FYST530 Quantum Mechanics II

Problem set 12 Return before 15 on Wednesday 18.4. to the box labeled FYST530

1. (a) Show that the Fock space 1-particle operator, $F^{(1)} = \sum_{\mu,\nu} \langle \mu | \hat{f} | \nu \rangle a^{\dagger}_{\mu} a_{\nu}$, acts as

$$F^{(1)}|1_{\nu_1}1_{\nu_2}\cdots 1_{\nu_N}\rangle = \sum_{\mu}\sum_{i=1}^N \langle \mu|\hat{f}|\nu_i\rangle|1_{\nu_1}1_{\nu_2}\cdots 0_{\nu_i}1_{\mu}\cdots 1_{\nu_N}\rangle$$

in a system of N fermions.

(b) Taking the ansatz

$$F^{(2)} = K \sum_{\mu\mu'\nu\nu'} \langle \mu\mu' | \hat{g} | \nu\nu' \rangle a^{\dagger}_{\mu} a^{\dagger}_{\mu'} a_{\nu} a_{\nu'}$$

for the Fock space 2-particle operator, and requiring that

$$F^{(2)}|1_{\nu_1}1_{\nu_2}\rangle = \sum_{\mu_1\mu_2} \langle \mu_1\mu_2 |\hat{g}|\nu_1\nu_2\rangle |1_{\mu_1}1_{\mu_2}\rangle,$$

show that $K = -\frac{1}{2}$.

2. The fermionic pair distribution function $g_{\sigma\sigma'}(\mathbf{x} - \mathbf{x}')$ for the ground state $|F\rangle$ of the Fermi gas is defined as

$$\left(\frac{n}{2}\right)^2 g_{\sigma\sigma'}(\mathbf{x} - \mathbf{x}') \equiv \langle F | \Psi^{\dagger}(\mathbf{x}, \sigma) \Psi^{\dagger}(\mathbf{x}', \sigma') \Psi(\mathbf{x}', \sigma') \Psi(\mathbf{x}, \sigma) | F \rangle,$$

where $\Psi(\mathbf{x}, \sigma)$ s are fermionic field operators. Show that

(a) If $\sigma \neq \sigma'$, then

$$g_{\sigma\sigma'}(\mathbf{x}-\mathbf{x}')=1.$$

(b) If $\sigma = \sigma'$, then

$$g_{\sigma\sigma'}(\mathbf{x}-\mathbf{x}')=1-\left[G_{\sigma}(\mathbf{x}-\mathbf{x}')\right]^2,$$

where $G_{\sigma}(\mathbf{x} - \mathbf{x}')$ is the single-particle correlator

$$G_{\sigma}(\mathbf{x} - \mathbf{x}') \equiv \frac{2}{n} \langle F | \Psi^{\dagger}(\mathbf{x}, \sigma) \Psi(\mathbf{x}', \sigma) | F \rangle.$$

3. Consider a system of N identical <u>bosons</u>, spin-0, enclosed to a box with volume V, in the state

$$|\phi\rangle = |n_{\mathbf{p}_0} n_{\mathbf{p}_1} n_{\mathbf{p}_2}, ...\rangle \qquad n_{\mathbf{p}_i} = 0, 1, 2, 3, ...$$

(a) Show that

$$\langle \phi | \Psi^{\dagger}(\mathbf{x}) \Psi(\mathbf{x}) | \phi \rangle = \frac{N}{V} \equiv n.$$

(b) Show that the pair distribution function is $({f k}\equiv{f p}/\hbar)$

$$\begin{split} n^2 g(\mathbf{x} - \mathbf{x}') &\equiv \langle \phi | \Psi^{\dagger}(\mathbf{x}) \Psi^{\dagger}(\mathbf{x}') \Psi(\mathbf{x}') \Psi(\mathbf{x}) | \phi \rangle \\ &= n^2 + \left| \frac{1}{V} \sum_k e^{-i\mathbf{k} \cdot (\mathbf{x} - \mathbf{x}')} n_{\mathbf{k}} \right|^2 - \frac{1}{V^2} \sum_k n_{\mathbf{k}} (n_{\mathbf{k}} + 1). \end{split}$$

- (c) Consider two examples of $g(\mathbf{x} \mathbf{x}')$:
 - (i) Suppose all the bosons occupy the same state, $n_{\mathbf{p}_0}=N.$ Show that

$$n^2 g(\mathbf{x} - \mathbf{x}') = \frac{N(N-1)}{V^2}.$$

(ii) Suppose the particles are distributed over many momentum values with Gaussian distribution

$$n_{\mathbf{p}} = \frac{(2\pi)^3 n}{(\sqrt{\pi}\Delta)^3} e^{-(\mathbf{k} - \mathbf{k}_0)^2 / \Delta^2}$$

Show that if the density of particles n, and the width Δ of the momentum distribution are held fixed, then, in the limit of large volume V, the pair distribution function is given by

$$n^2 g(\mathbf{x} - \mathbf{x}') = n^2 \left(1 + e^{-\frac{\Delta^2}{2}(\mathbf{x} - \mathbf{x}')^2} \right)$$

Sketch the behaviour of $g(\mathbf{x} - \mathbf{x}')$ as a function of $(\mathbf{x} - \mathbf{x}')$. Can you see the bosonic nature of the particles showing up in $g(\mathbf{x} - \mathbf{x}')$?

4. Suppose the Hamilton operator for system of N fermions (in a box with volume V) is of the form

$$H = \sum_{i=1}^{N} \frac{\hat{\mathbf{p}}_i^2}{2m} + \sum_{i < j} V_2(\mathbf{x}_i - \mathbf{x}_j),$$

which is invariant in translations. Show that in this case the plane-waves

$$\varphi_{\mathbf{k},s}(\mathbf{x},\sigma) = \frac{1}{\sqrt{V}} e^{i\mathbf{x}\cdot\mathbf{k}} \delta_{\sigma s}$$

are solutions for the Hartree-Fock equations. Compute the single-particle energies in terms of Fouries components $V_{\bf q}$ of V_2 :

$$V_2(\mathbf{x}) = \sum_q e^{i\mathbf{q}\cdot\mathbf{k}} V_{\mathbf{q}}$$

Compute also the total energy in terms of spins and momenta of the occupied states.