

Points of stationary phase.

$$\dot{x} = \frac{p}{m} t' \quad |p| \leq m\eta \quad \text{on supports of } \rho \quad (18)$$

Then, if $|\frac{\dot{x}}{t}| \geq \eta$

$$|\check{f}_E(x)| \leq C_m (1 + |x| + |E|)^{-m}$$

$$|\check{f}_E(x)| \leq C_m (1 + |x| \eta^2 \rho + |E| \eta^2 \rho)^{-m} \eta^{4\rho}$$

$$\left| \frac{x \eta^2 \rho}{t \eta^2 \rho} \right| \geq \eta \Rightarrow |x| \geq \eta |E| \eta^2 \rho.$$

Finally

$$\int_{|x| \geq \eta |E| \eta^2 \rho} |\check{f}_E(x)|^{-\alpha} dx \leq C_m \eta^{4\rho} \int_{|x| \geq \eta |E| \eta^2 \rho} (1 + |x| \eta^2 \rho + |E| \eta^2 \rho)^{-m} dx$$

$$|x| \geq \eta |E| \eta^2 \rho$$

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and estimating the integral we obtain the result.

Corollary

For any $Q \in \mathbb{R}^2$, $f \in C_0^\infty(B_{m\eta})$, $0 \leq \rho < 1$,

$\rho = 1, 2, 3, \dots$ there is a C_ρ such that

$$\|F(\rho_i - Q - \sqrt{E}t) \rho^{\frac{1}{2}}\| \leq C_\rho \left(\frac{\rho - m\eta}{\eta \rho} \right) \quad (19)$$

$$F(|x - Q| \leq \frac{|\sqrt{E}t|}{\rho}) \leq C_\rho (1 + |\sqrt{E}t|)^{-\rho}$$

For $\nu > (\rho \eta)^{\frac{1}{\rho}}$ with $\nu = |\sqrt{E}t|$.

Proof

$$\cos(\tilde{M}, M + \sqrt{E}) - \eta \sqrt{E} |E| \geq \frac{|E|}{\rho} - \eta \sqrt{E} |E| \geq 0$$

$$x \in \tilde{M}, y \in M$$

$$|x - y - \sqrt{E}| = |x - Q - \sqrt{E} + (Q - y)|$$

$$\geq \frac{|E|}{\rho} - \frac{|E|}{\rho} = \frac{|E|}{\rho}$$

$$\frac{|E|}{\rho} \geq \eta \sqrt{E} |E| \Rightarrow \sqrt{E} \geq \rho \eta$$

Define $\delta = \frac{1}{\rho \eta} - 1$. Then,

$$I \leq C_\rho (1 + \rho \sqrt{E} + \eta \sqrt{E} |E|)^{-\rho}$$

$$\rho \sqrt{E} + \eta \sqrt{E} |E| = \sqrt{E} (\rho + \eta \sqrt{E} |E|) =$$

$$= \sqrt{E} \cos(\tilde{M}, M + \sqrt{E}) \geq \frac{|E|}{\rho}$$

Finally

$$I \leq C_\rho (1 + |E|)^{-\rho}$$

□

3 The Hamiltonian

The potential V satisfies.

Essential condition: Definition of V

V is $V_{rel} = 0$ -small, i.e. it is H_0 -bounded with relative bound zero (or just < 1)

$$\|V\psi\| \leq \epsilon \|H_0\psi\| + K\|\psi\|, \quad \forall \epsilon > 0, \text{ and}$$

$$\int_0^\infty dR \|V(x) (H_0 + I)^{-1} F(|x| > R)\| < \infty.$$

This is equivalent to the more intuitive condition

$$\int_0^\infty dR \|V(x) F(|x| > R) (H_0 + I)^{-1}\| < \infty.$$

Let us prove this

Denote

$$h(R, \beta) = \|V (H_0 + \beta)^{-1} F(|x| > R)\|, \quad \beta \in \mathbb{C} \setminus [0, \infty)$$

Take $\beta = \gamma$. $\gamma \in \mathbb{C}_0^\infty$, $\gamma(x) = 1, |x| > 1, \gamma(x) = 0, |x| \leq 1/2, 0 \leq \gamma \leq 1$

$$j_R(x) = \gamma(x/R), \quad j_R(x) = 1, |x| > R,$$

$$j_R(x) = 0, |x| \leq \frac{R}{2}.$$

$$h_1(R, \beta) = \|V (R_0 + \beta)^{-1} j_R\|$$

$$h_2(R, \beta) = \|V j_R (R_0 + \beta)\|$$

Then

↑ explain abuse of notation.

