

Spectral rigidity and tensor tomography

Jyväskylä analysis seminar

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Slides and papers are available at <http://users.jyu.fi/~jojapeil>

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Prelude

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- This is an inverse spectral problem. A hard one.
- There is a weaker version of the spectral problem: the spectral rigidity problem.
- Can we solve the simpler problem if we assume the Earth to be spherically symmetric?
- Yes!

Outline

- 1 Seismic spectral data
 - The spectrum of free oscillations
 - The spectrum of periodic orbits
 - The goal
- 2 Spectra of a manifold with boundary
- 3 Different forms of uniqueness
- 4 Spherical symmetry
- 5 The main results
- 6 Anisotropy and geometry
- 7 Non-linear and linear gauge freedom
- 8 Ray transforms

The spectrum of free oscillations

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- The set of these frequencies is the spectrum of free oscillations.

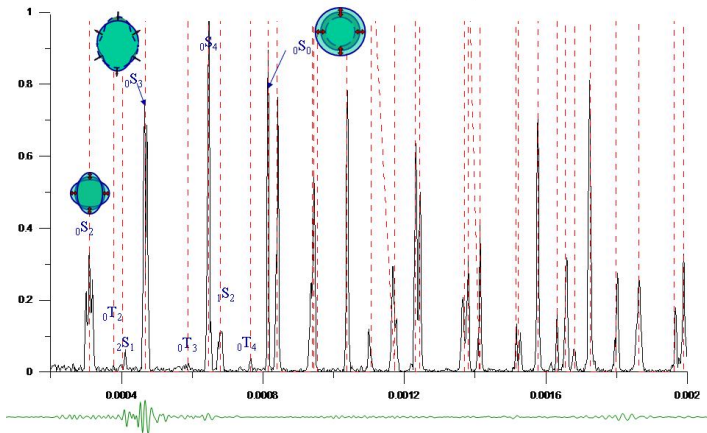
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- The amplitudes of different modes vary between different events, but the frequencies are always the same.
- The set of these frequencies is the spectrum of free oscillations.
- About 10 000 first frequencies are known.

The spectrum of free oscillations

Sumatra Earthquake: spectrum

Membach, SG C021, 20041226 08h00-20041231 00h00



Spectrum of free oscillations from an earthquake.

The spectrum of periodic orbits

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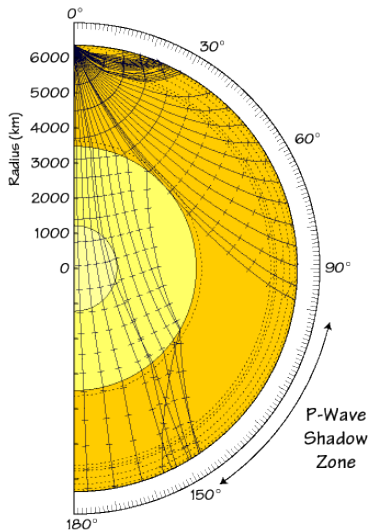
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- The set of all lengths of periodic seismic wave paths is the “length spectrum” of the Earth.

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- Some of the wave paths are periodic. Every periodic wave path has a length (in time).
- The set of all lengths of periodic seismic wave paths is the “length spectrum” of the Earth.
- Originally the length spectrum was just a mathematical tool, but it turns out it can be measured directly using deep earthquakes.

The spectrum of periodic orbits



Seismic wave paths and the P-wave shadow zone. (Wikimedia Commons)

The goal

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This problem only makes sense within a given model.

We want to reconstruct the Earth in the natural Cartesian coordinates.

- 1 Seismic spectral data
- 2 Spectra of a manifold with boundary
 - Manifolds with boundary
 - The spectrum of the Laplacian
 - The length spectrum
- 3 Different forms of uniqueness
- 4 Spherical symmetry
- 5 The main results
- 6 Anisotropy and geometry
- 7 Non-linear and linear gauge freedom
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Manifolds with boundary

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- In practice, the Earth is the closed unit ball $M = \bar{B}(0, 1) \subset \mathbb{R}^3$. The anisotropic sound speed is modeled with a Riemannian metric g on M .
- Physically, this corresponds to omitting S-waves and including only elliptic anisotropy.

