} [\Omega_{\text{gw}}(f)]^{1/2}$$



LISA proposal, Amaro-Seoane et al 2017

# Estimated LISA prospects



- Estimate signal-to-noise ratio  $r$ 

$$\rho^2 = T_{\text{obs}} \int df \left( \frac{\Omega_{\text{gw}}(f)}{\Omega_{\text{noise}}(f)} \right)^2$$
  - Observation time 4 years
  - Neglect foregrounds
    - White Dwarf binaries (annual variation)
    - LIGO BHB precursors (negligible)

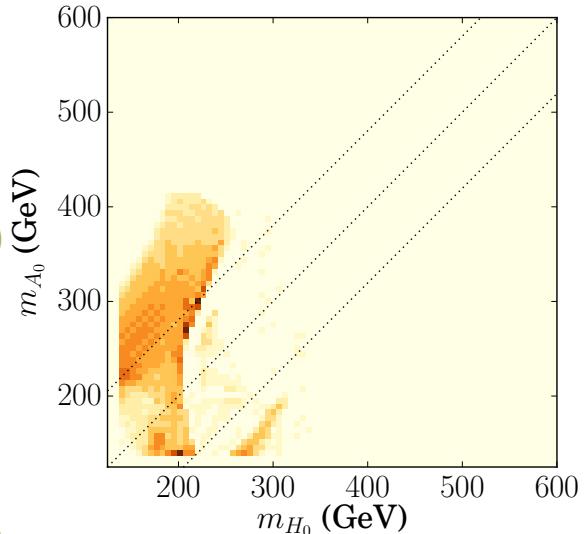
- E.g. (favourable cases):
  - $\beta/H_n = 100$ ,  $T_n = 100$  GeV
- NB  $\alpha > 0.1$  highly uncertain  
Cutting, Hindmarsh, Weir (in prep.)
  - But important region for LISA  
Ellis, Lewicki, No (2018)

# Phase transitions: future challenges

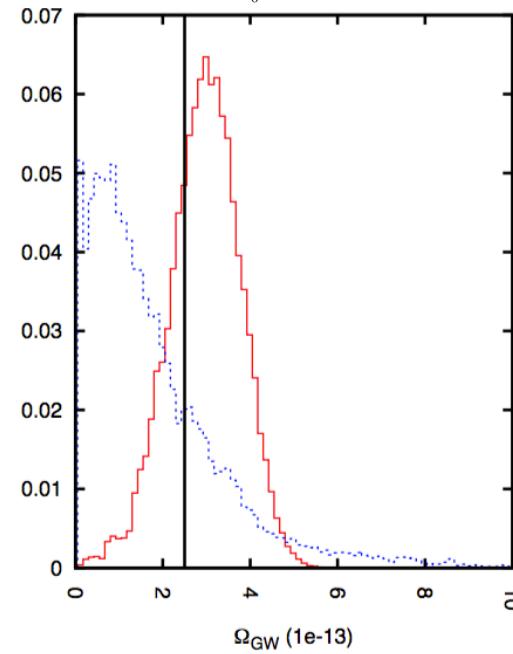
- Full characterization of GW spectrum from phase transitions
  - Strong transitions Cutting, MH, Weir (in prep)
  - Magnetic fields? Pol et al (2019), Zhang et al (2019)
- Phase transition parameters from underlying particle physics models
  - E.g. transition strength in DR & non-perturbative 3D Moore, Rummukainen (2000)  
Laine, Nardini, Rummukainen (2012)  
Gould et al (2019)  
Kainulainen et al (2019)
- Connection to collider (LHC ...) data
  - E.g.  $\lambda_{hhh}$  Noble, Perelstein 2007; Patel, Ramsey-Musolf 2013
  - $A^0 \rightarrow H^0 Z$  Dorsch, Huber, Mimasu, No (2016)  
Andersen et al (2017)
- Distinguishing phase transitions from astrophysical GW foregrounds

Adams, Cornish (2013), Hashino et al (2018)