$$h_1(R, \beta) = \|F(|x| > R) (R_0 + \beta)^{-1} V\| \leq \|j_R (R_0 + \beta)^{-1} V\|$$

$$= \|V (R_0 + \beta)^{-1} j_R\| = h_2(R, \beta).$$

$$h_1(R, \beta) = \|j_R (R_0 + \beta)^{-1} V\| \leq \|F(|x| > R) (R_0 + \beta)^{-1} V\|$$

$$= h_2(R, \beta)$$

low,

$$h(R, \beta) \leq h_1(R, \beta) \leq h\left(\frac{R}{2}, \beta\right) \quad \text{and}$$

$$h(R, \beta) \in L^1(0, \infty) \Leftrightarrow h_2(R, \beta) \in L^1(0, \infty).$$

Moreover

$$|h_1(R, \beta) - h_2(R, \beta)| \leq \|V(R_0(\beta)) J_R - V J_R(R_0(\beta))\|$$

$$= \|V(R_0(\beta)) [J_R, H_0] R_0(\beta)\|$$

$$[J_R, H_0] = \frac{1}{2m} \left[\frac{1}{R^2} (\Delta J)_R + \frac{2}{R} (\nabla J)_R \cdot \nabla \right]$$

Then

$$|h_1(R, \beta) - h_2(R, \beta)| \leq \frac{1}{2m} \frac{1}{R^2} \|V(R_0(\beta)) (\Delta J)_R\| + \frac{1}{mR} \sum_{i=1}^n \|V(R_0(\beta)) (\nabla_i J)_R\| \| \partial_i R_0(\beta) \|$$

$$\leq \frac{C_2}{R^2} h_2\left(\frac{R}{2}, \beta\right) + \frac{C_3}{R} h_1\left(\frac{R}{2}, \beta\right) \leq \frac{C_3}{R} h_1\left(\frac{R}{2}, \beta\right)$$

It follows that

$$h_2(R, \beta) \in L^1(0, \infty) \Rightarrow h_2(R, \beta) \in L^1(0, \infty),$$

Furthermore,

$$|h_1(R, \beta) - h_2(R, \beta)| \leq \|V(R_0(\beta)) J_{R_2} [J_{R_1}, H_0] R_0(\beta)\|$$

define

$$\leq \|V [R_0(\beta), J_{R_2}] [J_{R_1}, H_0] R_0(\beta)\|$$

$$\leq \|V J_{R_2} R_0(\beta) [J_{R_1}, H_0] R_0(\beta)\| \leq$$

$$\leq \|V(R_0, \beta) [\int_{\beta_2}^{\beta_1} H_0] R_0, \beta) [\int_{\beta_2}^{\beta_1} H_0] R_0, \beta) \|$$

$$+ \frac{C_2}{R} h_2(R, \beta) \leq \frac{C_2}{R^2} + \frac{C_2}{R} h_2(R, \beta)$$

Then, $h_2(R, \beta) \in L^1 \Rightarrow h_2(R, \beta) \in L^1$

So, we have proved that

$$h_1(R, \beta) \in L^1(0, \infty) \Leftrightarrow h_2(R, \beta) \in L^1(0, \infty) \Leftrightarrow$$

$$h_2(R, \beta) \in L^1(0, \infty)$$

Define $h_3(R, \beta) = \| F(|x| > R) V(x, R, \beta) \|$
 We have that

$$h_2(R, \beta) \leq h_3(R, \beta), \text{ and also}$$

$$h_3(R, \beta) \leq h_2(R, \beta)$$

Hence

$$h_2(R, \beta) \in L^1(0, \infty) \Leftrightarrow h_3(R, \beta) \in L^1(0, \infty)$$

In consequence

$$h_1(R, \beta) \in L^1(0, \infty) \Leftrightarrow h_3(R, \beta) \in L^1(0, \infty)$$

Finally

$$|h_3(R, \beta_2) - h_3(R, \beta_1)|$$

$$\leq \| V(x) F(|x| > R) (R, \beta_2) - R, \beta_1) \|$$

$$= |\beta_2 - \beta_1| \| V(x) F(|x| > R) R, \beta_2) R, \beta_1) \|$$

$$\in C_2 h(R, \beta_r), \quad r=1,2.$$

This implies that

$$h_3(R, \beta_1) \in L^1(0, \infty) \Leftrightarrow h_3(R, \beta_2) \in L^1(0, \infty)$$

$\forall \beta_1, \beta_2 \in \mathbb{R} \setminus \{0, \infty\}$, and the same is true for $h(R, \beta)$.

