



JYVÄSKYLÄN YLIOPISTO  
UNIVERSITY OF JYVÄSKYLÄ

## Communication between theory and practice via deep learning and “deep teaching”

Math + X Symposium on Data Science and Inverse Problems in Geophysics

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Based on joint work with

**Maarten de Hoop and Vitaly Katsnelson**

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- Describing how deep learning might help experimentalists and theorists learn from each other.
- Finding a reliable way to reconstruct as much information as possible from the spectrum of free oscillations.
- The audience telling me whether any of this is possible.

# Outline

- 1 The problem
  - The spectrum of free oscillations
  - Geometry
  - Reconstructing geometry from spectrum
- 2 A mathematical result
- 3 Deep learning deep structures
- 4 Appendix A: Anisotropy and geometry
- 5 Appendix B: Ray transforms in low regularity and drums
- 6 Appendix C: More detailed spectral rigidity

# The spectrum of free oscillations



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- The oscillations are excited by large earthquakes.
- The amplitudes of different modes vary between different events, but the frequencies are always the same.
- The set of these frequencies is the spectrum of free oscillations.
- About 10 000 first frequencies are known.



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- Compare to the change of paradigm from Newton's to Einstein's gravity: The old view is that the Earth is on a curved path because of a force exerted by the Sun. The new view is that the Earth goes straight in a geometry influenced by the Sun.
- In this “elastic geometry” the distance between two points is the shortest travel time. All other geometrical quantities we need can be derived from length.
- There is one geometry for each wave speed. Due to several polarizations there are several geometries. In the simplest case there is a P-geometry and an S-geometry.

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This problem is hard in theory (mathematical proof of uniqueness) and practice (reliable numerical reconstruction from real data).

## Problem (Meta problem)

*How can theoretical and practical endeavors with the problem benefit from each other?*

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- We would like to reconstruct the geometry in the Cartesian coordinates. Alas, this is not possible in anisotropic elasticity without severe constraints. The elastic model is not sensitive to the underlying Euclidean geometry and is therefore blind to it.
- Once one has reconstructed the elastic geometry, one would like to interpret it physically and chemically. I believe it is beneficial to separate the two steps: first find the geometry, then interpret it.

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- 2 A mathematical result
  - A simplified model
  - Spectral rigidity of the round Earth
  - Limitations
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- A spherically symmetric Earth is  $M = \bar{B}(0, 1) \subset \mathbb{R}^3$  with  $g = c^{-2}e$  and  $c(x) = c(|x|)$ .
- If  $g$  is a rotation invariant Riemannian metric on  $M$ , there is a radial (more complicated if  $n = 2$ ) diffeomorphism  $\phi: M \rightarrow M$  so that  $\phi^*g$  is of this form.

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- The modes of free oscillations are modeled as scalar Neumann eigenfunctions of the Laplace–Beltrami operator of  $(M, g)$ .
- Reconstructing a Riemannian manifold  $(M, g)$  from the Neumann spectrum of  $\Delta_g$

$\approx$

Reconstructing elliptically inhomogeneous elastic object from the spectrum of free oscillations.

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## Definition

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Intuitively, this means that geodesics (straight lines) curve outwards.

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The Earth does not fully satisfy our assumptions:

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- In addition, the shear wave speed vanishes in the liquid outer core.

Apart from these problems (jumps and liquid) both shear and pressure wave speeds do satisfy the Herglotz condition everywhere.

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*Let  $M$  be the closed unit ball in  $\mathbb{R}^3$ . Let  $c_s(r)$  be a family of radial sound speeds depending  $C^\infty$ -smoothly on both  $s \in (-\varepsilon, \varepsilon)$  and  $r \in [0, 1]$ . Assume each  $c_s$  satisfies the Herglotz condition and a generic geometrical condition.*

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*If each  $c_s$  gives rise to the same spectrum (of the corresponding Laplace–Beltrami operator), then  $c_s = c_0$  for all  $s$ .*

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This simple model of the round Earth is spectrally rigid!

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- Our model only included P waves. A full linear elastic model should come with polarizations.
- The outer core is liquid, not solid.
- There are interfaces where the sound speed jumps. The real elastic geometry is non-smooth.
- The proof only works with strong assumptions, but the same auxiliary quantities might be useful to look at in any setting.

