

Particle Physics, Part 2

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Recall from Wednesday

- “Diracology”: calculate spin-summed and averaged squared amplitude $\overline{|\mathcal{M}|^2}$
- Sum over final state spins
- Average over initial state spins
- Leptonic tensors $L_{\mu\nu}$, write as traces, use $\sum_s u(p)\bar{u}(p) = \not{p} + m$
- Traces of gamma matrices
- QED \rightarrow QCD: Invariant in local $SU(3)$ transforms $\exp(i \sum_{a=1}^8 \alpha^a t^a)$
- $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + ig[A_\mu, A_\nu]$ is 3x3 matrix, $A_\mu = A_\mu^a t^a$, gluon color a
- Quark spinors also carry a 3-component color vector
- QCD Lagrangian

Quiz

Two-jet production in photon-proton scattering can be used to probe the quark and gluon densities. Draw some diagrams that produce two jets (remember: quarks and gluons become jets) and probe the gluon density of the proton:

$$\gamma + g \rightarrow q + \bar{q}$$

How can you produce two jets in photon + quark scattering?

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How can you produce two jets in photon + quark scattering?

Photon $\rightarrow q\bar{q}$ splitting, then quark and/or (lowest order: or) antiquark couples to the gluon from the proton.

A quark can absorb the photon, and emit a gluon (γ and g not identical, no cross diagram!)

Quiz

Consider a process $q(i)\bar{q}(j) \rightarrow q(k)\bar{q}(l)$ via an s channel gluon. Here i, j, k, l refer to the colors. Calculate the color factor C (which multiplies the invariant amplitude $|\mathcal{M}|^2$).

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Summation over i, j, k, l and gluon color indices a, b is implicit. Note that the gluon propagator forces color at both ends to be the same. We use $(t_{ij}^a)^* = t_{ji}^a$, as $t^{a\dagger} = t^a$.

$$\begin{aligned} C &= \frac{1}{N_c^2} (t_{ji}^a t_{kl}^a) (t_{ji}^b t_{kl}^b)^* = \frac{1}{N_c^2} (t_{ji}^a t_{ij}^b) (t_{kl}^a t_{lk}^b) = \frac{1}{N_c^2} \text{Tr}(t^a t^b) \text{Tr}(t^a t^b) = \frac{1}{N_c^2} \frac{1}{2} \delta^{ab} \frac{1}{2} \delta^{ab} \\ &= \frac{1}{N_c^2} \frac{1}{2 \cdot 2} 8 = \frac{2}{9} \end{aligned}$$

Same result as in the t channel case.

We will discuss these shortly....

- Why it is possible to probe proton structure (distance scale \ll proton radius) in DIS?
- What is the interpretation (in the parton model) of Bjorken- x , defined as

$$x = \frac{Q^2}{2P \cdot q}$$

- If the structure functions satisfy $F_1 = \frac{1}{2x}F_2$, and $F_2 \sim$ number of quarks (Callan-Gross relation), why does it tell us that the quarks are spin- $\frac{1}{2}$ particles?

Recall from Monday

- QCD vertices: qqg , ggg and ggq
- Feynman rules, for qqg similar to $qq\gamma$ but an additional t_{ij}^a
- $a = 1, 2, 3, 4, 5, 6, 7, 8$ is the gluon color, and quarks carry a color index $i = 1, 2, 3$
- Cross section calculations in QCD: as in QED, but additional color factor (factorizes, can calculate separately)
- Coupling constants run, α_{em} large at large scales (small distances)
 $\alpha_s \rightarrow 0$ at large scales (asymptotic freedom)
- Determining α_s : QCD corrections to scattering processes at different Q scales
- Collider data: quarks are spin- $\frac{1}{2}$, gluons spin-1 (angular distributions), # of colors $N_c = 3$

Rich QCD phenomenology, we won't cover here in detail, read from the lecture notes!

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Uncertainty principle: virtual photon has momentum $\sim q$, wavelength $\sim 1/q$. At HERA, one reaches high $Q^2 = -q^2$, corresponding to $1/Q \sim 1/(100 \text{ GeV}) \sim 1/500 \text{ fm}$. Compare with proton radius $\sim 1 \text{ fm}$.

Lecture notes p. 334

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Fraction of the proton momentum carried by the quark participating in the scattering in DIS (in the frame where the proton has large longitudinal momentum).

Lecture notes p. 338

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If quarks are spin- $\frac{1}{2}$ fermions, the angular distribution in DIS should be the same as in $e + \mu \rightarrow e + \mu$ scattering (p. 336). The parametrized DIS cross section (p. 335) has this form, if $F_1 = \frac{1}{2x}F_2$, normalized by \sim number of quarks.

Quiz

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First calculate weak isospin, note that $u_L = \begin{pmatrix} u_L \\ 0 \end{pmatrix}$:

$$\hat{T}_3 u_L = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} u_L \\ 0 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 \cdot u_L \\ 0 \end{pmatrix},$$

so the weak isospin $T_3 = 1/2$. Thus, the weak hypercharge reads

$$Y = 2(Q_{\text{em}} - T_3) = 2 \left(\frac{2}{3} - \frac{1}{2} \right) = \frac{1}{3}.$$