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Andersen et al 2017

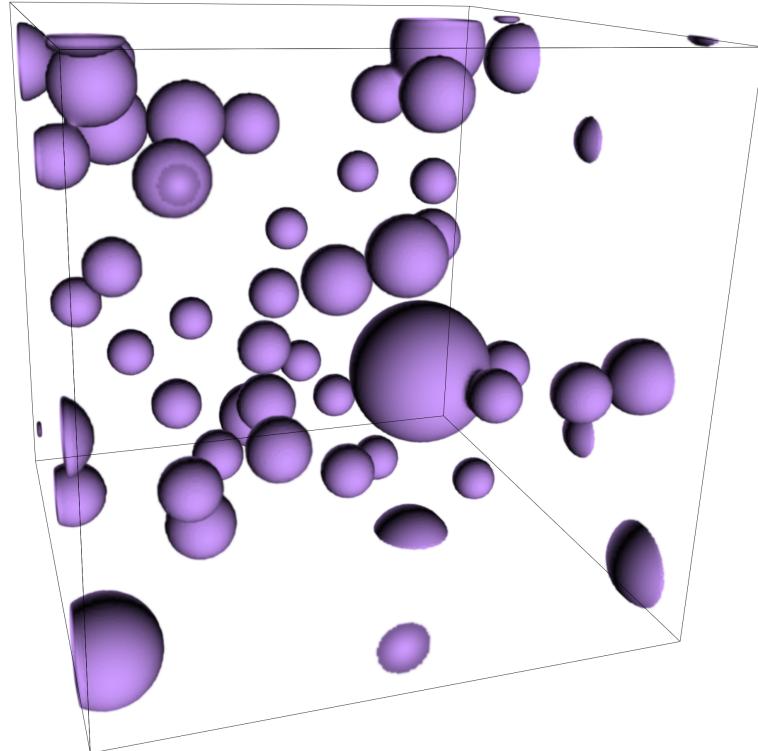


Adams, Cornish (2013)

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# Summary

- GWs probe of physics of early universe at very high energy
- LISA will probe physics of Higgs phase transition from 2034
- Measure/constrain phase transition parameters
  - $T_n$  = nucleation temperature
  - $\alpha$  = (scalar potential)/(thermal energy)
  - $R_*$  = mean bubble centre separation
  - $v_w$  = bubble wall speed
  - $g_{\text{eff}}$  = effective d.o.f.
- Need 1: accurate calculations of GW power spectrum from parameters
- Need 2: accurate calculations of parameters from underlying particle physics models
- Need 3: reliable methods to extract parameters from real data
- Lots of work to do
  - only 16 years to go!



# Numerical simulation of phase transitions

- Ingredients:

Ignatius et al (1994), Kurki-Suonio, Laine (1996)

- Higgs field  $-\ddot{\phi} + \nabla^2\phi - \frac{\partial V}{\partial\phi} = \eta W(\dot{\phi} + V^i\partial_i\phi)$ 
  - $\eta$  coupling to fluid (models energy transfer, friction)
  - $V_i$  velocity,  $W$   $\gamma$ -factor

- Relativistic fluid equation

$$\dot{E} + \partial_i(EV^i) + P[\dot{W} + \partial_i(WV^i)] - \frac{\partial V}{\partial\phi}W(\dot{\phi} + V^i\partial_i\phi) = \eta W^2(\dot{\phi} + V^i\partial_i\phi)^2.$$

$$\dot{Z}_i + \partial_j(Z_iV^j) + \partial_iP + \frac{\partial V}{\partial\phi}\partial_i\phi = -\eta W(\dot{\phi} + V^j\partial_j\phi)\partial_i\phi.$$

- $E/W$  energy density,  $P$  pressure,  $Z_i$  momentum density,

- Discretisation

Wilson & Matthews (2003)

Different approach: Giblin, Mertens (2013)