By Kato-Rellich Theorem

$H = H_0 + V$ is self-adjoint on

$$D(H_0) = \mathbb{R}^2$$

Lemma

Suppose that $V \in C_2$ and that for some $0 < p < 1$ and all $f \in C_0^\infty(\mathbb{R}^n)$

$$(1+R)^p \|V(x)g(\bar{p})F(|x|, R)\| \in L^1(0, \infty)$$

Then, for any $f \in C_0^\infty(B_{M\eta})$ there is a function h with $(1+R)^p h(R) \in L^1(0, \infty)$ such that for every $\bar{v} \in \mathbb{R}^n$ satisfying

$$|\bar{v}| = \eta, \quad |\bar{v}| = |\bar{v}'|, \quad \text{we have}$$

$$\|V(x) e^{itH_0} f(\bar{p} - \eta \bar{v}) (1+x^2)^{-3/2}\|$$

$$= \|V(x + t\bar{v}) e^{itH_0} f(\bar{p}) (1+x^2)^{-3/2}\| \leq h(|t\bar{v}|).$$

Proof:

$$I = \int_{\mathbb{R}^n} |V(x)| e^{i t h_0} f(\bar{p} - m\bar{v}) (1+x^2)^{-3/2} dx$$

$$= \int_{\mathbb{R}^n} |V(x)| e^{i t h_0} e^{i m\bar{v} \cdot x} f(\bar{p}) e^{-i m\bar{v} \cdot x} (1+x^2)^{-3/2} dx$$

$$= \int_{\mathbb{R}^n} |V(x)| e^{-i \bar{p} \cdot \bar{v} t} e^{i t h_0} f(\bar{p}) (1+x^2)^{-3/2} dx$$

$$= |V(x + \bar{v} t)| e^{i t h_0} f(\bar{p}) (1+x^2)^{-3/2}$$

Take $g \in C_0^\infty(\mathbb{R}^n)$ with $g \equiv 1$ on the support of f . Then

$$\| |V(x)| g(\bar{p} - m\bar{v} t) \| = \| |V(x)| e^{i m\bar{v} \cdot x} g(\bar{p}) e^{-i m\bar{v} \cdot x} \|$$

$= \| |V(x)| g(\bar{p}) \|$ is independent of t

Recall our Corollary:

$$\| F((x - \bar{v} t) \cdot \frac{1}{4} \frac{1}{h_0}) e^{i t h_0} g(\bar{p} - m\bar{v} t) F((x) \cdot \frac{1}{8} \frac{1}{h_0}) \|$$

$$\leq C (1 + |t|)^{-2}$$

$$I = \int_{\mathbb{R}^n} |V(x)| g(\bar{p} - m\bar{v} t) \left[F((x - \bar{v} t) \cdot \frac{1}{4} \frac{1}{h_0}) + F((x - \bar{v} t) \cdot \frac{1}{4} \frac{1}{h_0}) \right]$$

$$e^{i t h_0} f(\bar{p} - m\bar{v} t) \left(F(|x| \leq \frac{1}{8} \frac{1}{h_0}) + F(|x| \geq \frac{1}{8} \frac{1}{h_0}) \right)$$

$$(1+x^2)^{-3/2} \leq I_1 + I_2 + I_3$$

$$I_1 = \int_{|x-\bar{u}| \geq \frac{1}{4}|\bar{u}|} V(x) g(|\bar{p}-m\bar{v}|) e^{-\epsilon h_0} f(|\bar{p}-m\bar{v}|) F(|x-\bar{u}| \geq \frac{1}{4}|\bar{u}|) (1+x^2)^{-3/2} dx$$

$$I_2 = \int_{|x-\bar{u}| \leq \frac{1}{4}|\bar{u}|} V(x) g(|\bar{p}-m\bar{v}|) e^{-\epsilon h_0} f(|\bar{p}-m\bar{v}|) F(|x-\bar{u}| \leq \frac{1}{4}|\bar{u}|) (1+x^2)^{-3/2} dx$$

$$I_3 = \int_{|x-\bar{u}| \geq \frac{1}{4}|\bar{u}|} V(x) g(|\bar{p}-m\bar{v}|) e^{-\epsilon h_0} f(|\bar{p}-m\bar{v}|) F(|x-\bar{u}| \geq \frac{1}{4}|\bar{u}|) (1+x^2)^{-3/2} dx$$

I₁ The classically forbidden region

$$I_1 \leq \int V(x) g(|\bar{p}-m\bar{v}|) e^{-\epsilon h_0} (1+|\bar{u}|)^{-\ell} dx, \quad \ell=1, 2, 3$$

I₂ Classically allowed region.
Here the potential is small.

$$I_2 \leq C \int V(x) g(|\bar{p}-m\bar{v}|) F(|x-\bar{u}| \leq \frac{1}{4}|\bar{u}|) dx$$

$$\leq \int V(x) g(|\bar{p}-m\bar{v}|) F(|x-\bar{u}| \leq \frac{3}{4}|\bar{u}|) dx := h_2(|\bar{u}|)$$

with $(1+|\bar{u}|)^{\ell} h_2(|\bar{u}|) \in L^1(0, \infty)$

I₃ The weight is small

$$I_3 \leq C (1+|\bar{u}|)^{-3/2}$$

□