The spectrum of the Laplacian

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$$\Delta_g u(x) = c(x)^n \operatorname{div}(c(x)^{2-n} \nabla u(x)).$$

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$$\Delta_g u(x) = c(x)^n \operatorname{div}(c(x)^{2-n} \nabla u(x)).$$

- The spectrum of free oscillations is the Neumann spectrum of the Laplace–Beltrami operator Δ_g .

The length spectrum

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- Seismic waves reflect at the surface, so they are in fact billiard trajectories or broken rays.
- The length spectrum of (M, g) is the set of all lengths of the periodic broken rays.

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- 1 Seismic spectral data
- 2 Spectra of a manifold with boundary
- 3 Different forms of uniqueness
 - Difficulties
 - Diffeomorphisms and coordinates
 - Global uniqueness
 - Local uniqueness
 - Spectral rigidity
- 4 Spherical symmetry
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Difficulties

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- Proving this conjecture is difficult for two reasons:
 - 1 The required tools do not yet exist on general manifolds with boundary.
 - 2 The conjecture is false.

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- If $\phi: M \rightarrow M$ is a diffeomorphism, then (M, g) and (M, ϕ^*g) give the same spectrum.
- We can take any change of coordinates whatsoever and use it to distort the metric, but the spectrum stays the same.
- Physically: There are preferred and natural Cartesian coordinates. But the anisotropic model is not “sensitive to the underlying Euclidean geometry”, so the Cartesian coordinates cannot be recognized. It is impossible to find the metric (anisotropic sound speed) in Cartesian coordinates from spectral data.

Problem

Let g_1 and g_2 be two Riemannian metrics on a manifold M with boundary. If they give the same spectrum, is there a diffeomorphism $\phi: M \rightarrow M$ so that $g_1 = \phi^ g_2$?*

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Problem

Let g_1 and g_2 be two Riemannian metrics on a manifold M with boundary. Suppose g_1 is very close to g_2 . If they give the same spectrum, is there a diffeomorphism $\phi: M \rightarrow M$ so that $g_1 = \phi^ g_2$?*

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This is still too hard.

Problem

Let g_s be family of Riemannian metrics on a manifold M with boundary, depending on a parameter $s \in (-\varepsilon, \varepsilon)$. If they all give the same spectrum, are there a diffeomorphisms $\phi_s: M \rightarrow M$ so that $g_0 = \phi_s^ g_s$?*

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This is within reach!

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 - Negatively curved surfaces: Guillemin–Kazhdan 1980.
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- We have adapted similar ideas of proof to manifolds with boundary.

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 - Spherically symmetric manifolds
 - The Herglotz condition
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Spherically symmetric manifolds

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- If g is a rotation invariant Riemannian metric on M , there is a radial (more complicated if $n = 2$) diffeomorphism $\phi: M \rightarrow M$ so that ϕ^*g is radially conformally Euclidean.
- The Earth is spherically symmetric to a good approximation, but the best (elliptically anisotropic) radial model might not be conformally Euclidean. After a radial change of coordinates the metric becomes conformal — and Cartesian coordinates are lost.

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Definition

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- All spheres $\{r = \text{constant}\}$ are strictly convex. (Foliation condition!)
- The manifold is non-trapping and has strictly convex boundary.

The Herglotz condition

- The radial Preliminary Reference Earth Model (PREM) is not $C^{1,1}$. Both pressure and shear wave speeds have jump discontinuities.

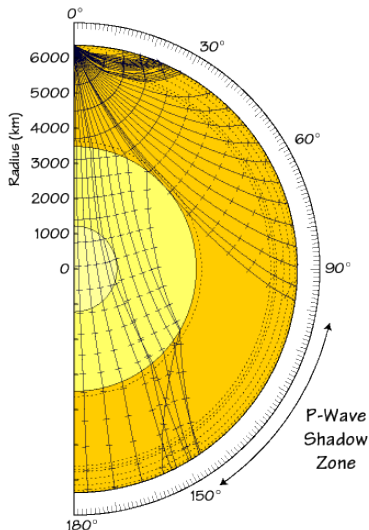
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- In addition, the shear wave speed vanishes in the liquid outer core.
- Apart from these problems (jumps and liquid) both shear and pressure wave speeds do satisfy the Herglotz condition everywhere.

