

1. Show that the critical temperature for forming the condensate in the non-relativistic limit ( $Q/V \ll m^3$ , this is only an approximative condition) can be expressed as

$$T_c = \frac{2\pi}{m} \left( \frac{Q}{V} \right)^{2/3} \zeta\left(\frac{3}{2}\right)^{-2/3}$$

and in the relativistic limit ( $Q/V \gg m^3$ ) as

$$T_c = \left( \frac{3}{m} \frac{Q}{V} \right)^{1/2}.$$

2. Show carefully the following identities on fermionic integrals

$$\begin{aligned} \int \prod_{i=1}^N d\theta_i^* d\theta_i e^{-\theta_i^* A_{ij} \theta_j} &= \det A \\ \int \prod_{i=1}^N d\theta_i^* d\theta_k \theta_l^* e^{-\theta_i^* A_{ij} \theta_j} &= \det A (A^{-1})_{kl}. \end{aligned}$$

3. Show that in the low temperature limit  $m/T \gg 1$

$$J_T^\mp(m, T) \equiv \mp T \int \frac{d^3 p}{(2\pi)^3} \log(1 \mp e^{-\beta \omega_p}) \rightarrow J_T^\mp = T n(T),$$

where  $\omega_p^2 = \mathbf{p}^2 + m^2$  and  $n$  is the number density. Moreover, show that the  $J_T^-$ -function can be expanded as:

$$J_T^-(m, T) = -\frac{m^2 T^2}{2\pi^2} \sum_{n=1}^{\infty} \frac{1}{n^2} K_2\left(\frac{nm}{T}\right)$$

where  $K_2(x)$  is the Bessel function of the second kind.

4. Compute also the related integral:

$$I^-(m, T) \equiv \oint \Delta_0(\omega_n, \omega_p) = \int \frac{d^3 \mathbf{p}}{(2\pi)^3} \frac{1}{2\omega} (1 + 2n_B(\omega_p)) \equiv I_0^-(m) + I_T^-(m, T)$$

Show that the thermal part of this integral has the expression

$$I_T^-(m, T) = \frac{mT}{2\pi^2} \sum_{n=1}^{\infty} \frac{1}{n} K_1\left(\frac{nm}{T}\right).$$

Finally show also that  $mI_T^-(m, T) = \partial_m J_T^-(m, T)$ .

5. Show that in the high-temperature limit

$$J_T^-(m, T) = \frac{\pi^2 T^4}{90} - \frac{m^2 T^2}{24} + \frac{m^3 T}{12\pi} + \frac{m^4}{2(4\pi)^2} \left[ \log\left(\frac{me^{\gamma_E}}{4\pi T}\right) - \frac{3}{4} \right] - \frac{m^6 \zeta(3)}{3(4\pi)^4 T^2} + \mathcal{O}\left(\frac{m^8}{T^4}\right)$$

for bosonic integral. Note that the third term arises purely from the zero mode. Similarly for Fermions show that

$$J_T^+(m, T) = \frac{7}{8} \frac{\pi^2 T^4}{90} - \frac{m^2 T^2}{48} - \frac{m^4}{2(4\pi)^2} \left[ \log\left(\frac{me^{\gamma_E}}{\pi T}\right) - \frac{3}{4} \right] + \frac{7m^6 \zeta(3)}{3(4\pi)^4 T^2} + \mathcal{O}\left(\frac{m^8}{T^4}\right).$$

The two first 2-3 terms in the expansions are easy, but the logarithmic corrections are much harder. The difficulty comes from the fact that  $J$ 's are not analytic around  $m = 0$ . Look at section 2.3 of the book by Laine and Vuorinen or the paper by Doland and Jackw, which you find temporarily from the course homepage.