

$$1. K(L, R, U) = \sqrt{R U L^+}^{-1} R \sqrt{U}$$

$$a) K(L_1, R_1, R_2 U L_2^+) K(L_2, R_2, U)$$

$$= \sqrt{R_1 R_2 U L_2^+ L_1^+}^{-1} R_1 \sqrt{R_2 U L_2^+} \sqrt{R_2 U L_2^+}^{-1} R_2 \sqrt{U}$$

$$= \sqrt{R_1 R_2 U (L_1 L_2)^+}^{-1} R_1 R_2 \sqrt{U}$$

$$= K(L_1 L_2, R_1 R_2, U)$$

b) isospm transform $R = L = V$

$$K(V, V, U) = \sqrt{V U V^+}^{-1} V \sqrt{U}$$

observe:

$$\left(\sqrt{V U V^+}\right)^2 = V U V^+ = V \sqrt{U} V^+ V \sqrt{U} V^+$$

$$\Rightarrow \sqrt{V U V^+} = V \sqrt{U} V^+$$

$$\Rightarrow K(V, V, U) = (V \sqrt{U} V^+)^{-1} V \sqrt{U}$$

$$= (V^+)^{-1} \sqrt{U}^{-1} V^{-1} V \sqrt{U} = V \sqrt{U}^{-1} \sqrt{U} = V$$

Thus $\Psi \rightarrow \Psi' = K(V, V, U) \Psi = V \Psi$ independent of U .

$$2. \quad \mathcal{L}_{\pi_N} = \bar{\Psi} \left(i \not{\partial} - m_N + \frac{g_A}{2} \gamma^\mu \gamma_5 u_\mu \right) \Psi$$

$$D_\mu = \partial_\mu + \Gamma_\mu$$

$$\Gamma_\mu = \frac{1}{2} \left(u^\dagger \partial_\mu u + u \partial_\mu u^\dagger \right)$$

$$a) \quad u = \exp \left[i \frac{\bar{\tau} \cdot \bar{\phi}}{2F} \right] = 1 + i \frac{\bar{\tau} \cdot \bar{\phi}}{2F} - \frac{\phi^2}{8F^2}, \quad \phi = \bar{\tau} \cdot \bar{\phi}$$

$$a) \quad (\bar{\tau} \cdot \bar{\phi})(\tau \cdot \phi) = \phi^2 (= \bar{\phi} \cdot \bar{\phi}) \quad \text{and} \quad \tau_i^\dagger = \tau_i$$

we get

$$\begin{aligned} 2 \Gamma_\mu &= \left(1 - i \frac{\bar{\tau} \cdot \bar{\phi}}{2F} - \frac{\phi^2}{8F^2} \right) \left(\frac{i}{2F} \bar{\tau} \cdot \partial_\mu \bar{\phi} - \frac{\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi}{8F^2} \right) \\ &+ \left(1 + i \frac{\bar{\tau} \cdot \bar{\phi}}{2F} - \frac{\phi^2}{8F^2} \right) \left(-\frac{i}{2F} \bar{\tau} \cdot \partial_\mu \bar{\phi} - \frac{\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi}{8F^2} \right) \\ &= \cancel{\frac{i}{2F} \bar{\tau} \cdot \partial_\mu \bar{\phi}} + \frac{(\bar{\tau} \cdot \bar{\phi})(\bar{\tau} \cdot \partial_\mu \bar{\phi})}{4F^2} - \frac{\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi}{8F^2} \\ &\quad - \cancel{\frac{i}{2F} \bar{\tau} \cdot \partial_\mu \bar{\phi}} + \frac{(\bar{\tau} \cdot \bar{\phi})(\bar{\tau} \cdot \partial_\mu \bar{\phi})}{4F^2} - \frac{\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi}{8F^2} + o(\phi^3) \end{aligned}$$

$$\Gamma_\mu = \frac{(\bar{\tau} \cdot \bar{\phi})(\bar{\tau} \cdot \partial_\mu \bar{\phi})}{4F^2} - \frac{1}{8F^2} (\partial_\mu \phi \cdot \phi + \bar{\phi} \cdot \partial_\mu \phi)$$