The Herglotz condition



The Herglotz condition is satisfied: ray paths curve outwards. (Wikimedia Commons)

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 - Length spectral rigidity
 - Ideas behind the proof
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Theorem (de Hoop–I.–Katsnelson, 2017)

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Let M be the closed unit ball in \mathbb{R}^3 . Let $c_s(r)$ be a family of radial sound speeds depending C^∞ -smoothly on both $s \in (-\varepsilon, \varepsilon)$ and $r \in [0, 1]$. Assume each c_s satisfies the Herglotz condition and a generic geometrical condition.

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This simple model of the round Earth is spectrally rigid!

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Let M be the closed unit ball in \mathbb{R}^3 . Let g_s be a family of rotation invariant metrics depending C^∞ -smoothly on $s \in (-\varepsilon, \varepsilon)$. Suppose each g_s is non-trapping with strictly convex boundary and assume a generic geometrical condition.

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If the spectra of the Laplace–Beltrami operators Δ_{g_s} are all equal, then there is a family of radial diffeomorphisms $\phi_s: M \rightarrow M$ so that $\phi_s^ g_s = g_0$ for all s . That is, the manifolds (M, g_s) are isometric.*

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Let M be the closed unit ball in \mathbb{R}^n , $n \geq 2$. Let $c_s(r)$ be a family of radial sound speeds depending $C^{1,1}$ -smoothly on both $s \in (-\varepsilon, \varepsilon)$ and $r \in [0, 1]$. Assume each c_s satisfies the Herglotz condition and a generic geometrical condition.

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If the length spectra of the manifolds (M, g_s) are all equal, then there is a family of radial (or more general if $n = 2$) diffeomorphisms $\phi_s: M \rightarrow M$ so that $\phi_s^ g_s = g_0$ for all s . That is, the manifolds (M, g_s) are isometric.*

Ideas behind the proof

Lemma (Trace formula)

Let $\lambda_0 < \lambda_1 \leq \lambda_2 \leq \dots$ be the positive eigenvalues of the Laplace–Beltrami operator. Define a function $f: \mathbb{R} \rightarrow \mathbb{R}$ by

$$f(t) = \sum_{k=0}^{\infty} \cos \left(\sqrt{\lambda_k} \cdot t \right).$$

Assume that the radial sound speed c satisfies some generic condition.

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In particular, the spectrum determines the length spectrum.

Similar “trace formulas” and related results are known on closed manifolds (eg. Duistermaat–Guillemin 1975) and a weaker version on some manifolds with boundary (eg. Guillemin–Melrose 1979).

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Corollary

Spectral rigidity follows from length spectral rigidity.

Lemma

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$$\frac{d}{ds}\ell_s(\gamma_s) = \frac{1}{2} \int_{\gamma_s} \frac{d}{ds}c_s^{-2}.$$

In particular, if the length spectrum does not depend on s , then $\frac{d}{ds}c_s^{-2}$ integrates to zero over (almost) all periodic broken rays.

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Assume the Herglotz condition. A radially symmetric function is uniquely determined by its integrals over (almost) all periodic broken rays.

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This concludes the proof.

Remark: No proof works without spherical symmetry.

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- 5 The main results
- 6 Anisotropy and geometry
 - Elliptic and general elastic anisotropy
 - Pressure and shear waves
 - Anisotropy and coordinates
 - Our model
- 7 Non-linear and linear gauge freedom
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Elliptic and general elastic anisotropy

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 - Riemannian manifolds are a very special subclass of Finsler manifolds.
- A material is isotropic if sound speed is independent of direction. This can be modeled by a conformally Euclidean metric.

Pressure and shear waves

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- To model elastic waves in general anisotropy, one needs a manifold with two Finsler metrics, one for pressure and one for shear waves.
- In fact, the shear wave speed might not even be a Finsler metric in the traditional sense.

Anisotropy and coordinates

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- If g (or F) is a Riemannian (or Finsler) metric on M , then the pullback ϕ^*g (or ϕ^*F) is different Riemannian (Finsler) metric that behaves exactly the same for any boundary or spectral measurements.
- A fully anisotropic model can *never* be reconstructed from boundary measurements uniquely. The data is always invariant under changes of coordinates.

