



JYVÄSKYLÄN YLIOPISTO  
UNIVERSITY OF JYVÄSKYLÄ

# Dix's inverse problem on elastic Finsler manifolds

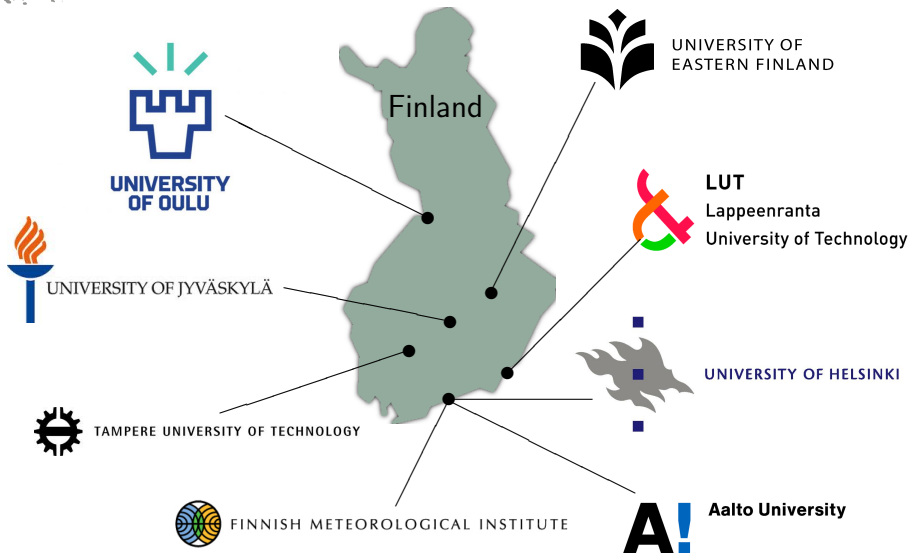
Geo-Mathematical Imaging Group  
2019 project review

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Based on joint work with  
**Maarten de Hoop and Matti Lassas**

# Finnish Centre of Excellence in Inverse Modelling and Imaging 2018-2025



# Goals

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- Understanding how and why to model elastic waves in terms of Finsler geometry.

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- Understanding how and why to model elastic waves in terms of Finsler geometry.
- Seeing how geometrical tools can be used to solve Dix's problem in great generality.

- 1 Elastic geometry
- 2 Dix's inverse problem

# Elastic waves

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In linear elasticity the displacement vector  $u(x, t)$  satisfies the elastic wave equation

$$\partial_j [c_{ijkl}(x) \partial_k u_l(x, t)] - \rho(x) \partial_t^2 u_i(x, t) = 0,$$

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The singularities can be described in geometrical terms as point particles or wave fronts.

# Propagation of singularities

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The slowness vector  $p$  (inverse phase velocity), position  $x$ , and polarization  $A$  of the singularity have to satisfy

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The possible values of  $p$  are on the slowness surface defined by the equation

$$\det[\Gamma(x, p) - I] = 0.$$

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Microlocal analysis (or physical reasoning) tells how singularities propagate: They follow a Hamiltonian flow related to the slowness surface.

If the slowness surface is smooth and strictly convex ( $qP!$ ), then this coincides with the geodesic flow of a Finsler manifold, where the Finsler metric is the Legendre transform of the slowness surface.

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If the slowness surfaces are ellipsoids, we have a Riemannian manifold. But we do not assume any kind of isotropy!

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- The slowness surface corresponds to a co-Finsler metric.
- Taking the Legendre transform gives a Finsler metric which corresponds to group velocities.
- The singularities of elastic waves (= elastic waves as point particles) follow geodesics of this Finsler geometry.



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Gains:

- 1 Freedom to have any shape of slowness surfaces, not necessarily arising from Hookean elasticity.
- 2 Access to powerful tools of differential geometry.

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One can define (and write) a new Finsler metric  $\phi^*F$ , the pullback of  $F$  over  $\phi$ . All geodesics and everything behaves exactly the same —  $F$  and  $\phi^*F$  look exactly alike from the surface.

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Not all Finsler metrics arise from a stiffness tensor. When does  $\phi^* F$  do?

# Outline

- 1 Elastic geometry
- 2 Dix's inverse problem

# The inverse problem for a geometer

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The Riemannian version was solved by de Hoop, Holman, Iversen, Lassas, and Ursin (2015). Moving to Finsler geometry means that we can deal with arbitrary (non-elliptic) anisotropy.

We make no assumptions on symmetry, isotropy, or homogeneity.

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## Theorem (de Hoop–I.–Lassas)

*Let  $(M, F)$  be any Finsler manifold and  $U \subset M$  an open set. Suppose  $F$  is fiberwise analytic. If we know the smooth metric “spheres” (wave fronts from point sources) together with radii (travel times) in the set  $U$ , we can find the universal cover of  $(M, F)$ .*

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All Finsler metrics arising from elasticity are fiberwise analytic. This is a good example of useful additional structure that comes from good modelling — the model is general enough to handle all anisotropy but narrow enough to exclude nonphysical oddities.

# Ideas and methods



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- We can find the geometry along any geodesic, and this information must be turned into a global geometrical description. This requires using analyticity and building an atlas (a collection of coordinate charts).

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