# Outline

- 1 The problem
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  - Data combination
  - Extracting useful quantities
  - Dealing with data blow-up
  - Deep teaching
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# Data combination

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- Mathematical models, solutions to inverse problems, and explicit reconstruction algorithms tend to be simple-minded: a single type of data is used to find a single quantity.
- All data can be used together in an algorithm based on machine learning or iteration. This method is rarely provably reliable.



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- This reliable reconstruction from partial data is useful *data* for a machine learning algorithm.
- Also less final data can be useful. We need not know how to use it.
- What are useful (and reliably computable) quantities?

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In a weak sense in spherical symmetry:  
The spectrum determines the geometry.



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- 1 From eigenvalues  $\lambda_0, \lambda_1, \lambda_2, \dots$  compute the function

$$f(t) = \sum_{k=0}^{\infty} \cos(\sqrt{\lambda_k} \cdot t).$$

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- 4 The periodic broken ray transform can be inverted explicitly.

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- The points where  $f$  behaves badly (singularities).
- The length spectrum (also accessible in other ways).

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If the length of the interval is  $L$ , the eigenvalues are

$$\lambda_k = \left( \frac{k\pi}{L} \right)^2, \quad k = 0, 1, 2, \dots$$

Suppose  $L = \frac{1}{2}$  and we have measured the numbers  $0, 4\pi^2, 16\pi^2, \dots$

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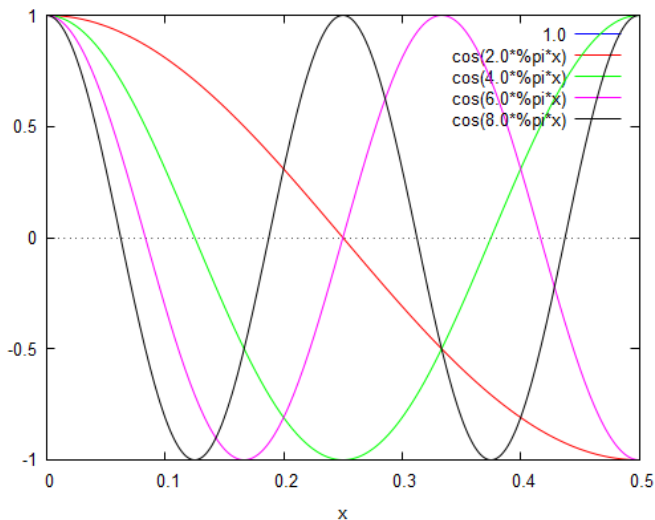
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We compute and plot the trace function

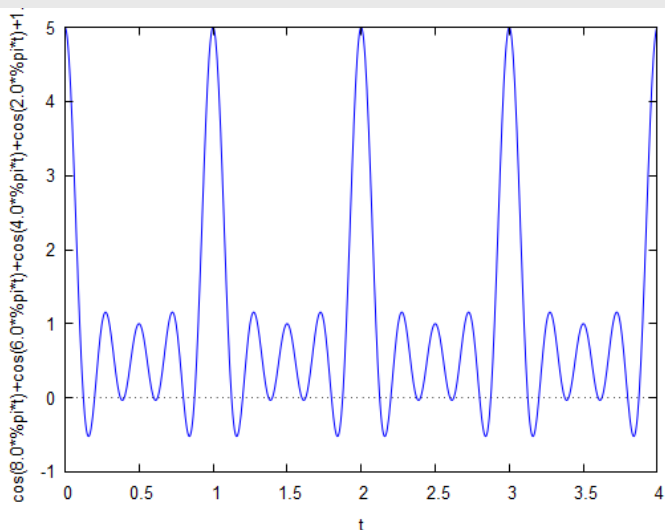
$$f(t) = \sum_{k=0}^{\infty} \cos(\sqrt{\lambda_k} \cdot t).$$

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Eigenfunctions for  $k = 0, 1, 2, 3, 4$ .

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Trace function computed from  $k = 0, 1, 2, 3, 4$ .

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- The singularities give the length spectrum.
- The length spectrum determines the geometry (length) of the interval.
- Perhaps the trace would be a useful quantity to extract for machine learning with spectral data. The singularities contain geometrical information, and we want to reconstruct the geometry.
- Perhaps other theorists have other potentially useful quantities to throw at a machine and improve performance.