Similarly we get $u_\mu = i \left[u^\dagger \partial_\mu u - u \partial_\mu u^\dagger \right] = - \frac{\tau \cdot \partial_\mu \bar{\phi}}{F}$

Plus these into the Lagrangian and get the interaction terms

$$\mathcal{L}_{\text{int}} = \bar{\Psi} \left[i \partial_\mu + \frac{i}{4F^2} (\bar{\tau} \cdot \bar{\phi}) (\bar{\tau} \cdot \partial_\mu \bar{\phi}) - \frac{i}{8F^2} (\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi) - m_N + \frac{g_A}{2} \gamma^\mu \gamma_5 \left(- \frac{\bar{\tau} \cdot \partial_\mu \bar{\phi}}{F} \right) \right]$$

$$= \bar{\Psi} (i \partial_\mu - m_N) \Psi - \frac{i}{2} \cdot \frac{g_A}{F} \gamma^\mu \gamma_5 \bar{\tau} \cdot \partial_\mu \phi \Psi + \frac{i}{4F^2} \bar{\Psi} \gamma^\mu \left[(\bar{\tau} \cdot \bar{\phi}) (\bar{\tau} \cdot \partial_\mu \bar{\phi}) - \frac{\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi}{2} \right]$$

(typo in the problem text, missing i and wrong sign)

$$b) \quad (\bar{\tau} \cdot \bar{A}) (\bar{\tau} \cdot \bar{B}) = \bar{A} \cdot \bar{B} \mathbb{1} + i (\bar{A} \times \bar{B}) \cdot \bar{\tau}$$

$$\text{Now } (\bar{\tau} \cdot \bar{\phi}) (\bar{\tau} \cdot \partial_\mu \bar{\phi}) = \bar{\phi} \cdot \partial_\mu \bar{\phi} + i \epsilon_{abc} T^a \bar{\phi}^b (\partial_\mu \bar{\phi})^c$$

and

$$\frac{i}{2} (\partial_\mu \phi \cdot \phi + \phi \cdot \partial_\mu \phi) = \frac{i}{2} \left[(\bar{\tau} \cdot \partial_\mu \bar{\phi}) (\bar{\tau} \cdot \bar{\phi}) + (\bar{\tau} \cdot \bar{\phi}) (\bar{\tau} \cdot \partial_\mu \bar{\phi}) \right]$$

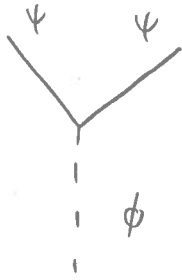
$$= \frac{i}{2} \left[\underbrace{\partial_\mu \bar{\phi} \cdot \bar{\phi} + \bar{\phi} \cdot \partial_\mu \bar{\phi}}_{2 \bar{\phi} \cdot \partial_\mu \bar{\phi}} + i (\partial_\mu \bar{\phi} \times \bar{\phi}) + i (\bar{\phi} \times \partial_\mu \bar{\phi}) \right]$$

$$= \bar{\phi} \cdot \partial_\mu \bar{\phi}$$

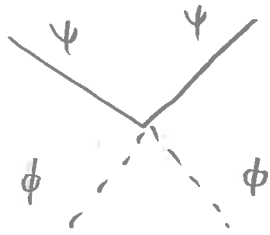
Thus the 2nd interaction term reads

$$\frac{i}{4F^2} \bar{\Psi} \gamma^\mu \left[i \epsilon_{abc} T^a \bar{\phi}^b (\partial_\mu \bar{\phi})^c \right] \Psi = \underline{\underline{-\frac{1}{4F^2} \bar{\Psi} \gamma^\mu \epsilon_{abc} T^a \bar{\phi}^b (\partial_\mu \bar{\phi})^c \Psi}}$$

$$c. \quad -\frac{1}{2} \frac{g_A}{F} \bar{\Psi} \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\phi} \Psi$$



$$-\frac{1}{4F^2} \bar{\Psi} \gamma^\mu \epsilon_{abc} \vec{\tau}^b \vec{\phi}^c (\partial_\mu \vec{\phi})^a \Psi$$