- Metric perturbation

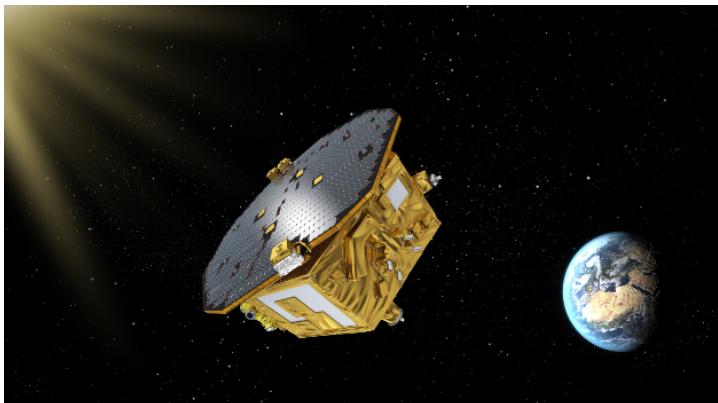
$$\ddot{u}_{ij} - \nabla^2u_{ij} = 16\pi GT_{ij}$$

Garcia-Bellido, Figueroa, Sastre (2008)

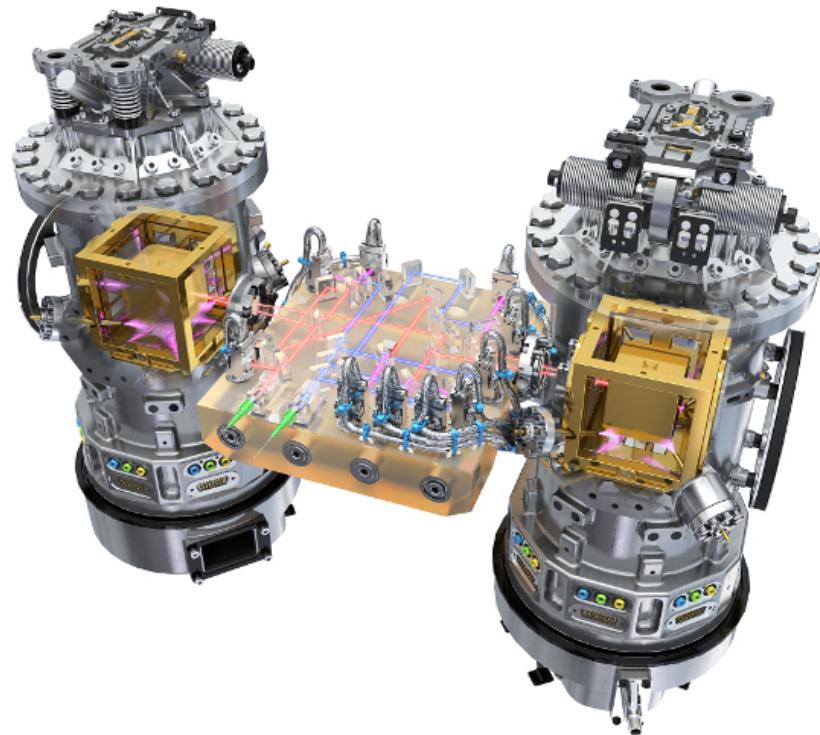
$$\dot{h}_{ij}^{tt}(\mathbf{k}) = \Lambda_{ijkl}^{tt}(\mathbf{k})\dot{u}_{ij}(\mathbf{k})$$