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- The learning process can be unnecessarily (unfeasibly?) heavy.

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For example, we do not know whether the trace function is useful in more complicated spectral inverse problems.

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Examples:

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- 2 Deep teaching: If a machine learning algorithm learns what data is useful, it can tell it to those collecting data and working with theory.

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- 4 Appendix A: Anisotropy and geometry
  - Elliptic and general elastic anisotropy
  - Pressure and shear waves
  - Anisotropy and coordinates
  - Our model
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  - Riemannian manifolds are a very special subclass of Finsler manifolds.
- A material is isotropic if sound speed is independent of direction. This can be modeled by a conformally Euclidean metric.

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- To model elastic waves in general anisotropy, one needs a manifold with two Finsler metrics, one for pressure and one for shear waves.
- In fact, the shear wave speed might not even be a Finsler metric in the traditional sense.

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- A fully anisotropic model can *never* be reconstructed from boundary measurements uniquely. The data is always invariant under changes of coordinates.

# Anisotropy and coordinates

- Let  $\phi: M \rightarrow M$  be a diffeomorphism of a manifold that keeps the boundary fixed.
- If  $g$  (or  $F$ ) is a Riemannian (or Finsler) metric on  $M$ , then the pullback  $\phi^*g$  (or  $\phi^*F$ ) is different Riemannian metric that behaves exactly the same for boundary measurements.
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- A fully anisotropic model can *never* be reconstructed from boundary measurements uniquely. The data is always invariant under changes of coordinates.
- The best one can hope for is reconstruction up to changes of coordinates.
- The Earth is spherically symmetric to a good approximation, but the best (elliptically anisotropic) radial model might not be conformally Euclidean. After a radial change of coordinates the metric becomes conformal — and Cartesian coordinates are lost.

# Our model



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- Spherical symmetry.
- Reconstruction possible in the natural Cartesian coordinates. — No gauge freedom.

# Outline

- 1 The problem
- 2 A mathematical result
- 3 Deep learning deep structures
- 4 Appendix A: Anisotropy and geometry
- 5 Appendix B: Ray transforms in low regularity and drums
  - X-ray transforms
  - Periodic broken ray transforms
  - Hearing the shape of a drum
- 6 Appendix C: More detailed spectral rigidity

## Theorem (de Hoop–I.(2017))

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Earlier similar results:

- The X-ray transform (Radon et al.): Euclidean metric ( $c$  is constant).
- Mukhometov, 1977: Smooth simple metrics (simplicity is stronger than Herglotz).
- Sharafutdinov, 1997:  $C^\infty$  metrics and  $C^\infty$  functions.



# Periodic broken ray transforms

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*Let  $M$  be a rotation symmetric non-trapping manifold with a  $C^{1,1}$  metric and strictly convex boundary and dimension at least three. Assume that there are not too many conjugate points at the boundary. The integrals of a function  $f \in L^p(M)$ ,  $p > 3$ , over all periodic broken rays determines the even part of the function.*

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*Very little can be recovered of the odd part.*

Tools used:

- Planar average ray transform.
- Abel transform.
- Funk transform.
- Fourier series.

# Hearing the shape of a drum

# Hearing the shape of a drum

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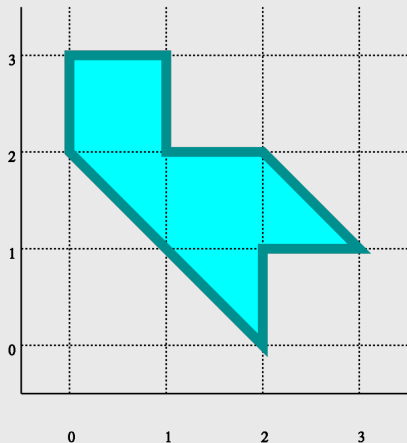
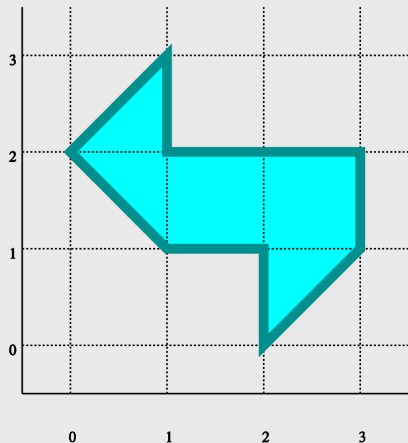
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# Hearing the shape of a drum