3. Hadronic states $\langle B | \dots | A \rangle$

Momentum eigenvalues p_A^μ p_B^μ

transl. inv.

a)

$$\langle B | \partial_\mu A_i^\mu(x) | A \rangle \stackrel{\text{transl. inv.}}{=} \partial_\mu \underbrace{\langle B | e^{i\hat{p}\cdot x} A_i^\mu(0) e^{-i\hat{p}\cdot x} | A \rangle}_{(e^{-i\hat{p}\cdot x} | B \rangle)^\dagger = (e^{-i p_B \cdot x} | B \rangle)^\dagger = \langle B | e^{i p_B \cdot x}}$$

$$= \partial_\mu \left[e^{i \overbrace{(p_B - p_A) \cdot x}^q} \langle B | A_i^\mu(0) | A \rangle \right]$$

$$= e^{i q \cdot x} i q_\mu \langle B | A_i^\mu(0) | A \rangle$$

$$= e^{i q \cdot x} i \langle B | q_\mu A_i^\mu(0) | A \rangle$$

Thus

$$\underbrace{\langle B | \partial_\mu A_i^\mu(x) | A \rangle}_{m_q P_i(x)} = e^{i q \cdot x} i \langle B | q_\mu A_i^\mu(0) | A \rangle$$

$$m_q \langle B | P_i(x) | A \rangle = i e^{i q \cdot x} \langle B | q_\mu A_i^\mu(0) | A \rangle$$

Similarly using the rotational invariance for

$$P_i(x) = e^{i\hat{p}\cdot x} P_i(0) e^{-i\hat{p}\cdot x} \quad \text{we get}$$

$$m_q e^{i q \cdot x} \langle B | P_i(0) | A \rangle = e^{i q \cdot x} i \langle B | q_\mu A_i^\mu(0) | A \rangle$$

$$h) \langle N(p') | A_i^\mu(0) | N(p) \rangle = \bar{u}(p') \left[\gamma^\mu G_A(t) + \frac{(p'-p)^\mu}{2m_N} G_P(t) \right] \gamma_5 \frac{\tau_i}{2} u(p)$$

Now

$$\frac{M_\pi^2 \bar{F}_\pi}{M_\pi^2 - t} G_{\pi N}(t) = i \bar{u}(p') \gamma_5 \tau_i u(p)$$

$$= m_q \langle N(p') | P_i(0) | N(p) \rangle$$

$$= i q_\mu \langle N(p') | A_i^\mu(0) | N(p) \rangle$$

$$= i \bar{u}(p') \left[(p'-p) G_A(t) + \frac{(p'-p)^2}{2m_N} G_P(t) \right] \gamma_5 \frac{\tau_i}{2} u(p)$$

use DE: $\bar{u}(p') \not{p}' = m_N u(p)$

$\not{p} u(p) = m_N u(p)$

and $\not{p} \gamma_5 = -\gamma_5 \not{p}$

$$= i \bar{u} \left[(m_N + m_N) G_A(t) + \frac{t}{2m_N} G_P(t) \right] \gamma_5 \frac{\tau_i}{2} u(p)$$

$$= i \bar{u}(p') \frac{1}{2} \left[2m_N G_A(t) + \frac{t}{2m_N} G_P(t) \right] \gamma_5 \tau_i u(p)$$

Thus

$$\frac{M_\pi^2 \bar{F}_\pi}{M_\pi^2 - t} G_{\pi N}(t) = m_N G_A(t) + \frac{t}{4m_N} G_P(t)$$