- Let $\phi: M \rightarrow M$ be a diffeomorphism of a manifold that keeps the boundary fixed.
- If g (or F) is a Riemannian (or Finsler) metric on M , then the pullback ϕ^*g (or ϕ^*F) is different Riemannian (Finsler) metric that behaves exactly the same for any boundary or spectral measurements.
- A fully anisotropic model can *never* be reconstructed from boundary measurements uniquely. The data is always invariant under changes of coordinates.
- The best one can hope for is reconstruction up to changes of coordinates.

Our model

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- Isotropic P-wave speed. — Conformally Euclidean metric.
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- Reconstruction possible in the natural Cartesian coordinates. — No gauge freedom.

Outline

- 1 Seismic spectral data
- 2 Spectra of a manifold with boundary
- 3 Different forms of uniqueness
- 4 Spherical symmetry
- 5 The main results
- 6 Anisotropy and geometry
- 7 Non-linear and linear gauge freedom
 - Non-linear gauge freedom
 - Linear gauge freedom
 - Interplay
 - Symmetry and conformal variations
- 8 Ray transforms

Non-linear gauge freedom

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- Correct question: Are there diffeomorphisms $\phi_s: M \rightarrow M$ depending smoothly on s so that $g_0 = \phi_s^* g_s$?
- The freedom of diffeomorphisms is a gauge freedom. The metric can only ever be recovered up to gauge.

Linear gauge freedom

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- Correct question: If a symmetric two-tensor field f integrates to zero over every closed geodesic, is there a one-form h so that $f = \sigma \nabla h$?
- The freedom of having a potential h is a gauge freedom. The tensor field can only ever be reconstructed up to gauge.

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- Each ϕ_s is a diffeomorphism and $\phi_s^* g_s = g_0$.

Symmetry and conformal variations

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- This leads to the usual X-ray tomography problem (on a closed manifold): Is a scalar function uniquely determined by its integrals over all closed geodesics?
- There is no gauge freedom!
- Every rotation symmetric metric can be made conformally Euclidean.

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 - Tensor tomography
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 - Periodic broken ray transforms
 - Periodic slabs

Tensor tomography

Theorem (Paternain–Salo–Uhlmann, 2014–2015)

Let (M, g) be a closed manifold of dimension $n \geq 2$ whose geodesic flow is Anosov. If $n \geq 3$, assume also that the terminator value is $> \frac{2(n+1)}{n+2}$.

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A smooth symmetric second rank tensor field f integrates to zero over all closed geodesics if and only if there is a one-form h so that $f = \sigma \nabla h$.

X-ray transforms

Theorem (de Hoop–I., 2017)

Let M be a rotation symmetric non-trapping manifold with a piecewise $C^{1,1}$ metric and strictly convex boundary. Then the geodesic X-ray transform is injective on $L^2(M)$.

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Earlier similar results:

- The X-ray transform (Radon et al.): Euclidean metric (c is constant).
- Mukhometov, 1977: Smooth simple metrics (simplicity is stronger than Herglotz).
- Sharafutdinov, 1997: C^∞ metrics and C^∞ functions.

Periodic broken ray transforms

Theorem (de Hoop–I., 2017)

Let M be a rotation symmetric non-trapping manifold with a $C^{1,1}$ metric and strictly convex boundary and dimension at least three. Assume that there are not too many conjugate points at the boundary. The integrals of a function $f \in L^p(M)$, $p > 3$, over all periodic broken rays determine the even part of the function.

Very little can be recovered of the odd part.

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Very little can be recovered of the odd part.

Tools used:

- Planar average ray transform.
- Abel transform.
- Funk transform.
- Fourier series.

Periodic slabs

Theorem (I.-Uhlmann, 2017)

A scalar L^2 function on $[0, 1] \times \mathbb{T}^n$ integrates to zero over all lines if and only if it is of the form $f(x, y) = h(x)$ with $\int_0^1 h(x) dx = 0$.

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There is a kernel for scalar functions!

We can also characterize the kernel of the ray transform for tensor fields: it contains tensors of the above form plus a potential. Symmetry increases gauge freedom.

End

Thank you.

Slides are available at <http://users.jyu.fi/~jojapeil>.