- If we know that the metric on  $M \subset \mathbb{R}^n$  is (conformally) Euclidean, this ambiguity due to diffeomorphisms goes away.
- This is why the answer to Kac's famous question "Can you hear the shape of a drum?" is not trivially "No!".
- ... but it is non-trivially "No!" if there are no geometrical restrictions.

# Hearing the shape of a drum



These two drums sound exactly alike. (Wikimedia Commons)



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# A simplified model

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- In practice, the Earth is the closed unit ball  $M = \bar{B}(0, 1) \subset \mathbb{R}^3$ . The anisotropic sound speed is modeled with a Riemannian metric  $g$  on  $M$ .
- Physically, this corresponds to omitting S-waves and including only elliptic anisotropy.

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- Let  $M = \bar{B}(0, 1) \subset \mathbb{R}^n$  be the closed unit ball and  $c(x) = c(|x|)$  a  $C^{1,1}$  sound speed.

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- If  $g$  is a rotation invariant Riemannian metric on  $M$ , there is a radial (more complicated if  $n = 2$ ) diffeomorphism  $\phi: M \rightarrow M$  so that  $\phi^*g$  is radially conformally Euclidean.



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- The spectrum of free oscillations is the Neumann spectrum of the Laplace–Beltrami operator  $\Delta_g$ .

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## Definition

A radial sound speed  $c(r)$  satisfies the Herglotz condition if

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- In addition, the shear wave speed vanishes in the liquid outer core.
- Apart from these problems (jumps and liquid) both shear and pressure wave speeds do satisfy the Herglotz condition everywhere.

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This simple model of the round Earth is spectrally rigid!



# Spectral rigidity of the round Earth

## Corollary (de Hoop–I.–Katsnelson (2017))

*Let  $M$  be the closed unit ball in  $\mathbb{R}^3$ . Let  $g_s$  be a family of rotation invariant metrics depending  $C^\infty$ -smoothly on  $s \in (-\varepsilon, \varepsilon)$ . Suppose each  $g_s$  is non-trapping with strictly convex boundary and assume a generic geometrical condition.*

*If the spectra of the Laplace–Beltrami operators  $\Delta_{g_s}$  are all equal, then there is a family of radial diffeomorphisms  $\phi_s: M \rightarrow M$  so that  $\phi_s^* g_s = g_0$  for all  $s$ . That is, the manifolds  $(M, g_s)$  are isometric.*

# Outline of the proof

## Lemma (Trace formula)

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Let  $\lambda_0 < \lambda_1 \leq \lambda_2 \leq \dots$  be the positive eigenvalues of the Laplace–Beltrami operator. Define a function  $f: \mathbb{R} \rightarrow \mathbb{R}$  by

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In particular, the spectrum determines the length spectrum.

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Similar “trace formulas” and related results are known on closed manifolds (eg. Duistermaat–Guillemin 1975) and a weaker version on some manifolds with boundary (eg. Guillemin–Melrose 1979).

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## Corollary

*Spectral rigidity follows if we can prove length spectral rigidity.*

Theorem (de Hoop–I.–Katsnelson (2017))



# Outline of the proof

## Theorem (de Hoop–I.–Katsnelson (2017))

*Let  $M$  be the closed unit ball in  $\mathbb{R}^n$ ,  $n \geq 2$ . Let  $c_s(r)$  be a family of radial sound speeds depending  $C^{1,1}$ -smoothly on both  $s \in (-\varepsilon, \varepsilon)$  and  $r \in [0, 1]$ . Assume each  $c_s$  satisfies the Herglotz condition and a generic geometrical condition.*

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In particular, if the length spectrum does not depend on  $s$ , then  $\frac{d}{ds} c_s^{-2}$  integrates to zero over (almost) all periodic broken rays.



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This concludes the proof.

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- No planet is precisely spherically symmetric.
- The proof only works with strong assumptions, but the same auxiliary quantities might be useful to look at in any setting.

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