# LISA Pathfinder

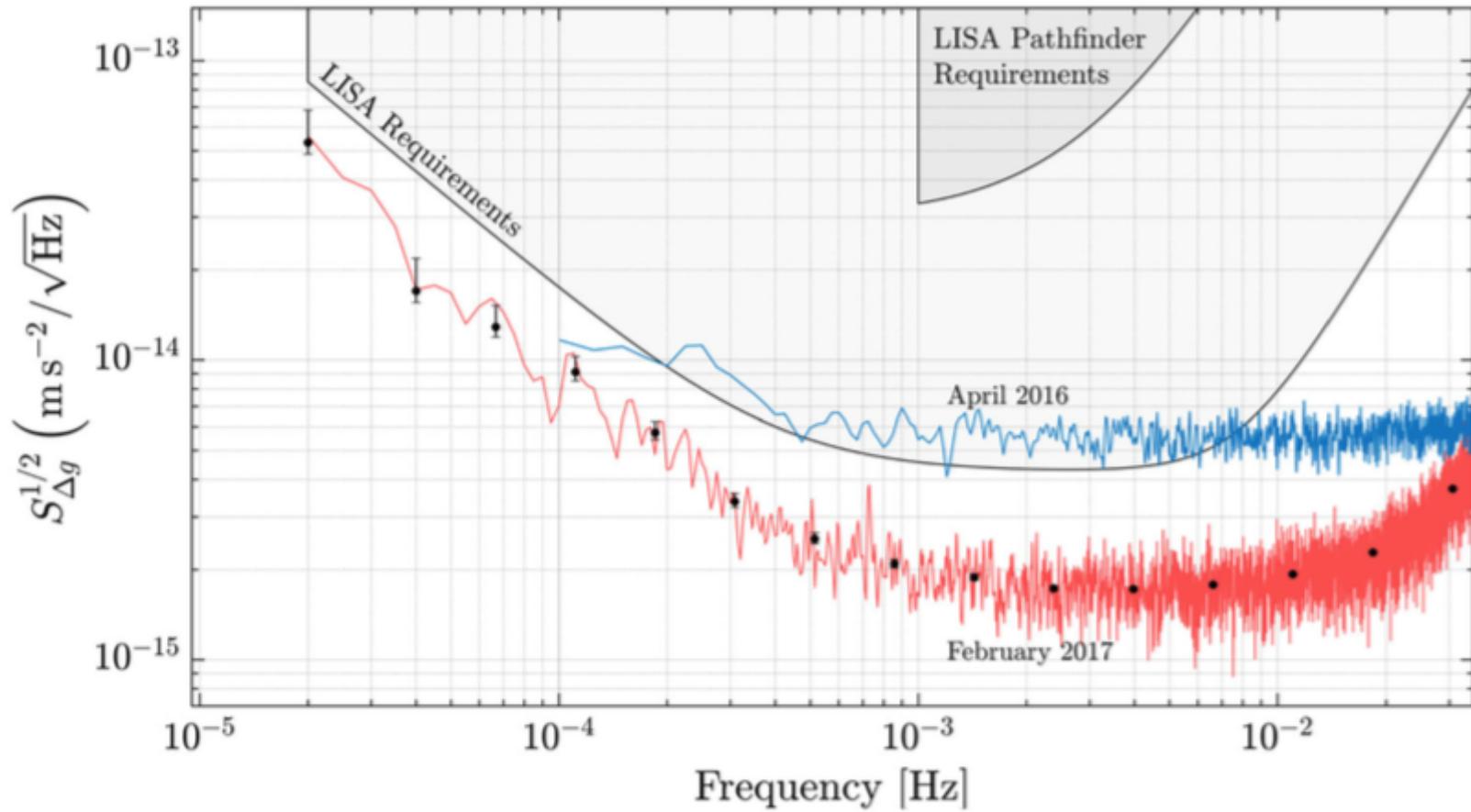
- Technology tester for LISA
- Test masses in free-fall at L1
- Control and measure motion
- Launched 3/12/15
- Masses released 3/2/16
- Mission end 30/6/17
- Quantified acceleration noise  
– how good is the free fall



