

2. The Free Propagation

(13)

We denote by $\vec{p} = -i\nabla$ the operator of momentum.

Then,

$$H_0 = \frac{-\Delta}{2m} = \frac{\vec{p}^2}{2m} \quad \mathcal{D}(H_0) = \mathcal{L}^2$$

We define operators $f(\vec{p})$, for Borel function f by functional calculus or as $f(\vec{p}) = \mathcal{F}^{-1} f(\vec{p}) \mathcal{F}$ where \mathcal{F} denotes the Fourier transform

$$(\mathcal{F}g)(\vec{p}) = \frac{1}{(2\pi)^{n/2}} \int e^{i\vec{p}\cdot\vec{x}} g(\vec{x}) d\vec{x}$$

Note that

$$e^{i\vec{p}\cdot\vec{v}t} f(\vec{p}) e^{-i\vec{p}\cdot\vec{v}t} = f(\vec{p} + \vec{v}t) \quad (1)$$

$$e^{im\vec{v}\cdot\vec{x}} f(\vec{p}) e^{-im\vec{v}\cdot\vec{x}} = f(\vec{p} + m\vec{v}) \quad (2)$$

In particular

$$e^{im\vec{v}\cdot\vec{x}} e^{i\vec{p}\cdot\vec{x}} e^{-im\vec{v}\cdot\vec{x}} = e^{im\vec{v}^2/2} e^{i\vec{p}\cdot\vec{x}} e^{i\vec{p}\cdot\vec{x}} \quad (3)$$

$$\vec{v} = \vec{v}t$$

Proof of (2):

$$e^{im\vec{v}\cdot\vec{x}} f(\vec{p}) e^{-im\vec{v}\cdot\vec{x}} g(\vec{x}) = e^{im\vec{v}\cdot\vec{x}} \mathcal{F}^{-1} f(\vec{p}) \mathcal{F} e^{-im\vec{v}\cdot\vec{x}} g(\vec{x})$$

$$\mathcal{F}^{-1} e^{im\vec{v}\cdot\vec{x}} g = \frac{1}{(2\pi)^{n/2}} \int e^{i\vec{p}\cdot\vec{x}} e^{im\vec{v}\cdot\vec{x}} g(\vec{x}) d\vec{x}$$

$$= \hat{\mathcal{F}}(p - m\bar{v}), \quad \hat{\mathcal{F}} = \mathcal{F} \circ \tau_{\bar{v}} \quad (14)$$

$$e^{i m \bar{v} \cdot x} f(\bar{p}) = e^{i m \bar{v} \cdot x} \mathcal{F}(\lambda) = \frac{1}{(2\pi)^{n/2}} e^{i m \bar{v} \cdot x} \int e^{i p \cdot x} f(p)$$

$$\hat{\mathcal{F}}(p - m\bar{v}) dp = \frac{1}{(2\pi)^{n/2}} \int e^{i(p - m\bar{v}) \cdot x} f(p - m\bar{v} + m\bar{v})$$

$$\hat{\mathcal{F}}(p - m\bar{v}) dp = \mathcal{F}^{\wedge} f(p + m\bar{v}) \quad \mathcal{F}^{\wedge} f = f(\bar{\cdot} + m\bar{v}) \mathcal{F}$$

Note that

$$\mathcal{F} e^{i m \bar{v} \cdot x} f(x) = \hat{\mathcal{F}}(p - m\bar{v})$$

$e^{i m \bar{v} \cdot x}$ gives a "boost" $m\bar{v}$ to the state \mathcal{F}



$e^{i m \bar{v} \cdot x} f(x)$ is a state with velocity $\approx \bar{v}$

Notation

$B_{m\eta}(m\bar{v})$: open ball of center $m\bar{v}$ and radius $m\eta$.

Lemma

For any $f \in C_0^\infty(B_{m\eta})$, for some $\eta > 0$, and any $l = 1, 2, \dots$, there is a constant C_l such

$$\| F(x \in \tilde{M}) e^{i \epsilon H_0} f\left(\frac{p - m\tilde{v}}{v\beta}\right) F(x \in M) \| \quad (14)$$

$$\leq C_\epsilon (1 + \epsilon v \beta + \eta v^2 \beta |\epsilon|)^{-\ell}$$

For every $\tilde{v} \in \mathbb{R}^n$, $\epsilon \neq 0$, $v > 0$, $\beta \in \mathbb{R}$ and any measurable sets $M, \tilde{M} \subset \mathbb{R}^n$ such that $\epsilon := \text{dist}(M, M + \tilde{v}\epsilon) - \eta v \beta |\epsilon| > 0$.

Before we give the proof we try to understand the meaning of the Lemma.

