

1. As on p. 247-8 in the lecture notes, let's study the group $SU(N)$.

- (a) Knowing that the $SU(N)$ group matrices can be written in the form

$$U(\alpha) = \exp \left[i \sum_{i=1}^d \alpha_i T_i \right],$$

where T_i are the group generator matrices and α_i are real parameters, convince yourself that the generators have to be hermitian and traceless. You can use the facts $\det(e^A) = e^{\text{Tr}(A)}$ and $e^A e^B = e^{A+B+\frac{1}{2}[A,B]+\dots}$. The purpose of this problems is that you understand why the color symmetry group $SU(3)$ generators (The Gell-Mann -matrices) satisfy these conditions.

- (b) The generators are linearly independent matrices and fulfil $\text{Tr}(T_i T_j) = \lambda \delta_{ij}$. Using the Lie algebra for this group and the properties of the trace, show that the structure constants can be expressed as $C_{ljk} = -\frac{i}{\lambda} \text{Tr}(T_l [T_j, T_k])$ and that C_{ljk} are fully antisymmetric in any mutual exchange of indices.

2. As suggested on p.249 in the lecture notes,

- (a) Generalize the Noether theorem for the globally $U(1)$ symmetric Lagrangian

$$\mathcal{L}(\phi, \phi^*, \partial_\mu \phi, \partial_\mu \phi^*) = \partial_\mu \phi^* \partial^\mu \phi - m^2 \phi^* \phi - V_I(\phi^* \phi),$$

describing a charged scalar field (see p. 249), i.e. find the form of the conserved 4-current, analogous to that on p. 236.

- (b) Using the $U(1)$ transformation $U = e^{i\theta}$ (expand), find the variations $\delta\phi$ and $\delta\phi^*$ and based on your result above, show that you arrive at the same conserved 4-current as given on p. 249.

3. Using the general properties of the Dirac matrices $\alpha^1, \alpha^2, \alpha^3$ and β (see p. 255 in the lecture notes), and the definitions of the Dirac gamma matrices γ^μ (p. 256), verify the following, representation-independent, results:

- (a) $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu} \mathbf{1}_4$, known as the Clifford algebra,
- (b) $\gamma^i = -\gamma^0 \gamma^i \gamma^0$,
- (c) $\gamma^{0\dagger} = \gamma^0$ ja $\gamma^{i\dagger} = \gamma^0 \gamma^i \gamma^0 = -\gamma^i$,
- (d) $\gamma^{\mu\dagger} = \gamma^0 \gamma^\mu \gamma^0$.

4. Using the explicit 2-block forms for the free spin- $\frac{1}{2}$ particle & antiparticle spinors $u^{(r)}(p)$ and $v^{(r)}(p)$ ($r, s = 1, 2$), which we derived in the lecture notes (p.259-263), show that with the chosen normalization $\int_V d^3x \rho = 2E_p$ (p. 264), we have the following

- (a) normalization constant: $N(\mathbf{p}) = \sqrt{E_p + m}$
 (Hints: Do the spinor and matrix multiplications using the 2-block forms. Recall the Pauli spin-matrix property $(\vec{\sigma} \cdot \vec{p})^2 = \vec{p}^2$)
- (b) orthogonality relations: $u^{(r)\dagger}(p)u^{(s)}(p) = v^{(r)\dagger}(p)v^{(s)}(p) = 2E_p\delta_{rs}$
- (c) projection operators:
 (Hint: do the multiplication by γ^0 only after performing the spin sums)

$$\sum_{s=1,2} u^{(s)}(p)\bar{u}^{(s)}(p) = \not{p} + m$$

$$\sum_{s=1,2} v^{(s)}(p)\bar{v}^{(s)}(p) = \not{p} - m$$

5. Let's practice drawing Feynman graphs for scatterings in a given theory

- (a) Let's first consider a ϕ^3 -theory for an interacting neutral spin-0 particle, with a Lagrangian

$$\mathcal{L} = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{3!}\phi^3.$$

Draw all possible topologically different Feynman diagrams for $2 \rightarrow 2$ scattering in this theory, first in the lowest order (LO) in λ , then in the next-to-lowest order (NLO). Note that in the latter case there are quite many diagrams, be systematic in sorting these out.

- (b) Let's then consider some QED scatterings. Draw all Feynman graphs which contribute to the invariant amplitude in the lowest order in e (figure out which power of e corresponds to the lowest order in each case!) for the following scatterings

i $e^+ + e^- \rightarrow e^+ + e^-$

ii $e^+ + e^- \rightarrow \mu^+ + \mu^-$

iii $e^+ + e^- \rightarrow e^+ + e^- + \mu^+ + \mu^-$

Label the lines and draw the required arrows.