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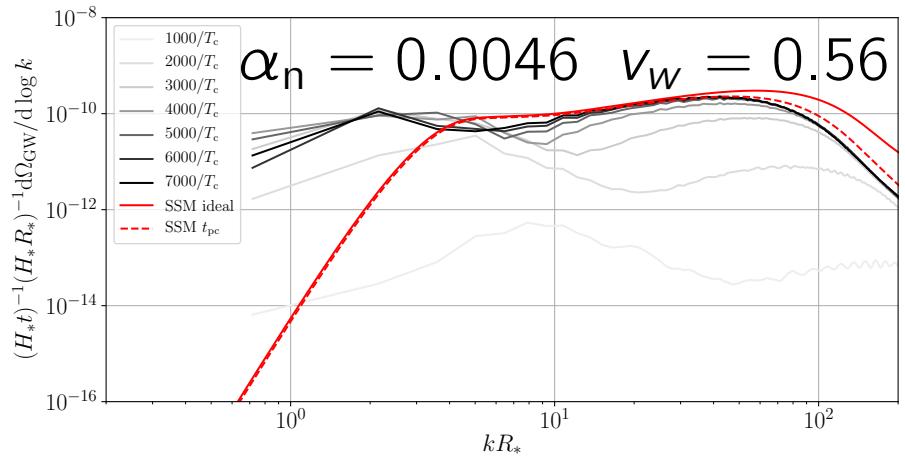
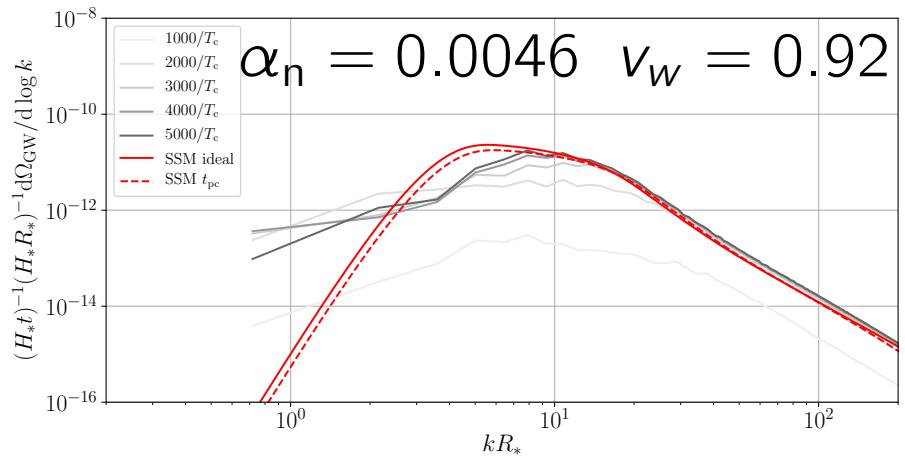
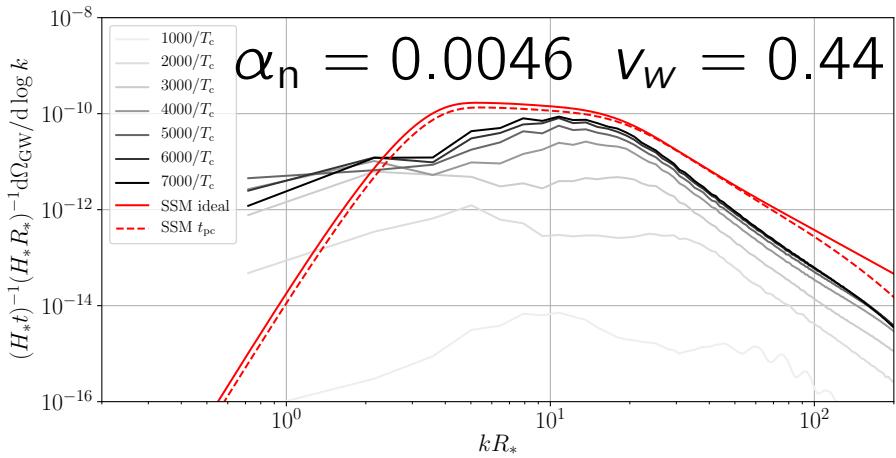
# LISA Pathfinder acceleration noise



# Sound shell model vs. simulations $P_{\text{gw}}$

- Solid: self-similar sound shell
- Dash: evolving sound shell at peak collision time
- Simultaneous nucleation

MH et al in prep 2019



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