For simplicity take $\beta = 0$.

$$\| F(x \in \tilde{M}) e^{i \epsilon H_0} f\left(\frac{p - m\tilde{v}}{v\beta}\right) F(x \in M) \|$$

$$\leq C_\epsilon (1 + \epsilon v \beta |\epsilon|)^{-\ell}$$

$$\text{dist}(M, M + \tilde{v}\epsilon) \geq \eta v \beta |\epsilon|$$

other side.

f localizes on $|p - m\tilde{v}| \leq m \eta$
 the velocity differs from \tilde{v} by less than η

Proof.

$$\| F(x \in \tilde{M}) e^{i \epsilon H_0} f\left(\frac{p - m\tilde{v}}{v\beta}\right) F(x \in M) \| =$$

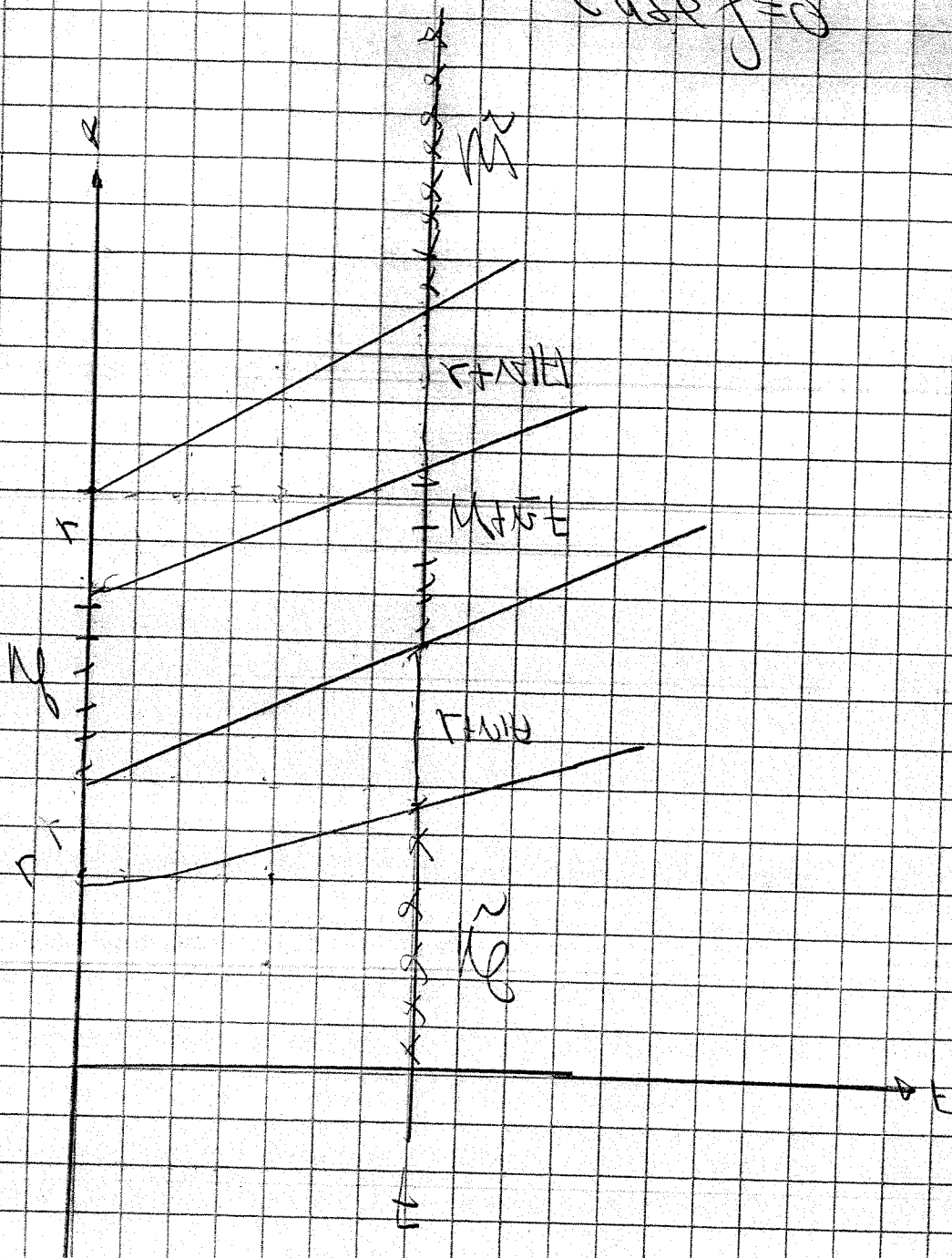
$$= \| F(x \in \tilde{M}) e^{i \epsilon H_0} e^{i m \tilde{v} \cdot x} f\left(\frac{p}{v\beta}\right) e^{-i m \tilde{v} \cdot x} F(x \in M) \|$$

$$= \| F(x \in \tilde{M}) e^{i m \tilde{v} \cdot x} e^{i \epsilon H_0} e^{-i m \tilde{v} \cdot x} f\left(\frac{p}{v\beta}\right) F(x \in M) \|$$

$$= \| F(x \in \tilde{M}) e^{i \epsilon H_0} f\left(\frac{p}{v\beta}\right) e^{-i \tilde{v} \cdot \tilde{v} \epsilon} F(x \in M) e^{i \tilde{v} \cdot \tilde{v} \epsilon} \|$$

Case $f=0$

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$$\text{dist}(\vec{M}, M + \vec{E}) = r + |E|, \quad r \geq 0$$

$$= \| F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \|^2 \quad (18)$$

$$= \| F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M + \vec{v}\epsilon) \|^2$$

Replace $M \rightarrow M + \vec{v}\epsilon$, and we have to prove that

$$\| F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \|^2 \leq$$

$$\leq C_e (1 + \gamma v_S + \eta v_S^2 |\epsilon|)^{-\rho},$$

$$r := \text{const}(\tilde{M}, M) - \eta v_S |\epsilon| \geq 0.$$

We have that

$$= \| F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \Phi \|^2 =$$

$$= \left(\int_{\tilde{M}} F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \Phi \right)$$

$$F(x \in \tilde{M}) e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \Phi \Big|_{\tilde{M}}$$

$$e^{i\epsilon H_0} f\left(\frac{\vec{p}}{v_S}\right) F(x \in M) \Phi = \frac{1}{(2\pi)^{n/2}} \int_{\tilde{M}} \tilde{f}_k(x=y) \Phi(y) dy$$

$$\tilde{f}_k(x) = \frac{1}{(2\pi)^{n/2}} \int e^{i\vec{p}\cdot\vec{x}} e^{-i\vec{p}\cdot\vec{z}/2m} f\left(\frac{\vec{p}}{v_S}\right) dz.$$

$$I = \frac{1}{(2\pi)^n} \int_{\tilde{M}} dx \int_M dy \tilde{f}_k(x-y) \Phi(y) dy.$$

$$\int_M dz \tilde{f}_k(x-z) \Phi(z) dz.$$

$$\text{But } \int_{\mathbb{R}^n} |\phi(y)| |\phi(z)| \leq |\phi(y)| + |\phi(z)| \quad (1)$$

$$I \leq \frac{1}{2} \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \left[|\tilde{f}_E(x-y)| |\tilde{f}_E(x-z)| (|\phi(y)|^2 + |\phi(z)|^2) \right] dx dy dz$$

Maxim

$$|x-y| \geq \text{const}(M, \tilde{w}) = \frac{r}{2} + \eta \sqrt{S|E|}$$

replace $x-y = \tilde{x}$, $x-z = \tilde{z}$ $y \rightarrow y$ and $x-y = \tilde{x}$ $x-z = \tilde{x}$ $z = z$ and we get

$$I \leq \frac{1}{(2\pi)^n} \left[\int_{\mathbb{R}^n} |\tilde{f}_E(x)|^2 dx \right]^2 \|\phi\|^2$$

Max, $r + \eta \sqrt{S|E|}$

We have to estimate $\tilde{f}_E(x)$ using (non) stationary phase

Hörmander, Reed-Simon III Corollary to Theorem XI.14, page 38.

$$\tilde{f}_E(x) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i p \cdot x} e^{-i \frac{p^2}{2m} t} f\left(\frac{p}{\sqrt{S}}\right) dp$$

$$\tilde{f}_E(x) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i p \cdot x \sqrt{S}} e^{-i \frac{p^2}{2m} t \sqrt{S}^2} f(p) dp$$

$$\tilde{f}_E(x) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i p \cdot x} e^{-i \frac{p^2}{2m} t} f(p) dp$$

$$x = x \sqrt{S} \quad t = t \sqrt{S}^2$$

Points of stationary phase.

$$\dot{x} = \frac{p}{m} t' \quad |p| \leq m\eta \quad \text{on supports of } \tilde{f} \quad (18)$$

Then, if $|\frac{\dot{x}}{t'}| \geq \eta$

$$|\tilde{f}_E(x)| \leq C_m (1 + |x| + |E|)^{-m}$$

$$|\tilde{f}_E(x)| \leq C_m (1 + |x| \eta^2 \delta + |E| \eta^2 \delta)^{-m} \eta^{4p}$$

$$\left| \frac{x \eta^2 \delta}{t' \eta^2 \delta} \right| \geq \eta \Rightarrow |x| \geq \eta |E| \eta^2 \delta.$$

Finally

$$\int_{|x| \geq \eta |E| \eta^2 \delta} |\tilde{f}_E(x)|^{-\alpha} dx \leq C_m \eta^{4p} \int_{|x| \geq \eta |E| \eta^2 \delta} (1 + |x| \eta^2 \delta + |E| \eta^2 \delta)^{-m} dx$$

$$|x| \geq \eta |E| \eta^2 \delta$$

$$|x| \geq \eta |E| \eta^2 \delta$$

and estimating the integral we obtain the result.

Corollary

For any $Q \in \mathbb{R}^2$, $f \in C_0^\infty(B_{m\eta})$, $0 \leq \rho < 1$,

$\rho = 1, 2, 3, \dots$ there is a C_ρ such that

$$\|F(\rho_i - Q - \sqrt{E}t) \otimes \tilde{f}(\frac{p - m\dot{v}}{\eta^2 \delta})\| \leq C_\rho$$

$$\|F(|x - Q| \leq \frac{|\sqrt{E}t|}{\delta})\| \leq C_\rho (1 + |\sqrt{E}t|)^{-\rho}$$

For $\nu > (\rho \eta)^{\frac{1}{\rho}}$ with $\nu = |\sqrt{E}t|$.